M1 "Sciences et Génies des Matériaux" & M1 franco-allemand "Polymères"

## Quantum Mechanics course

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The following questions are meant to help the students in reviewing their lecture notes.

- 1. Which mathematical functions are used for describing the state of a particle moving along the x axis in classical Newton and quantum mechanics, respectively? Write the fundamental time-dependent equations that these functions are supposed to fulfill.
- 2. Let  $\begin{bmatrix} \hat{A} \end{bmatrix} \equiv \{A_{ij}\}_{1 \le i,j \le N}$  be the matrix representation of a *linear* quantum operator  $\hat{A}$  in a given basis  $\{|u_i\rangle\}_{1 \le i \le N}$ . Show that, according to the linearity of  $\hat{A}$ , we can deduce from  $\begin{bmatrix} \hat{A} \end{bmatrix}$  the quantum state  $\hat{A} |\Psi\rangle$  that is obtained by applying  $\hat{A}$  to any state  $|\Psi\rangle = \sum_{j=1}^{N} C_j |u_j\rangle$ . Show that, if  $\{|u_i\rangle\}_{1 \le i \le N}$ is orthonormal, then  $C_i = \langle u_i |\Psi\rangle$  and  $A_{kj} = \langle u_k |\hat{A}| u_j\rangle$ . Deduce that the identity operator reads  $\hat{\mathbb{1}} = \sum_{j=1}^{N} |u_j\rangle \langle u_j|$  and that  $\hat{A} = \hat{\mathbb{1}}\hat{A}\hat{\mathbb{1}} = \sum_{i,j=1}^{N} A_{ij} |u_i\rangle \langle u_j|$ .
- 3. How does the wave function  $\Psi(\vec{r})$ , where  $\vec{r} \equiv (x, y, z)$  denotes a position that the particle under study could have in the three-dimensional real space, relate to the quantum state  $|\Psi\rangle$  of that particle? Explain the notation  $\Psi(\vec{r}') = \langle \vec{r}' | \Psi \rangle$ , where  $\vec{r}' \equiv (x', y', z')$  is some given fixed position. Let  $\hat{x}$  denote the *x* coordinate position operator. Why, according to the postulates of quantum mechanics, should we expect the following equality to be fulfilled:  $\hat{x} | \vec{r}' \rangle = x' | \vec{r}' \rangle$ ? Deduce that  $(\hat{x}\Psi) (\vec{r}') = \langle \vec{r}' | \hat{x} | \Psi \rangle = x' \times \Psi(\vec{r}')$ .
- 4. Let  $\hat{\mathcal{O}}$  denote a Hermitian quantum operator and  $|\Psi_{\mathfrak{G}_i}\rangle$  an eigenvector of  $\hat{\mathcal{O}}$  associated to the eigenvalue  $\mathfrak{o}_i$  (which is indexed with *i*). Show that  $\mathfrak{o}_i$  is a real number (*i.e.*, it is equal to its complex conjugate). Show that, if  $|\Psi_{\mathfrak{G}_j}\rangle$  is another eigenvector of  $\hat{\mathcal{O}}$  associated to a different eigenvalue  $\mathfrak{o}_j \neq \mathfrak{o}_i$ , then  $|\Psi_{\mathfrak{G}_j}\rangle$  and  $|\Psi_{\mathfrak{G}_i}\rangle$  are orthogonal. On the basis of these mathematical properties, explain why, in quantum mechanics, we associate to any observable  $\mathcal{O}$  a Hermitian operator  $\hat{\mathcal{O}}$ . We assume that, just before measuring  $\mathcal{O}$ , the system under study is in the quantum state  $|\Psi\rangle = \sum_k C_k |\Psi_{\mathfrak{G}_k}\rangle$ . What is the probability of measuring  $\mathfrak{o}_i$ ? Why should we always verify that  $|\Psi\rangle$  is normalized before evaluating that probability? If  $|\Psi\rangle$  is not normalized, how can we construct a normalized quantum state that is equivalent to  $|\Psi\rangle$ ?

- 5. Let  $\hat{A}$  and  $\hat{B}$  be two Hermitian operators. We assume that they share a *common* orthonormal basis of eigenvectors  $\{|u_i\rangle\}$ . We denote  $a_i$  and  $b_i$  the eigenvalues of  $\hat{A}$  and  $\hat{B}$ , respectively, that are associated to  $|u_i\rangle$ . Show that  $\hat{A}$  and  $\hat{B}$  commute, *i.e.*,  $[\hat{A}, \hat{B}] |u_i\rangle = 0$  for any basis vector  $|u_i\rangle$ . Explain why, in this case, the observables A and B (that are associated to  $\hat{A}$  and  $\hat{B}$ , respectively) can be measured simultaneously. If the value  $a_j$  is measured for A, which value will then be measured for B?
- 6. We represent the Hamiltonian of a quantum student in the two-state orthonormal basis  $\{|\mathfrak{O}\rangle, |\mathfrak{O}\rangle\}$  of "happy" and " not happy" states as follows:  $\begin{bmatrix} \hat{H} \end{bmatrix} = \begin{bmatrix} \alpha & \beta \\ \beta & \alpha \end{bmatrix}$ , where  $\beta < 0$ . Are  $|\mathfrak{O}\rangle$  and  $|\mathfrak{O}\rangle$  stationary states? Verify that the normalized eigenvectors of  $\hat{H}$  read  $\frac{1}{\sqrt{2}}(|\mathfrak{O}\rangle + |\mathfrak{O}\rangle)$  and  $\frac{1}{\sqrt{2}}(|\mathfrak{O}\rangle |\mathfrak{O}\rangle)$ . What is the probability of being happy for a quantum student that stands in its ground state?
- 7. Let  $\hat{H}$  denote the Hamiltonian operator of a quantum system and  $\hat{U}(t) = e^{-i\hat{H}t/\hbar}$ , where t denotes the time and  $i^2 = -1$ . We recall that, for any quantum operator  $\hat{A}$ , the exponential of  $\hat{A}$  reads  $e^{\hat{A}} \equiv \sum_{n=0}^{+\infty} \frac{\hat{A}^n}{n!}$ . We consider an orthonormal basis  $\{|\Psi_j\rangle\}$  of eigenvectors of  $\hat{H}$  and denote  $\{E_j\}$  the associated energies. Show that  $\hat{H} = \hat{H}\hat{\mathbb{1}} = \sum_k E_k |\Psi_k\rangle \langle \Psi_k|$  and  $\hat{U}(t) = \hat{U}(t)\hat{\mathbb{1}} = \sum_j e^{-iE_jt/\hbar} |\Psi_j\rangle \langle \Psi_j|$ .

Deduce that  $\hat{H}\hat{U}(t) = i\hbar \frac{d\hat{U}(t)}{dt}$  and conclude that  $|\Psi(t)\rangle = \hat{U}(t) |\Psi(t=0)\rangle$  is the quantum state of the system at time t if, at time t = 0, it is in the state  $|\Psi(t=0)\rangle$ . Why is  $\hat{U}(t)$  referred to as time evolution operator? Why are the eigenvectors of  $\hat{H}$  referred to as stationary states?

- 8. What is the general idea behind perturbation theory? How do we technically derive the perturbation expansion of the energies for a given Hamiltonian  $\hat{H}$ ?
- 9. Let  $\hat{H}_0$  and  $\hat{H}$  be two different Hamiltonian operators. We denote  $\hat{W} = \hat{H} \hat{H}_0$  and introduce the  $\alpha$ -dependent Hamiltonian  $\hat{H}(\alpha) = \hat{H}_0 + \alpha \hat{W}$ , where  $\alpha$  is a real number. Verify that  $\hat{H}(\alpha)$  is Hermitian. We denote  $E_i(\alpha)$  the eigenvalue of  $\hat{H}(\alpha)$  associated to the eigenvector  $|\Psi_i(\alpha)\rangle$  which is normalized for any value of  $\alpha$ . Prove the Hellmann–Feynman theorem:  $\frac{dE_i(\alpha)}{d\alpha} = \left\langle \Psi_i(\alpha) \middle| \frac{\partial \hat{H}(\alpha)}{\partial \alpha} \middle| \Psi_i(\alpha) \right\rangle$ . Deduce that, in perturbation theory, the first-order correction to the energy is determined solely from the perturbation operator  $\hat{W}$  and the unperturbed quantum state  $|\Psi_i(\alpha = 0)\rangle$ .