

**40-minute exam: Density functional theory of extended systems**

- a) **[2.5 pts]** We consider, for simplicity, the expression in *one dimension* of the Kohn–Sham (KS) potential at the local density level of approximation (LDA):  $v_{\text{LDA}}^{\text{KS}}(x) \equiv v(x) + \left. \frac{\partial e_{\text{xc}}(n)}{\partial n} \right|_{n=n(x)} + \int dx' \frac{n(x')}{|x' - x|}$ , where  $v(x)$  is the nuclear potential with periodicity  $a$  [i.e.,  $v(x + a) = v(x)$ ] of the extended system under study. We know from Bloch’s theorem that, when the two-electron repulsion is neglected, the (one-electron) solutions to the Schrödinger equation can be written as  $\varphi(x) \equiv \varphi_{j,k}(x) = e^{ikx} \times u_{j,k}(x)$  where  $k$  is a given wave number value,  $j$  is a quantum number, and  $u_{j,k}(x)$  is a function of the position  $x$  that has the periodicity  $a$  of the lattice. Does Bloch’s theorem also hold for the KS orbitals? **Hint:** Explain why the density  $n(x) \equiv \sum_{j,k}^{\text{occupied}} |\varphi_{j,k}(x)|^2$  is periodic. Deduce the periodicity of the KS-LDA potential and conclude.
- b) **[3 pts]** The fundamental energy gap of an  $N$ -electron system is evaluated from its neutral ( $N$ ), anionic ( $N + 1$ ), and cationic ( $N - 1$ ) ground-state energies as follows:  $E_g = E_{N+1} + E_{N-1} - 2E_N$ . Modelling the fundamental gap of materials accurately is of primary importance for condensed matter physicists. Why? Show that, in the particular case of non-interacting electrons, the fundamental gap, the so-called optical gap, and the HOMO-LUMO gap  $\varepsilon_{\text{L}} - \varepsilon_{\text{H}}$  boil down to the same quantity. Does this statement hold true for real (interacting) electronic systems? Why?
- c) **[2.5 pts]** It can be shown that the true fundamental gap can be evaluated exactly from the KS gap as follows,

$$E_g = \varepsilon_{\text{L}}^{\text{KS}} - \varepsilon_{\text{H}}^{\text{KS}} + \left[ \frac{\partial E_{\text{xc}}^{\xi}[n]}{\partial \xi} \right]_{n=n_N} \Big|_{\xi=0}, \quad (1)$$

where  $n_N$  denotes the  $N$ -electron ground-state density and  $E_{\text{xc}}^{\xi}[n]$  is the so-called  $N$ -centered ensemble exchange-correlation (xc) density functional. The latter describes the xc energy contribution to the  $N$ -centered ensemble energy  $E^{\xi} = \xi E_{N+1} + \xi E_{N-1} + (1 - 2\xi)E_N$ , where the ensemble weight  $\xi$  varies in the range  $0 \leq \xi \leq 1/2$ . Do standard xc functionals like LDA or PBE have an explicit dependence in  $\xi$ ? Can this be problematic for evaluating the fundamental gap of semiconductor materials, for example?

- d) **[3 pts]** We define as  $E_c^{\xi}[n] \approx 0$  and  $E_x^{\xi}[n] \approx (B + C\xi) \int d\mathbf{r} n^{\beta}(\mathbf{r})$  the ensemble *exchange-only* LDA functional, where the exponent  $\beta$  as well as the constant coefficients  $B$  and  $C$  are *a priori* unknown. Explain how we could determine  $B$  and  $\beta$  by considering the  $\xi = 0$  limit of the theory. The coefficient  $C$  might be determined by applying Eq. (1) to a simple system like the helium atom. How would you proceed? Would the resulting functional be applicable to extended systems? In your opinion, how useful/successful would such a functional be?