Second quantization

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ISTPC 2017 summer school, Aussois, France.

Electronic Hamiltonian in first quantization

***N*-electron** Hamiltonian within the Born–Oppenheimer approximation:

$$\hat{H} = \hat{T} + \hat{V}_{\rm ne} + \hat{W}_{\rm ee}$$

$$\hat{T} = \sum_{i=1}^{N} \hat{t}(i) \qquad \text{where} \quad \hat{t}(i) \equiv -\frac{1}{2}\nabla^2_{\mathbf{r}_i} \qquad \rightarrow \quad \text{kinetic energy}$$

$$\hat{V}_{\rm ne} = \sum_{i=1}^{N} \hat{v}_{\rm ne}(i) \qquad \text{where} \quad \hat{v}_{\rm ne}(i) \equiv -\sum_{A}^{\rm nuclei} \frac{Z_A}{|\mathbf{r}_i - \mathbf{R}_A|} \times \qquad \rightarrow \quad \text{electron-nuclei attraction}$$

$$\hat{W}_{\rm ee} = \frac{1}{2}\sum_{i\neq j}^{N} \hat{w}_{\rm ee}(i,j) \qquad \text{where} \quad \hat{w}_{\rm ee}(i,j) \equiv \frac{1}{|\mathbf{r}_i - \mathbf{r}_j|} \times \qquad \rightarrow \quad \text{electron-electron repulsion}$$

One-electron wavefunction

- Let us start with Schrödinger theory: the quantum state of a single electron is written as

$$|\Psi\rangle = \int \mathrm{d}\mathbf{r}\, \Psi(\mathbf{r})|\mathbf{r}\rangle$$

where $\Psi(\mathbf{r})$ is the one-electron wavefunction (orbital) and $|\mathbf{r}\rangle$ denotes the quantum state "the electron is at position $\mathbf{r}$". This choice of basis is known as "**r** representation".

- Orthonormalization condition: $\langle \mathbf{r}'|\mathbf{r}\rangle = \delta(\mathbf{r}' - \mathbf{r})$ ⟵ Dirac distribution

where "$\forall f$", $\int \mathrm{d}\mathbf{r}\, f(\mathbf{r})\delta(\mathbf{r}' - \mathbf{r}) = f(\mathbf{r}')$ $\leftrightarrow$ $\sum_i f_i \delta_{ii'} = f_{i'}$ ⟵ Kronecker delta

- Pauli theory: the spin of the electron is now considered as an additional degree of freedom. The quantum state of a single electron is then written as

$$|\Psi\rangle = \int \mathrm{d}\mathbf{r} \sum_{\sigma=\alpha,\beta} \Psi(\mathbf{r},\sigma)|\mathbf{r},\sigma\rangle$$

where $|\mathbf{r},\alpha\rangle$ denotes the quantum state "electron at position $\mathbf{r}$ with spin up" and $|\mathbf{r},\beta\rangle$ corresponds to the state "electron at position $\mathbf{r}$ with spin down".

Two-electron wavefunction

- In the non-relativistic case, a single electron will have a spin σ_0 which is either up or down. The corresponding wavefunction Ψ_{σ_0} can then be written as a spin-orbital $\Psi_{\sigma_0}(\mathbf{r},\sigma) = \Psi(\mathbf{r})\delta_{\sigma\sigma_0}$.

- With the notations $X \equiv (\mathbf{r},\sigma)$ and $\int \mathrm{d}X \equiv \int \mathrm{d}\mathbf{r} \sum_{\sigma=\alpha,\beta}$,

  a one-electron quantum state in Pauli theory is simply written as

$$|\Psi\rangle = \int \mathrm{d}X\, \Psi(X)|X\rangle$$

  with $\langle X'|X\rangle = \delta(X-X') = \delta_{\sigma\sigma'}\delta(\mathbf{r}'-\mathbf{r})$ and, consequently, $\langle X'|\Psi\rangle = \Psi(X')$.

- Two-electron case:

$$|\Psi\rangle = \int\int \mathrm{d}X_1\mathrm{d}X_2\, \Psi(X_1,X_2)|1\colon X_1, 2\colon X_2\rangle$$

  where the two-electron quantum state $|1\colon X_1, 2\colon X_2\rangle$ corresponds to "electron 1 in state $|X_1\rangle$ and electron 2 in state $|X_2\rangle$".

- Anti-symmetrization principle: a physical two-electron wavefunction should fulfill the condition

$$\boxed{\Psi(X_1, X_2) = -\Psi(X_2, X_1)}$$

thus leading to

$$|\Psi\rangle = -\int\int \mathrm{d}X_1 \mathrm{d}X_2 \, \Psi(X_2, X_1) |1\text{:}\, \boxed{X_1}, 2\text{:}\, X_2\rangle = -\underbrace{\int\int \mathrm{d}X_1 \mathrm{d}X_2 \, \Psi(X_1, X_2) |1\text{:}\, X_2, 2\text{:}\, \boxed{X_1}\rangle}_{|\Psi_{1\leftrightarrow 2}\rangle}$$

and

$$|\Psi\rangle = \frac{1}{2}\int\int \mathrm{d}X_1 \mathrm{d}X_2 \underbrace{\Big[\Psi(X_1, X_2) - \Psi(X_2, X_1)\Big]}_{0 \quad \text{if} \quad X_1 = X_2} |1\text{:}\, X_1, 2\text{:}\, X_2\rangle$$

Conclusion: the anti-symmetrization of the wavefunction ensures that electrons are indistinguishable and that they cannot be in the same quantum state (Pauli principle).

Slater determinants

- Let $\left\{\varphi_K(X)\right\}_K$ denote an orthonormal basis of (molecular) spin-orbitals. Two electrons that occupy the spin-orbitals $\varphi_I(X)$ and $\varphi_J(X)$ will be described by the (normalized) Slater determinant

$$\Phi_{IJ}(X_1, X_2) = \frac{1}{\sqrt{2}} \begin{vmatrix} \varphi_I(X_1) & \varphi_I(X_2) \\ \varphi_J(X_1) & \varphi_J(X_2) \end{vmatrix} = \frac{1}{\sqrt{2}} \left(\varphi_I(X_1)\varphi_J(X_2) - \varphi_I(X_2)\varphi_J(X_1) \right)$$

- Note that Slater determinants and, consequently, linear combinations of Slater determinants are anti-symmetric.

- Therefore, Slater determinants are convenient "building blocks" for computing the electronic wavefunction.

- Still, we may wonder if we really need this complicated expression obtained from the determinant (obviously things get worse for a larger number of electrons).

- Another drawback of the current formulation: both Slater determinant and Hamiltonian expressions depend on the number of electrons.

"What is occupied ?" rather than "Who occupies what ?"

- Since electrons are indistinguishable, there is no need to know that electron 1 occupies φ_I and electron 2 occupies φ_J or the other way around ...

- The important information is that spin-orbitals φ_I and φ_J are occupied and the remaining ones are empty.

- Second quantization is a formalism that relies on this idea. Let us tell the story ...

- At the beginning, there was "nothing" ... $|\text{vac}\rangle$ $\longleftarrow$ normalized "vacuum state"

- ... then was introduced the concept of annihilation of an electron occupying φ_I , that would obviously give zero when applied to the vacuum state: $\boxed{\forall\, I,\quad \hat{a}_I|\text{vac}\rangle = 0}$ **(rule 1)**

- ... and then came the concept of creation of an electron occupying φ_I: $\hat{a}_I^\dagger|\text{vac}\rangle \equiv |\varphi_I\rangle$

"What is occupied ?" rather than "Who occupies what ?"

- ... and then came the idea to create another electron occupying φ_J: $\quad \hat{a}_J^\dagger \hat{a}_I^\dagger |\text{vac}\rangle \equiv |\Phi_{IJ}\rangle$

- Note that the creation operator $\hat{a}_I^\dagger$ is the adjoint of the annihilation operator $\hat{a}_I$. This ensures, in particular, that one-electron and vacuum states are orthogonal:

$$\langle \varphi_I | \text{vac} \rangle = \langle \hat{a}_I^\dagger \text{vac} | \text{vac} \rangle = \langle \text{vac} | \hat{a}_I | \text{vac} \rangle = 0$$

- In order to have a representation that is equivalent to the one used in first quantization, we only need two more rules:

$$\boxed{\forall\, I, J, \quad [\hat{a}_I, \hat{a}_J]_+ = \hat{a}_I \hat{a}_J + \hat{a}_J \hat{a}_I = 0} \qquad \textbf{(rule 2)} \qquad \longrightarrow \qquad \left[\hat{a}_I^\dagger, \hat{a}_J^\dagger\right]_+ = [\hat{a}_J, \hat{a}_I]_+^\dagger = 0$$

$$\boxed{\forall\, I, J, \quad \left[\hat{a}_I, \hat{a}_J^\dagger\right]_+ = \hat{a}_I \hat{a}_J^\dagger + \hat{a}_J^\dagger \hat{a}_I = \delta_{IJ}} \qquad \textbf{(rule 3)}$$

- **Notation:** the anticommutator of $\hat{A}$ with $\hat{B}$ reads $\left[\hat{A}, \hat{B}\right]_+ = \left\{\hat{A}, \hat{B}\right\} = \hat{A}\hat{B} + \hat{B}\hat{A}$ while the commutator equals $\left[\hat{A}, \hat{B}\right] = \hat{A}\hat{B} - \hat{B}\hat{A}$.

"What is occupied ?" rather than "Who occupies what ?"

- **Rule 2** contains the indistinguishability of the electrons, $\hat{a}_J^\dagger \hat{a}_I^\dagger |\text{vac}\rangle = -\hat{a}_I^\dagger \hat{a}_J^\dagger |\text{vac}\rangle$,

  and the Pauli principle, $\hat{a}_I^\dagger \hat{a}_I^\dagger |\text{vac}\rangle = 0$.

- **Rule 3** ensures that you can only annihilate what has already been created (!),

  $\hat{a}_I \hat{a}_J^\dagger |\text{vac}\rangle = \delta_{IJ} |\text{vac}\rangle - \hat{a}_J^\dagger \hat{a}_I |\text{vac}\rangle = \delta_{IJ} |\text{vac}\rangle$

- It is now very easy to generate representations of Slater determinants for an arbitrary number N of electrons: multiply more creation operators !

$$|I_1 I_2 \ldots I_{N-1} I_N\rangle = \hat{a}_{I_1}^\dagger \hat{a}_{I_2}^\dagger \ldots \hat{a}_{I_{N-1}}^\dagger \hat{a}_{I_N}^\dagger |\text{vac}\rangle \quad \equiv \quad \frac{1}{\sqrt{N!}} \det\left[\varphi_{I_i}(X_j)\right]$$

EXERCISE: **(1)** Show that $|I_1 I_2 \ldots I_{N-1} I_N\rangle$ is normalized.

(2) Let us consider another state $|J_1 J_2 \ldots J_{N-1} J_N\rangle$ and assume that at least one of the occupied spin-orbitals (let us denote it φ_{J_k}) is not occupied in $|I_1 I_2 \ldots I_{N-1} I_N\rangle$. Show that the two states are orthogonal.

(3) The "counting" operator $\hat{N}$ is defined as $\hat{N} = \sum_I \hat{n}_I$ where $\hat{n}_I = \hat{a}_I^\dagger \hat{a}_I$. Show that

$$\hat{n}_I |I_1 I_2 \ldots I_{N-1} I_N\rangle = |I_1 I_2 \ldots I_{N-1} I_N\rangle \qquad \text{if} \quad I = I_k \qquad 1 \leq k \leq N$$

$$= 0 \qquad\qquad\qquad\qquad \text{otherwise}$$

and conclude that $\boxed{\hat{N}|I_1 I_2 \ldots I_{N-1} I_N\rangle = N|I_1 I_2 \ldots I_{N-1} I_N\rangle}$.

(4) Explain why states corresponding to different numbers of electrons are automatically orthogonal.

(5) Explain why any normalized state $|\Psi\rangle$ fulfills the condition $\boxed{0 \leq \langle\Psi|\hat{n}_I|\Psi\rangle \leq 1}$.

One-electron operators in second quantization

- Let $\hat{h}$ denote a one-electron operator ($\hat{t} + \hat{v}_{\rm ne}$ for example): it acts on the one-electron states $|\varphi_I\rangle$.

- Resolution of the identity: $\sum_I |\varphi_I\rangle\langle\varphi_I| = \hat{\mathbb{1}}$,

  which leads to the conventional representation $\hat{h} = \hat{\mathbb{1}}\ \hat{h}\ \hat{\mathbb{1}} = \sum_{I,J} \langle\varphi_I|\hat{h}|\varphi_J\rangle |\varphi_I\rangle\langle\varphi_J|$.

- Second-quantized representation: $\boxed{\hat{h} \equiv \sum_{I,J} \langle\varphi_I|\hat{h}|\varphi_J\rangle \hat{a}_I^\dagger \hat{a}_J}$

  Indeed,

$$\left(\sum_{I,J} \langle\varphi_I|\hat{h}|\varphi_J\rangle \hat{a}_I^\dagger \hat{a}_J\right) |\varphi_K\rangle = \left(\sum_{I,J} \langle\varphi_I|\hat{h}|\varphi_J\rangle \hat{a}_I^\dagger \hat{a}_J\right) \hat{a}_K^\dagger |\mathrm{vac}\rangle = \left(\sum_{I,J} \langle\varphi_I|\hat{h}|\varphi_J\rangle \hat{a}_I^\dagger\, \delta_{JK}\right) |\mathrm{vac}\rangle$$

$$= \sum_I \langle\varphi_I|\hat{h}|\varphi_K\rangle |\varphi_I\rangle = \hat{h}|\varphi_K\rangle$$

- What is convenient is that this second-quantized representation is valid for any number N of electrons:

$$\boxed{\sum_{i=1}^{N} \hat{h}(i) \quad \equiv \quad \sum_{I,J} \langle \varphi_I | \hat{h} | \varphi_J \rangle \hat{a}_I^\dagger \hat{a}_J \quad \equiv \quad \hat{h}}$$

The information about N has been completely transferred to the states. It does not appear in the operator anymore.

EXERCISE: Let us consider another orthonormal basis $\left\{ \tilde{\varphi}_K(X) \right\}_K$ of spin-orbitals that we decompose in the current basis as follows, $|\tilde{\varphi}_P\rangle = \sum_Q U_{QP} |\varphi_Q\rangle$.

(1) Show that the matrix $\mathbf{U}$ is unitary ($\mathbf{U}^\dagger = \mathbf{U}^{-1}$).

(2) Explain why $\hat{a}_{\tilde{P}}^\dagger = \sum_Q U_{QP}\, \hat{a}_Q^\dagger$ and show that $\sum_{I,J} \langle \tilde{\varphi}_I | \hat{h} | \tilde{\varphi}_J \rangle \hat{a}_{\tilde{I}}^\dagger \hat{a}_{\tilde{J}} \equiv \hat{h}$.

(3) Show how the diagonalization of $\hat{h}$ in the one-electron space leads automatically to the diagonalization in the N-electron space (use exercise page 10).

Two-electron operators in second quantization

- Let $\hat{w}$ denote a two-electron operator: it acts on two-electron states $|\varphi_I \varphi_J\rangle = |1\text{: } \varphi_I, 2\text{: } \varphi_J\rangle$.

- A complete anti-symmetrized basis should be used for describing the two electrons:

$$|IJ\rangle = \frac{1}{\sqrt{2}}\Big(|\varphi_I \varphi_J\rangle - |\varphi_J \varphi_I\rangle\Big) \equiv \hat{a}_I^\dagger \hat{a}_J^\dagger |\text{vac}\rangle \qquad \text{with} \quad I < J.$$

  Consequently, any two-electron anti-symmetrized state $|\Psi\rangle$ shoud fulfill the condition

$$\hat{P}_{\mathcal{A}}|\Psi\rangle = |\Psi\rangle \qquad \text{where} \quad \hat{P}_{\mathcal{A}} = \sum_{I<J} |IJ\rangle\langle IJ| \qquad \longleftarrow \quad \text{projection operator !}$$

- Projection of the two-electron operator onto the space of anti-symmetrized states:

$$\hat{w}_{\mathcal{A}} = \hat{P}_{\mathcal{A}} \hat{w} \hat{P}_{\mathcal{A}} = \sum_{I<J,K<L} \langle IJ|\hat{w}|KL\rangle \; |IJ\rangle\langle KL|$$

Two-electron operators in second quantization

EXERCISE: Prove that $\hat{w}_{\mathcal{A}} \equiv \frac{1}{2} \sum_{IJKL} \langle \varphi_I \varphi_J | \hat{w} | \varphi_K \varphi_L \rangle \, \hat{a}_I^\dagger \hat{a}_J^\dagger \hat{a}_L \hat{a}_K$

hint: apply $\hat{w}_{\mathcal{A}}$ and the proposed second-quantized representation to $|PQ\rangle \equiv \hat{a}_P^\dagger \hat{a}_Q^\dagger |\text{vac}\rangle \quad (P < Q)$. Conclude.

- What is convenient is that this second-quantized representation is valid for any number N of electrons and includes the projection onto anti-symmetrized states:

$$\frac{1}{2} \sum_{i \neq j}^{N} \hat{w}(i,j) \equiv \frac{1}{2} \sum_{IJKL} \langle \varphi_I \varphi_J | \hat{w} | \varphi_K \varphi_L \rangle \, \hat{a}_I^\dagger \hat{a}_J^\dagger \hat{a}_L \hat{a}_K \equiv \hat{w}$$

Summary

- In summary, the electronic Hamiltonian can be written in second quantization as follows,

$$\hat{H} = \sum_{IJ} \langle \varphi_I | \hat{h} | \varphi_J \rangle \hat{a}_I^\dagger \hat{a}_J + \frac{1}{2} \sum_{IJKL} \langle \varphi_I \varphi_J | \hat{w}_{\text{ee}} | \varphi_K \varphi_L \rangle \, \hat{a}_I^\dagger \hat{a}_J^\dagger \hat{a}_L \hat{a}_K$$

where $\langle \varphi_I | \hat{h} | \varphi_J \rangle = \int \mathrm{d}X \, \varphi_I^*(X) \Big(\hat{h} \varphi_J \Big)(X)$ ← one-electron integrals

$$\langle \varphi_I \varphi_J | \hat{w}_{\text{ee}} | \varphi_K \varphi_L \rangle = \int \int \mathrm{d}X_1 \mathrm{d}X_2 \, \varphi_I^*(X_1) \varphi_J^*(X_2) \Big(\hat{w}_{\text{ee}} \varphi_K \varphi_L \Big)(X_1, X_2)$$ ←two-electron integrals

- Note that this expression is also valid for a relativistic Hamiltonian. Two or four-component spinors should be used rather than spin-orbitals in conjunction with the Dirac (Breit) Coulomb Hamiltonian.

- The standard (non-relativistic) Hamiltonian will be used in the following.

Why "second" quantization ?

- Let us focus on the (one-electron) electron-nuclei local potential operator which, in second quantization, reads $\hat{V}_{\rm ne} = \sum_{i=1}^{N} \hat{v}_{\rm ne}(i) \equiv \sum_{IJ} \langle \varphi_I | \hat{v}_{\rm ne} | \varphi_J \rangle \hat{a}_I^\dagger \hat{a}_J$ where

$$\langle \varphi_I | \hat{v}_{\rm ne} | \varphi_J \rangle = \int {\rm d}X \; v_{\rm ne}(\mathbf{r}) \varphi_I^*(X) \varphi_J(X),$$

thus leading to

$$\hat{V}_{\rm ne} \equiv \int {\rm d}X \; v_{\rm ne}(\mathbf{r}) \underbrace{\left(\sum_I \varphi_I^*(X) \hat{a}_I^\dagger \right)}_{\hat{\Psi}^\dagger(X)} \underbrace{\left(\sum_J \varphi_J(X) \hat{a}_J \right)}_{\hat{\Psi}(X)} = \boxed{\int {\rm d}\mathbf{r} \, v_{\rm ne}(\mathbf{r}) \sum_\sigma \hat{\Psi}^\dagger(\mathbf{r},\sigma) \hat{\Psi}(\mathbf{r},\sigma) \equiv \hat{V}_{\rm ne}}$$

$\hat{\Psi}^\dagger(X)$ $\hat{\Psi}(X)$ $\leftarrow$ field operators

- For a single electron occupying the spin-orbital $\Psi(X) = \Psi(\mathbf{r},\sigma)$, the corresponding expectation value for the electron-nuclei potential energy equals

$$\langle \Psi | \hat{v}_{\rm ne} | \Psi \rangle = \int {\rm d}X \; v_{\rm ne}(\mathbf{r}) \Psi^*(X) \Psi(X) = \int {\rm d}\mathbf{r} \, v_{\rm ne}(\mathbf{r}) \sum_\sigma \Psi^*(\mathbf{r},\sigma) \Psi(\mathbf{r},\sigma).$$

- **Physical interpretation** of the field operators:

$$\hat{\Psi}^\dagger(X)|\text{vac}\rangle = \sum_I \varphi_I^*(X)\hat{a}_I^\dagger|\text{vac}\rangle = \sum_I \varphi_I^*(X)|\varphi_I\rangle = \sum_I |\varphi_I\rangle\langle\varphi_I|X\rangle = \boxed{|X\rangle = \hat{\Psi}^\dagger(X)|\text{vac}\rangle},$$

  which means that $\hat{\Psi}^\dagger(X) = \hat{\Psi}^\dagger(\mathbf{r},\sigma)$ creates an electron at position $\mathbf{r}$ with spin σ.

- Consequently, the density operator reads in second quantization $\boxed{\hat{n}(\mathbf{r}) = \sum_\sigma \hat{\Psi}^\dagger(\mathbf{r},\sigma)\hat{\Psi}(\mathbf{r},\sigma)}$,

  and the electron density associated with the normalized N-electron wavefunction Ψ is simply calculated as follows,

$$n_\Psi(\mathbf{r}) = \langle\Psi|\hat{n}(\mathbf{r})|\Psi\rangle.$$

- **Anticommutation rules:** $\left[\hat{\Psi}(X), \hat{\Psi}(X')\right]_+ = \sum_{IJ} \varphi_J(X)\varphi_I(X')\left[\hat{a}_J, \hat{a}_I\right]_+ = 0$ and

$$\begin{aligned}\left[\hat{\Psi}(X), \hat{\Psi}^\dagger(X')\right]_+ &= \sum_{IJ} \varphi_J(X)\varphi_I^*(X')\left[\hat{a}_J, \hat{a}_I^\dagger\right]_+ = \sum_{IJ} \varphi_J(X)\varphi_I^*(X')\delta_{IJ} \\ &= \sum_I \langle X|\varphi_I\rangle\langle\varphi_I|X'\rangle = \langle X|X'\rangle = \delta(X - X').\end{aligned}$$

EXERCISE:

(1) Show that each contribution to the electronic Hamiltonian reads in second quantization as follows,

$$\hat{T} \equiv -\frac{1}{2}\int \mathrm{d}\mathbf{r} \sum_{\sigma} \hat{\Psi}^{\dagger}(\mathbf{r},\sigma)\nabla^2_{\mathbf{r}}\hat{\Psi}(\mathbf{r},\sigma), \qquad \hat{V}_{\mathrm{ne}} \equiv \int \mathrm{d}\mathbf{r}\, v_{\mathrm{ne}}(\mathbf{r})\, \hat{n}(\mathbf{r}),$$

$$\hat{W}_{\mathrm{ee}} \equiv \frac{1}{2}\int\int \mathrm{d}\mathbf{r}\mathrm{d}\mathbf{r}' \sum_{\sigma\sigma'} \frac{\hat{\Psi}^{\dagger}(\mathbf{r},\sigma)\hat{\Psi}^{\dagger}(\mathbf{r}',\sigma')\hat{\Psi}(\mathbf{r}',\sigma')\hat{\Psi}(\mathbf{r},\sigma)}{|\mathbf{r}-\mathbf{r}'|}$$

$$= \frac{1}{2}\int\int \mathrm{d}\mathbf{r}\mathrm{d}\mathbf{r}' \frac{\hat{n}(\mathbf{r})\hat{n}(\mathbf{r}') - \delta(\mathbf{r}-\mathbf{r}')\hat{n}(\mathbf{r})}{|\mathbf{r}-\mathbf{r}'|}.$$

(2) At the non-relativistic level, real algebra can be used, $\varphi_I(X) = \varphi_{i\sigma}(\mathbf{r},\tau) = \phi_i(\mathbf{r})\delta_{\sigma\tau}$,

$$\hat{h} \equiv -\frac{1}{2}\nabla^2_{\mathbf{r}} + v_{\mathrm{ne}}(\mathbf{r})\times \quad \text{and} \quad \hat{w}_{\mathrm{ee}} \equiv \frac{1}{|\mathbf{r}_1-\mathbf{r}_2|}\times \ .$$

Show that the Hamiltonian, that is here a spin-free operator, can be rewritten in the basis of the molecular orbitals $\left\{\phi_p(\mathbf{r})\right\}_p$ as follows

$$\hat{H} = \sum_{p,q} h_{pq}\hat{E}_{pq} + \frac{1}{2}\sum_{p,q,r,s}\langle pr|qs\rangle\left(\hat{E}_{pq}\hat{E}_{rs} - \delta_{qr}\hat{E}_{ps}\right)$$

where $\hat{E}_{pq} = \sum_{\sigma} \hat{a}^{\dagger}_{p,\sigma} \hat{a}_{q,\sigma}$, $h_{pq} = \langle \phi_p | \hat{h} | \phi_q \rangle$ and

$$\langle pr|qs\rangle = \int\int \mathrm{d}\mathbf{r}_1 \mathrm{d}\mathbf{r}_2 \; \phi_p(\mathbf{r}_1)\phi_r(\mathbf{r}_2) \frac{1}{|\mathbf{r}_1 - \mathbf{r}_2|} \phi_q(\mathbf{r}_1)\phi_s(\mathbf{r}_2) = (pq|rs).$$

Model Hamiltonians: example of the Hubbard Hamiltonian

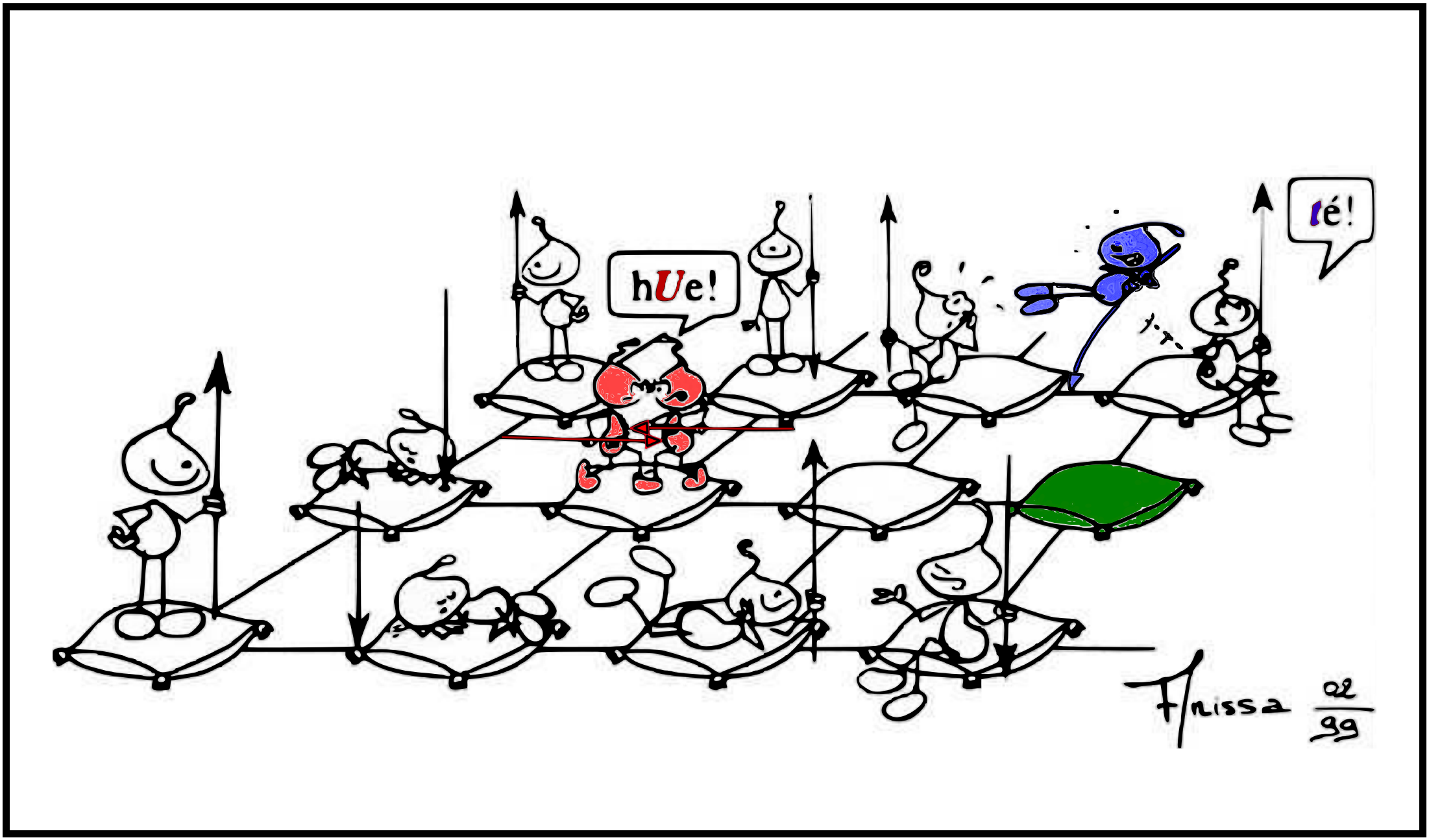
$$h_{ij} \longrightarrow -t\left(\delta_{i,j-1} + \delta_{i,j+1}\right) + \varepsilon_i \delta_{ij}$$

$$\langle ij|kl\rangle \longrightarrow U \delta_{ij}\delta_{ik}\delta_{lj}$$

$$\hat{E}_{ik}\hat{E}_{jl} - \delta_{kj}\hat{E}_{il} \longrightarrow \hat{n}_i\hat{n}_i - \hat{n}_i$$

where $\hat{n}_i = \hat{E}_{ii} = \hat{n}_{i\uparrow} + \hat{n}_{i\downarrow}$ so that $\hat{n}_i\hat{n}_i = 2\hat{n}_{i\uparrow}\hat{n}_{i\downarrow} + \hat{n}_i$

$$\hat{H} \longrightarrow \underbrace{-t \sum_{\langle i,j\rangle} \sum_{\sigma=\uparrow,\downarrow} \hat{a}^\dagger_{i,\sigma}\hat{a}_{j,\sigma}}_{\hat{\mathcal{T}}\text{ (hopping)}} + \underbrace{U \sum_i \hat{n}_{i\uparrow}\hat{n}_{i\downarrow}}_{\text{on-site repulsion}} + \underbrace{\sum_i \varepsilon_i \hat{n}_i}_{\text{local potential}}$$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