

Density-functional theory without density functionals: A quantum embedding perspective

***Sajanthan Sekaran^a, Saad Yalouz^a, Martin Rafael Gulin^a, Filip Cernatic^a,
Masahisa Tsuchiizu^b, Matthieu Saubanère^c, and Emmanuel Fromager^a***

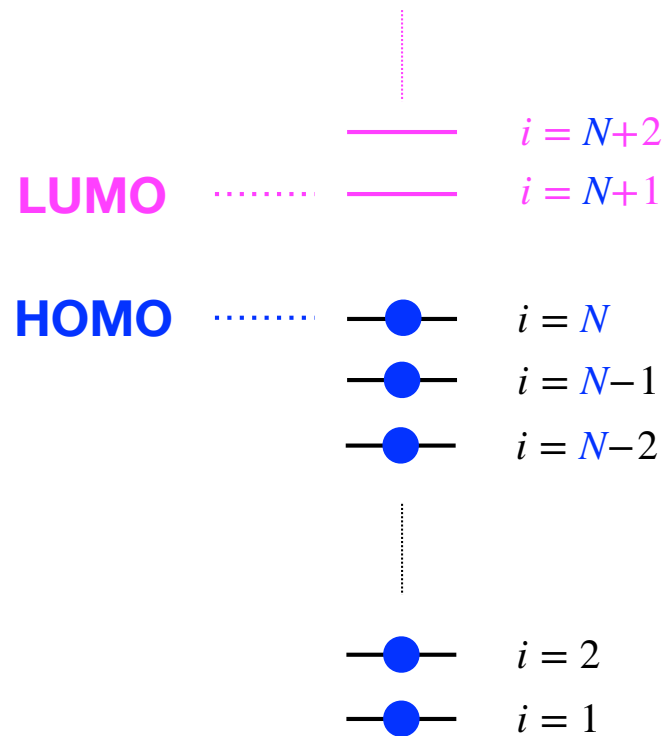
^a*Laboratoire de Chimie Quantique, Institut de Chimie de Strasbourg,
Université de Strasbourg, Strasbourg, France*

^b*Department of Physics, Nara Women's University, Nara, Japan*

^c*Institut Charles Gerhardt, CNRS/Université de Montpellier, France*

(Kohn-Sham) DFT for N -electron ground states

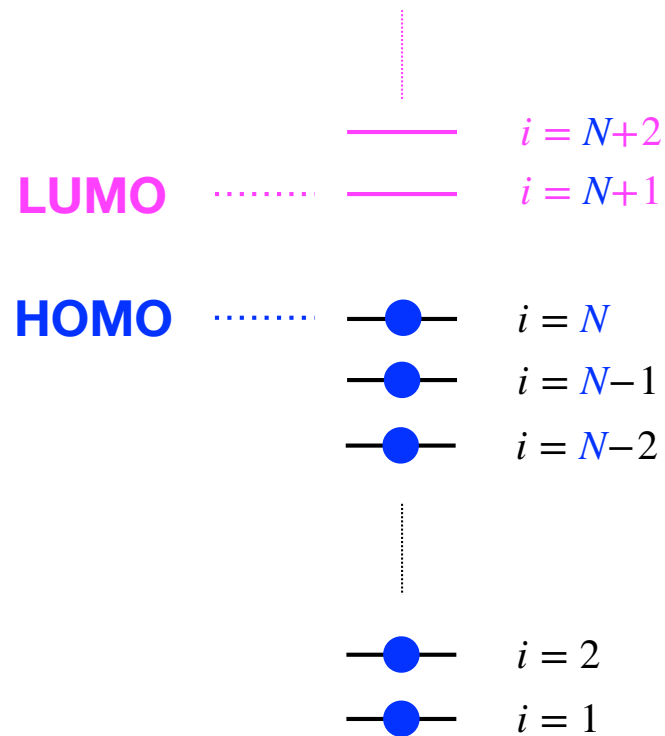
$$\left[-\frac{\nabla_{\mathbf{r}}^2}{2} + v_{\text{ext}}(\mathbf{r}) + v_{\text{Hxc}}(\mathbf{r}) \right] \varphi_i(\mathbf{r}) = \varepsilon_i \varphi_i(\mathbf{r})$$



(Kohn-Sham) DFT for N -electron ground states

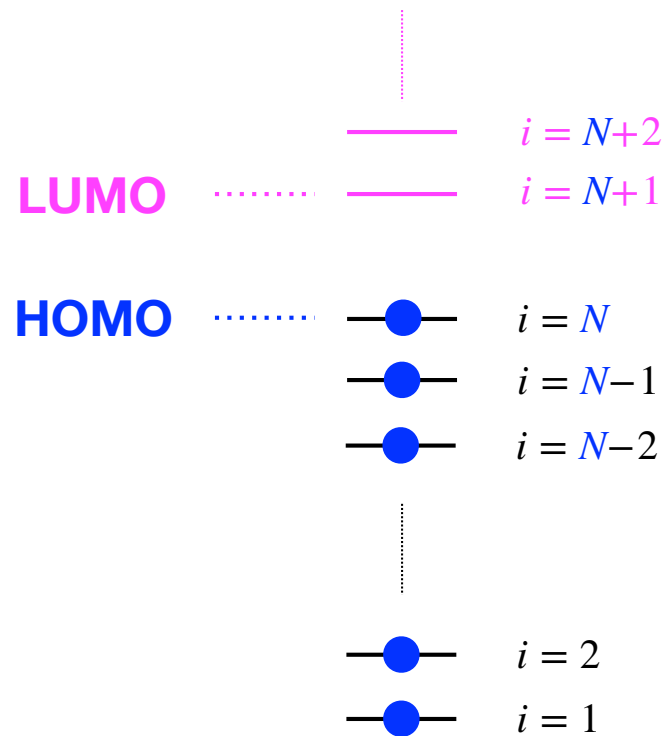
One-electron picture

$$\left[-\frac{\nabla_{\mathbf{r}}^2}{2} + v_{\text{ext}}(\mathbf{r}) + v_{\text{Hxc}}(\mathbf{r}) \right] \varphi_i(\mathbf{r}) = \varepsilon_i \varphi_i(\mathbf{r})$$



DFT for N -electron ground states

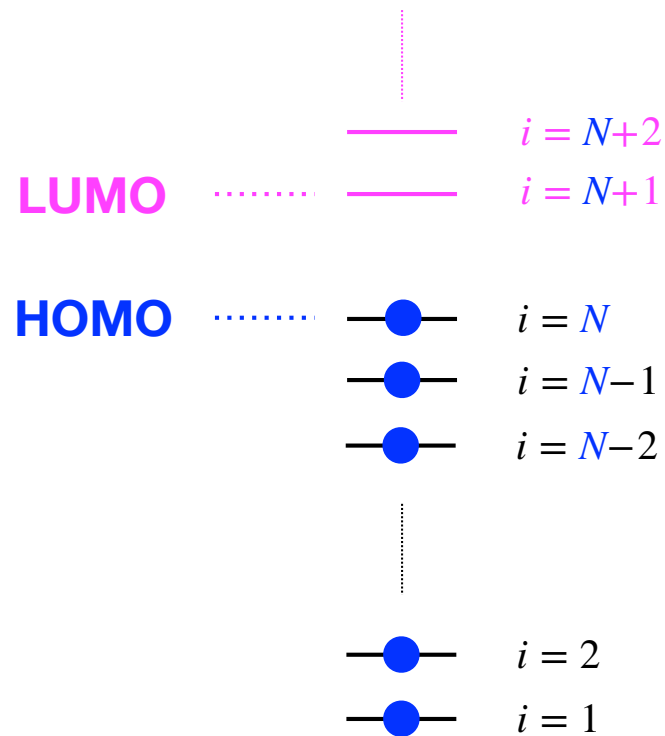
$$\mathbf{r} \equiv (x, y, z) \quad \left[-\frac{\nabla_{\mathbf{r}}^2}{2} + v_{\text{ext}}(\mathbf{r}) + v_{\text{Hxc}}(\mathbf{r}) \right] \varphi_i(\mathbf{r}) = \varepsilon_i \varphi_i(\mathbf{r})$$



$$\nabla_{\mathbf{r}}^2 \equiv \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2}$$

DFT for N -electron ground states

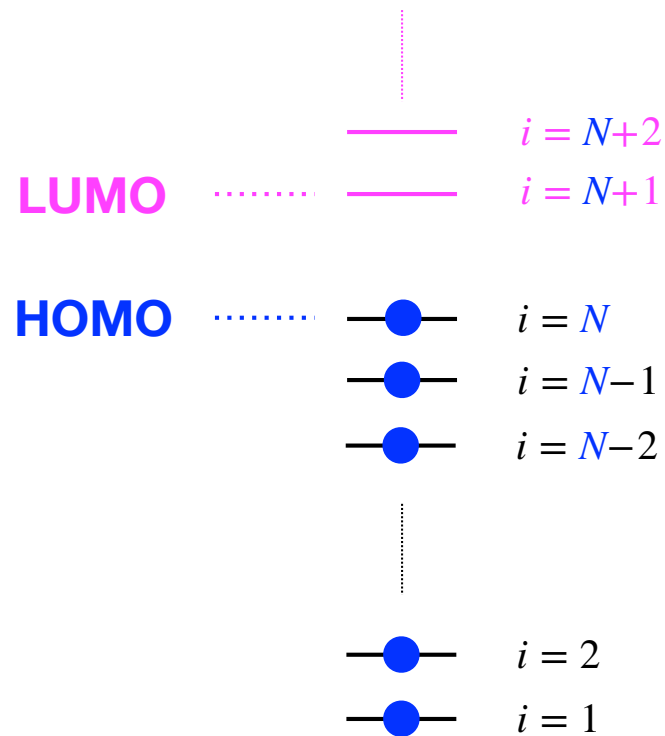
$$\left[-\frac{\nabla_{\mathbf{r}}^2}{2} + v_{\text{ext}}(\mathbf{r}) + v_{\text{Hxc}}(\mathbf{r}) \right] \varphi_i(\mathbf{r}) = \varepsilon_i \varphi_i(\mathbf{r})$$



DFT for N -electron ground states

“External” local and **multiplicative** potential

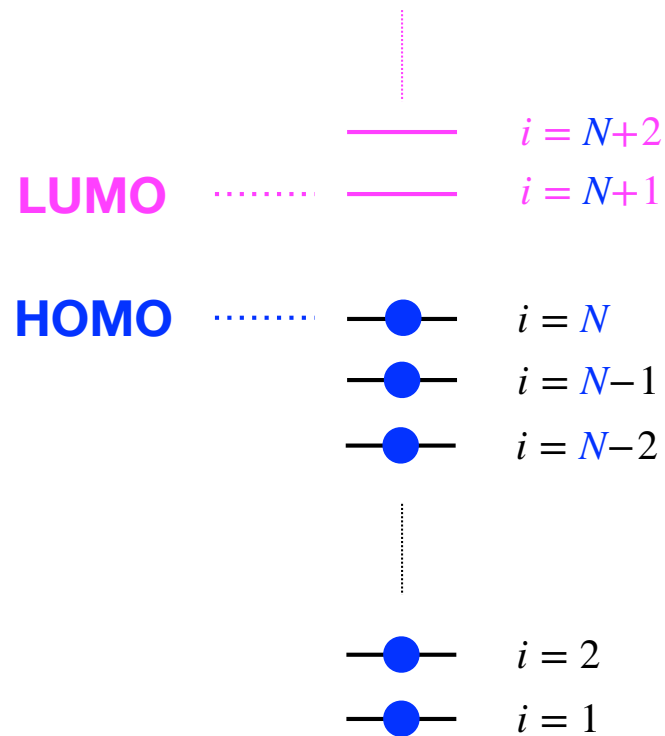
$$\left[-\frac{\nabla_{\mathbf{r}}^2}{2} + \overset{\downarrow}{v_{\text{ext}}(\mathbf{r})} + v_{\text{Hxc}}(\mathbf{r}) \right] \varphi_i(\mathbf{r}) = \varepsilon_i \varphi_i(\mathbf{r})$$



DFT for N -electron ground states

Nuclear potential in quantum chemistry calculations

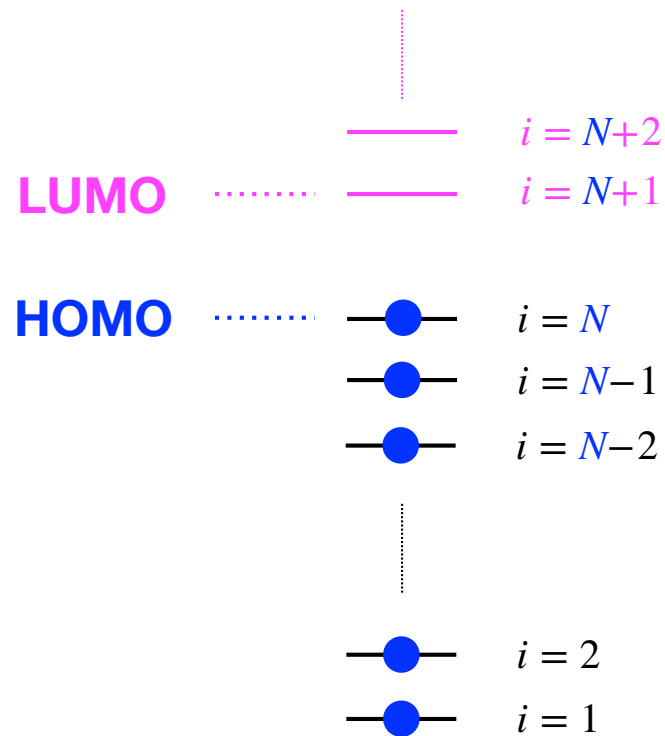
$$\left[-\frac{\nabla_{\mathbf{r}}^2}{2} + \overset{\downarrow}{v_{\text{ext}}(\mathbf{r})} + v_{\text{Hxc}}(\mathbf{r}) \right] \varphi_i(\mathbf{r}) = \varepsilon_i \varphi_i(\mathbf{r})$$



DFT for N -electron ground states

**Hartree-exchange-correlation
local (multiplicative) potential**

$$\left[-\frac{\nabla_{\mathbf{r}}^2}{2} + v_{\text{ext}}(\mathbf{r}) + v_{\text{Hxc}}(\mathbf{r}) \right] \varphi_i(\mathbf{r}) = \varepsilon_i \varphi_i(\mathbf{r})$$



DFT for N -electron ground states

$$\left[-\frac{\nabla_{\mathbf{r}}^2}{2} + v_{\text{ext}}(\mathbf{r}) + v_{\text{Hxc}}(\mathbf{r}) \right] \varphi_i(\mathbf{r}) = \varepsilon_i \varphi_i(\mathbf{r})$$

$$n_0^N(\mathbf{r}) = \sum_{i=1}^N |\varphi_i(\mathbf{r})|^2$$

*exact ground-state
density*

DFT for N -electron ground states

$$\left[-\frac{\nabla_{\mathbf{r}}^2}{2} + v_{\text{ext}}(\mathbf{r}) + v_{\text{Hxc}}(\mathbf{r}) \right] \varphi_i(\mathbf{r}) = \varepsilon_i \varphi_i(\mathbf{r})$$

Hartree-exchange-correlation potential

$$v_{\text{Hxc}}(\mathbf{r}) = \left. \frac{\delta E_{\text{Hxc}}[n]}{\delta n(\mathbf{r})} \right|_{n=n_0^N}$$

DFT for N -electron ground states

$$\left[-\frac{\nabla_{\mathbf{r}}^2}{2} + v_{\text{ext}}(\mathbf{r}) + v_{\text{Hxc}}(\mathbf{r}) \right] \varphi_i(\mathbf{r}) = \varepsilon_i \varphi_i(\mathbf{r})$$

Hartree-exchange-correlation potential

Hxc energy functional

$$v_{\text{Hxc}}(\mathbf{r}) = \frac{\delta E_{\text{Hxc}}[n]}{\delta n(\mathbf{r})} \Big|_{n=n_0^N}$$

DFT for N -electron ground states

$$\left[-\frac{\nabla_{\mathbf{r}}^2}{2} + v_{\text{ext}}(\mathbf{r}) + v_{\text{Hxc}}(\mathbf{r}) \right] \varphi_i(\mathbf{r}) = \varepsilon_i \varphi_i(\mathbf{r})$$

$$v_{\text{Hxc}}(\mathbf{r}) = \left. \frac{\delta E_{\text{Hxc}}[n]}{\delta n(\mathbf{r})} \right|_{n=n_0^N}$$

Self-consistency!

$$n_0^N(\mathbf{r}) = \sum_{i=1}^N |\varphi_i(\mathbf{r})|^2$$

DFT for N -electron ground states

$$\left[-\frac{\nabla_{\mathbf{r}}^2}{2} + v_{\text{ext}}(\mathbf{r}) + v_{\text{Hxc}}(\mathbf{r}) \right] \varphi_i(\mathbf{r}) = \varepsilon_i \varphi_i(\mathbf{r})$$

$$n_0^N(\mathbf{r}) = \sum_{i=1}^N |\varphi_i(\mathbf{r})|^2$$

*exact ground-state
density*

$$E_0 = -\frac{1}{2} \sum_{i=1}^N \int d\mathbf{r} \varphi_i(\mathbf{r}) \nabla_{\mathbf{r}}^2 \varphi_i(\mathbf{r}) + \int d\mathbf{r} v_{\text{ext}}(\mathbf{r}) n_0^N(\mathbf{r}) + E_{\text{Hxc}}[n] \Big|_{n=n_0^N}$$

*exact ground-state
energy*

DFT for N -electron ground states

$$\left[-\frac{\nabla_{\mathbf{r}}^2}{2} + v_{\text{ext}}(\mathbf{r}) + v_{\text{Hxc}}(\mathbf{r}) \right] \varphi_i(\mathbf{r}) = \varepsilon_i \varphi_i(\mathbf{r})$$

exact Hxc potential

$$v_{\text{Hxc}}(\mathbf{r}) = \left. \frac{\delta E_{\text{Hxc}}[n]}{\delta n(\mathbf{r})} \right|_{n=n_0^N}$$

$$n_0^N(\mathbf{r}) = \sum_{i=1}^N |\varphi_i(\mathbf{r})|^2$$

exact ground-state density

$$E_0 = -\frac{1}{2} \sum_{i=1}^N \int d\mathbf{r} \varphi_i(\mathbf{r}) \nabla_{\mathbf{r}}^2 \varphi_i(\mathbf{r}) + \int d\mathbf{r} v_{\text{ext}}(\mathbf{r}) n_0^N(\mathbf{r}) + E_{\text{Hxc}}[n] \Big|_{n=n_0^N}$$

exact ground-state energy

DFT for N -electron ground states

$$\left[-\frac{\nabla_{\mathbf{r}}^2}{2} + v_{\text{ext}}(\mathbf{r}) + v_{\text{Hxc}}(\mathbf{r}) \right] \varphi_i(\mathbf{r}) = \varepsilon_i \varphi_i(\mathbf{r})$$

exact Hxc potential

$$v_{\text{Hxc}}(\mathbf{r}) = \left. \frac{\delta E_{\text{Hxc}}[n]}{\delta n(\mathbf{r})} \right|_{n=n_0^N}$$

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exact ground-state density

$$E_0 = -\frac{1}{2} \sum_{i=1}^N \int d\mathbf{r} \varphi_i(\mathbf{r}) \nabla_{\mathbf{r}}^2 \varphi_i(\mathbf{r}) + \int d\mathbf{r} v_{\text{ext}}(\mathbf{r}) n_0^N(\mathbf{r}) + \left. E_{\text{Hxc}}[n] \right|_{n=n_0^N}$$

exact ground-state energy

DFT for N -electron ground states

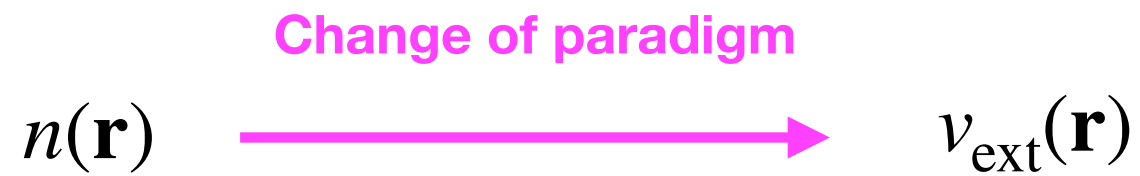
Once we have $E_{\text{Hxc}}[n]$ we can run a DFT calculation for *any system*

DFT for N -electron ground states

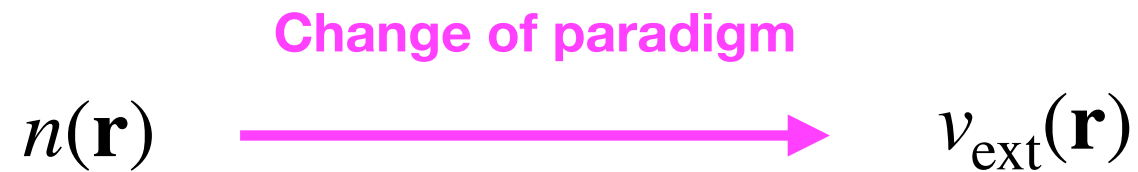
Once we have $E_{\text{Hxc}}[n]$ we can run a DFT calculation for *any system*

“Holy grail”

Local potential functional approach to DFT



Local potential functional approach to DFT



$$\left. \frac{\delta E_{\text{Hxc}}[n]}{\delta n(\mathbf{r})} \right|_{n=n_0^N} \equiv v_{\text{Hxc}}(\mathbf{r}) = v_{\text{Hxc}}[v_{\text{ext}}](\mathbf{r}) = ???$$

Local potential functional approach to DFT

Original N -electron Schrödinger equation

$$\hat{H}\Psi_0 = E_0\Psi_0$$

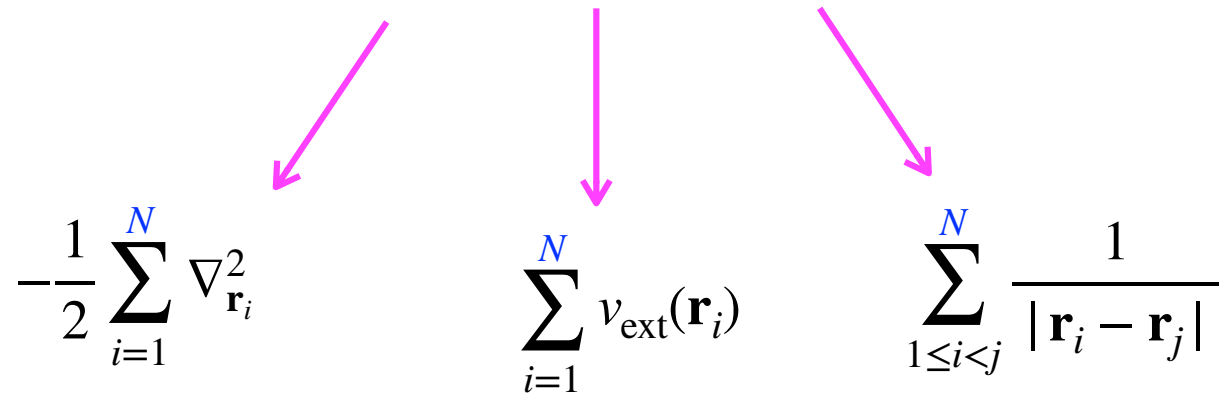
Local potential functional approach to DFT

Original N -electron Schrödinger equation

$$\hat{H}\Psi_0 = E_0\Psi_0$$

True physical Hamiltonian

$$\hat{H} = \hat{T} + \hat{V}_{\text{ext}} + \hat{W}_{\text{ee}}$$


$$-\frac{1}{2} \sum_{i=1}^N \nabla_{\mathbf{r}_i}^2$$

$$\sum_{i=1}^N v_{\text{ext}}(\mathbf{r}_i)$$

$$\sum_{1 \leq i < j}^N \frac{1}{|\mathbf{r}_i - \mathbf{r}_j|}$$

Local potential functional approach to DFT

Original N -electron Schrödinger equation

$$\hat{H}\Psi_0 = E_0\Psi_0$$

True physical Hamiltonian

$$\hat{H} = \hat{T} + \hat{V}_{\text{ext}} + \hat{W}_{\text{ee}}$$

$i = N+2$
 $i = N+1$
 $i = N$
 $i = N-1$
 $i = N-2$
 $i = 2$
 $i = 1$

$$\sum_{1 \leq i < j \leq N} \frac{1}{|\mathbf{r}_i - \mathbf{r}_j|}$$

Local potential functional approach to DFT

Original N -electron Schrödinger equation

$$\hat{H}\Psi_0 = E_0\Psi_0$$

True physical Hamiltonian

$$\hat{H} = \hat{T} + \hat{V}_{\text{ext}} + \hat{W}_{\text{ee}} \quad \textit{Known!}$$

Local potential functional approach to DFT

Original N -electron Schrödinger equation

$$\hat{H}\Psi_0 = E_0\Psi_0$$

True physical Hamiltonian

$$\hat{H} = \hat{T} + \hat{V}_{\text{ext}} + \hat{W}_{\text{ee}} \quad \textit{Known!}$$

*The Hxc potential is **completely unnecessary** here...*

Local potential functional approach to DFT

Original N -electron Schrödinger equation $\hat{H}\Psi_0 = E_0\Psi_0$

True physical Hamiltonian $\hat{H} = \hat{T} + \hat{V}_{\text{ext}} + \hat{W}_{\text{ee}}$ *Known!*

The Hxc potential is completely unnecessary here...

...unless Ψ_0 is constructed (approximately) from the Kohn-Sham orbitals.

Local potential functional approach to DFT

Original N -electron Schrödinger equation $\hat{H}\Psi_0 = E_0\Psi_0$

True physical Hamiltonian $\hat{H} = \hat{T} + \hat{V}_{\text{ext}} + \hat{W}_{\text{ee}}$ *Known!*

The Hxc potential is completely unnecessary here...

...unless Ψ_0 is constructed (approximately) from the Kohn-Sham orbitals.

The **approximate** one-electron **Kohn-Sham picture** and the **exact interacting picture** share the **same density** in this case

Local potential functional approach to DFT

$$\hat{H} = \hat{T} + \hat{V}_{\text{ext}} + \hat{V}_{\text{Hxc}} + \left(\hat{W}_{\text{ee}} - \hat{V}_{\text{Hxc}} \right)$$

Local potential functional approach to DFT

$$\hat{H} = \hat{T} + \hat{V}_{\text{ext}} + \hat{V}_{\text{Hxc}} + \left(\hat{W}_{\text{ee}} - \hat{V}_{\text{Hxc}} \right)$$



$$\sum_{i=1}^N v_{\text{Hxc}}(\mathbf{r}_i) \times$$

Local potential functional approach to DFT

$$\hat{H} = \hat{T} + \hat{V}_{\text{ext}} + \hat{V}_{\text{Hxc}} + \left(\hat{W}_{\text{ee}} - \hat{V}_{\text{Hxc}} \right)$$

Ground (and excited)
KS pictures

Local potential functional approach to DFT

$$\hat{H} = \hat{T} + \hat{V}_{\text{ext}} + \hat{V}_{\text{Hxc}} + \left(\hat{W}_{\text{ee}} - \hat{V}_{\text{Hxc}} \right)$$

Ground (and excited)
KS pictures

Slater determinants $\leftarrow \left\{ \Phi_I^{\text{KS}} \right\}_{I=0,1,2,\dots}$

Local potential functional approach to DFT

$$\hat{H} = \hat{T} + \hat{V}_{\text{ext}} + \hat{V}_{\text{Hxc}} + \underbrace{\left(\hat{W}_{\text{ee}} - \hat{V}_{\text{Hxc}} \right)}_{\text{“perturbation”}}$$

Local potential functional approach to DFT

$$\hat{H} = \hat{T} + \hat{V}_{\text{ext}} + \hat{V}_{\text{Hxc}} + \underbrace{\left(\hat{W}_{\text{ee}} - \hat{V}_{\text{Hxc}} \right)}_{\text{“perturbation”}}$$

$$E_0 = \langle \Phi_0^{\text{KS}} | \hat{H} | \Phi_0^{\text{KS}} \rangle + \sum_{I>0} \frac{\left| \langle \Phi_0^{\text{KS}} | \hat{W}_{\text{ee}} - \hat{V}_{\text{Hxc}} | \Phi_I^{\text{KS}} \rangle \right|^2}{\mathcal{E}_0^{\text{KS}} - \mathcal{E}_I^{\text{KS}}} + \dots$$

Local potential functional approach to DFT

$$\hat{H} = \hat{T} + \hat{V}_{\text{ext}} + \hat{V}_{\text{Hxc}} + \underbrace{\left(\hat{W}_{\text{ee}} - \hat{V}_{\text{Hxc}} \right)}$$

“perturbation”

$$E_0 = \langle \Phi_0^{\text{KS}} | \hat{H} | \Phi_0^{\text{KS}} \rangle + \sum_{I>0} \frac{\left| \langle \Phi_0^{\text{KS}} | \hat{W}_{\text{ee}} - \hat{V}_{\text{Hxc}} | \Phi_I^{\text{KS}} \rangle \right|^2}{\mathcal{E}_0^{\text{KS}} - \mathcal{E}_I^{\text{KS}}} + \dots$$

Görling-Levy perturbation theory

Local potential functional approach to DFT

$$\hat{H} = \hat{T} + \hat{V}_{\text{ext}} + \overset{???}{\hat{V}}_{\text{Hxc}} + \underbrace{\left(\overset{???}{\hat{W}}_{\text{ee}} - \overset{???}{\hat{V}}_{\text{Hxc}} \right)}_{\text{"perturbation"}}$$

$$E_0 = \overset{???}{\langle \Phi_0^{\text{KS}} |} \overset{???}{\hat{H}} \overset{???}{| \Phi_0^{\text{KS}} \rangle} + \sum_{I>0} \frac{\overset{???}{\left| \langle \Phi_0^{\text{KS}} |} \overset{???}{\hat{W}}_{\text{ee}} - \overset{???}{\hat{V}}_{\text{Hxc}} \overset{???}{| \Phi_I^{\text{KS}} \rangle} \right|^2}{\overset{???}{\mathcal{E}_0^{\text{KS}}} - \overset{???}{\mathcal{E}_I^{\text{KS}}}} + \dots$$

Local potential functional approach to DFT

$$\hat{H} = \hat{T} + \hat{V}_{\text{ext}} + \hat{V}_{\text{Hxc}} + \underbrace{\left(\hat{W}_{\text{ee}} - \hat{V}_{\text{Hxc}} \right)}_{\text{"perturbation"}}$$

$$E_0 = \langle \Phi_0^{\text{KS}} | \hat{H} | \Phi_0^{\text{KS}} \rangle + \sum_{I>0} \frac{\left| \langle \Phi_0^{\text{KS}} | \hat{W}_{\text{ee}} - \hat{V}_{\text{Hxc}} | \Phi_I^{\text{KS}} \rangle \right|^2}{\mathcal{E}_0^{\text{KS}} - \mathcal{E}_I^{\text{KS}}} + \dots$$

The evaluation of the energy fully relies on \hat{V}_{Hxc}

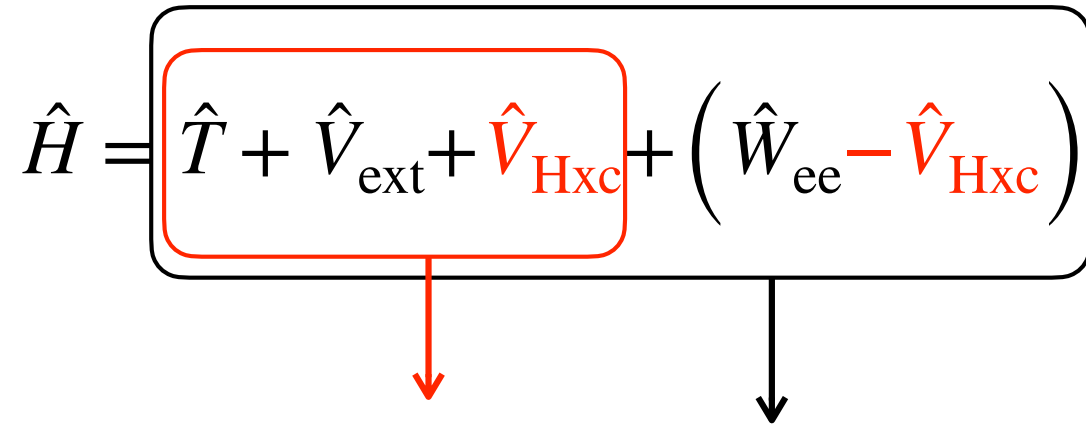
Local potential functional approach to DFT

$$\hat{H} = \hat{T} + \hat{V}_{\text{ext}} + \hat{V}_{\text{Hxc}}^{???) + \underbrace{\left(\hat{W}_{\text{ee}} - \hat{V}_{\text{Hxc}}^{???)}_{\text{“perturbation”}}$$

To-be-determined local potential
(hence the name of the approach)

The evaluation of the energy fully relies on \hat{V}_{Hxc}

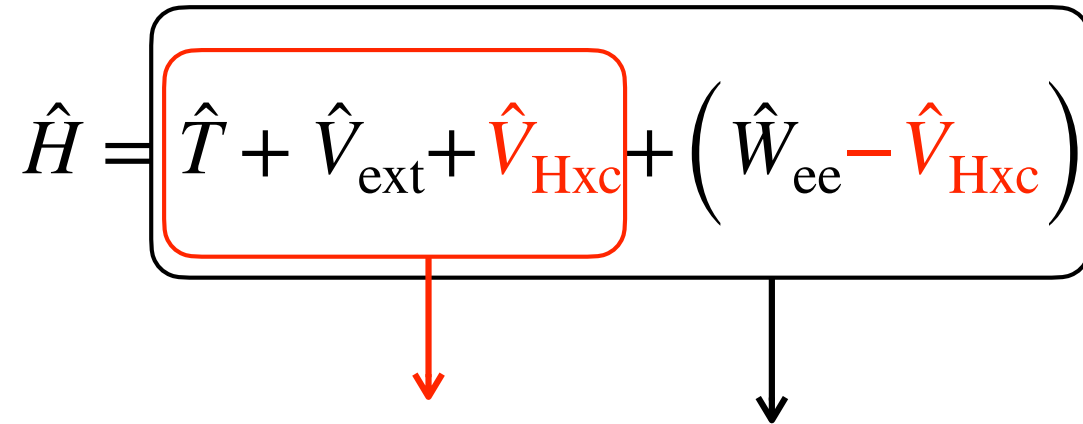
Local potential functional approach to DFT

$$\hat{H} = \hat{T} + \hat{V}_{\text{ext}} + \hat{V}_{\text{Hxc}} + \left(\hat{W}_{\text{ee}} - \hat{V}_{\text{Hxc}} \right)$$


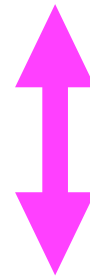
$$n_{\Phi_0^{\text{KS}}}(\mathbf{r}) = n_{\Psi_0}(\mathbf{r})$$

Exact density constraint

Local potential functional approach to DFT

$$\hat{H} = \hat{T} + \hat{V}_{\text{ext}} + \hat{V}_{\text{Hxc}} + \left(\hat{W}_{\text{ee}} - \hat{V}_{\text{Hxc}} \right)$$


$$n_{\Phi_0^{\text{KS}}}(\mathbf{r}) = n_{\Psi_0}(\mathbf{r})$$



Self-consistent scheme
(when approximations
are made)

$$n_{\Phi_0^{\text{KS}}[v_{\text{Hxc}}]}(\mathbf{r}) = n_{\Psi_0[v_{\text{Hxc}}]}(\mathbf{r})$$

Local potential functional approach to DFT

$$\hat{H} = \hat{T} + \hat{V}_{\text{ext}} + \hat{V}_{\text{Hxc}} + \left(\hat{W}_{\text{ee}} - \hat{V}_{\text{Hxc}} \right)$$

$n_{\Phi_0^{\text{KS}}[v_{\text{Hxc}}]}(\mathbf{r}) = n_{\Psi_0[v_{\text{Hxc}}]}(\mathbf{r})$

Approaching the ground state Ψ_0 of the full system is **computationally demanding**

Local potential functional approach to DFT

$$\hat{H} = \hat{T} + \hat{V}_{\text{ext}} + \hat{V}_{\text{Hxc}} + \left(\hat{W}_{\text{ee}} - \hat{V}_{\text{Hxc}} \right)$$

$n_{\Phi_0^{\text{KS}}[v_{\text{Hxc}}]}(\mathbf{r}) = n_{\Psi_0[v_{\text{Hxc}}]}(\mathbf{r})$

The approach becomes interesting if a *reduced-in-size approximation* to Ψ_0 can be used

Local potential functional approach to DFT

$$\hat{H} = \hat{T} + \hat{V}_{\text{ext}} + \hat{V}_{\text{Hxc}} + \left(\hat{W}_{\text{ee}} - \hat{V}_{\text{Hxc}} \right)$$
$$n_{\Phi_0^{\text{KS}}[v_{\text{Hxc}}]}(\mathbf{r}) = n_{\Psi_0[v_{\text{Hxc}}]}(\mathbf{r})$$

The approach becomes interesting if a **reduced-in-size approximation** to Ψ_0 can be used

Quantum embedding!



Condensed Matter > Strongly Correlated Electrons

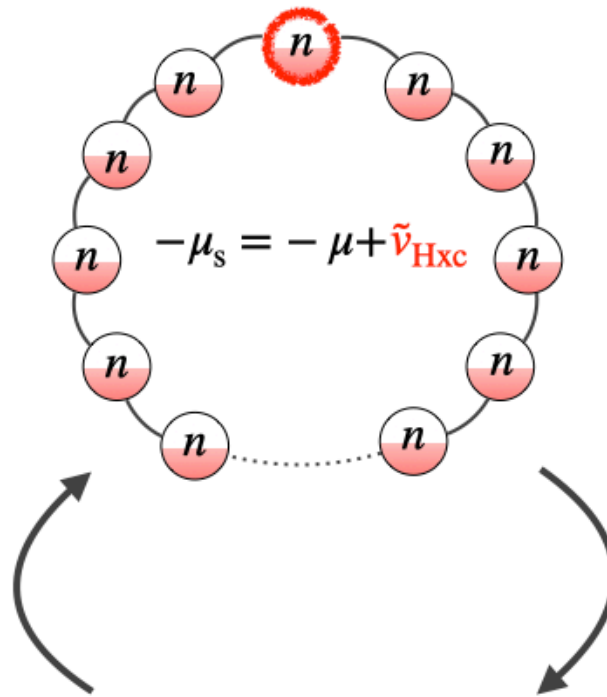
[Submitted on 16 Feb 2022]

Local Potential Functional Embedding Theory: A Self-Consistent Flavor of Density Functional Theory for Lattices without Density Functionals

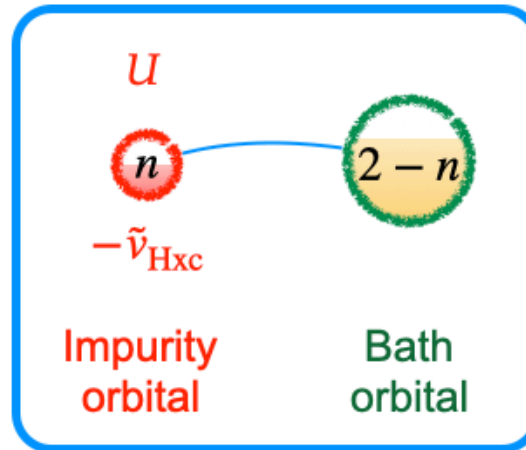
[Sajanthan Sekaran](#), [Matthieu Saubanère](#), [Emmanuel Fromager](#)

The recently proposed Householder transformed density-matrix functional embedding theory (Ht-DMFET) [Sekaran et al., Phys. Rev. B 104, 035121 (2021)], which is equivalent to (but formally simpler than) density matrix embedding theory (DMET) in the non-interacting case, is revisited from the perspective of density-functional theory (DFT). An in-principle-exact density-functional version of Ht-DMFET is derived for the one-dimensional Hubbard lattice with a single embedded impurity. On the basis of well-identified density-functional approximations, a local potential functional embedding theory (LPFET) is formulated and implemented. Even though LPFET performs better than Ht-DMFET in the low-density regime, in particular when electron correlation is strong, both methods are unable to describe the density-driven Mott-Hubbard transition, as expected. These results combined with our formally exact density-functional embedding theory reveal that a single statically embedded impurity can in principle describe the gap opening, provided that the complementary correlation potential (that describes the interaction of the embedding cluster with its environment, which is simply neglected in both Ht-DMFET and LPFET) exhibits a derivative discontinuity (DD) at half filling. The extension of LPFET to multiple impurities (which would enable to circumvent the modeling of DDs) and its generalization to quantum chemical Hamiltonians are left for future work.

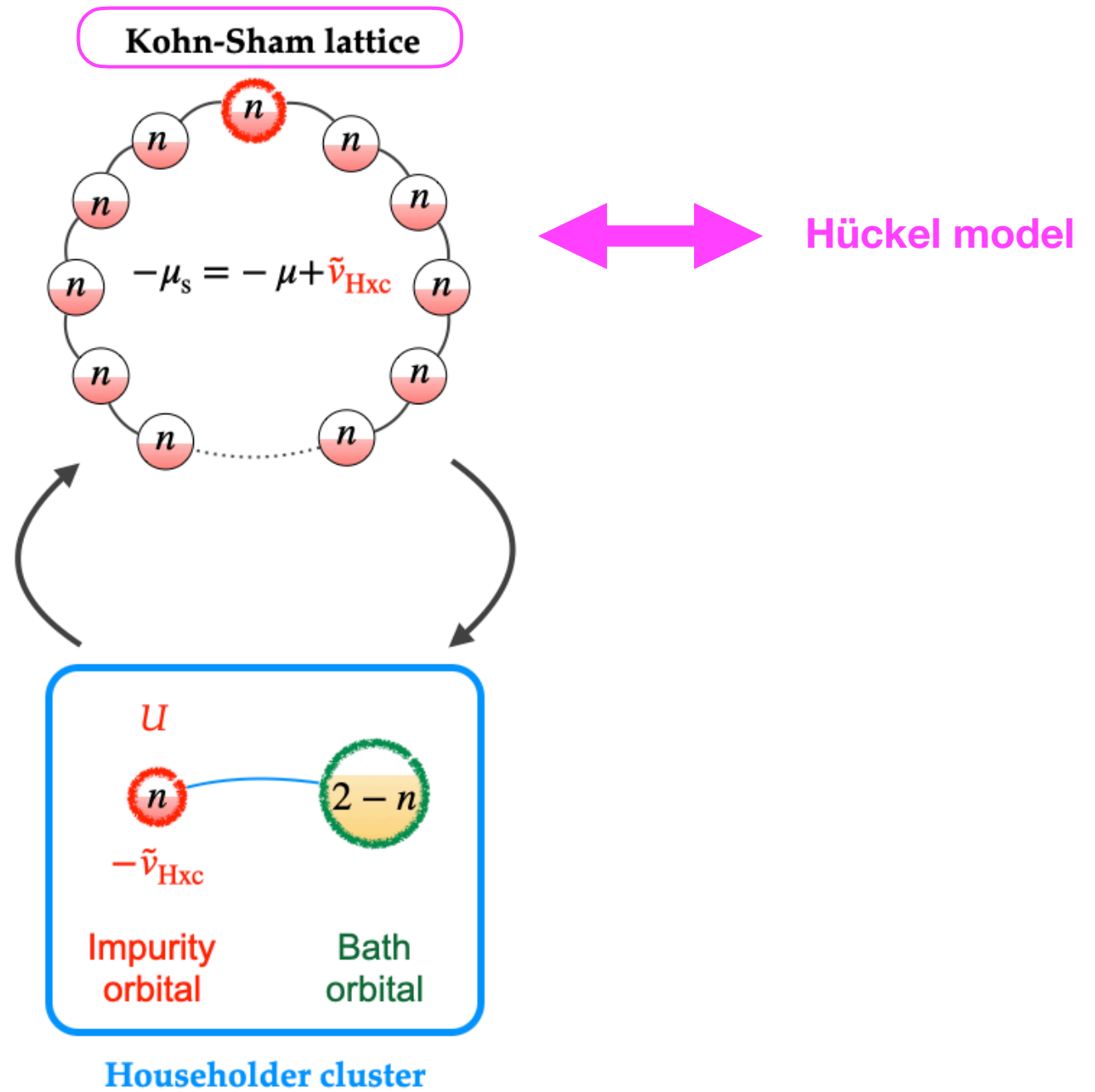
Kohn-Sham lattice



Application
to Hubbard lattices
(ring of hydrogen atoms)

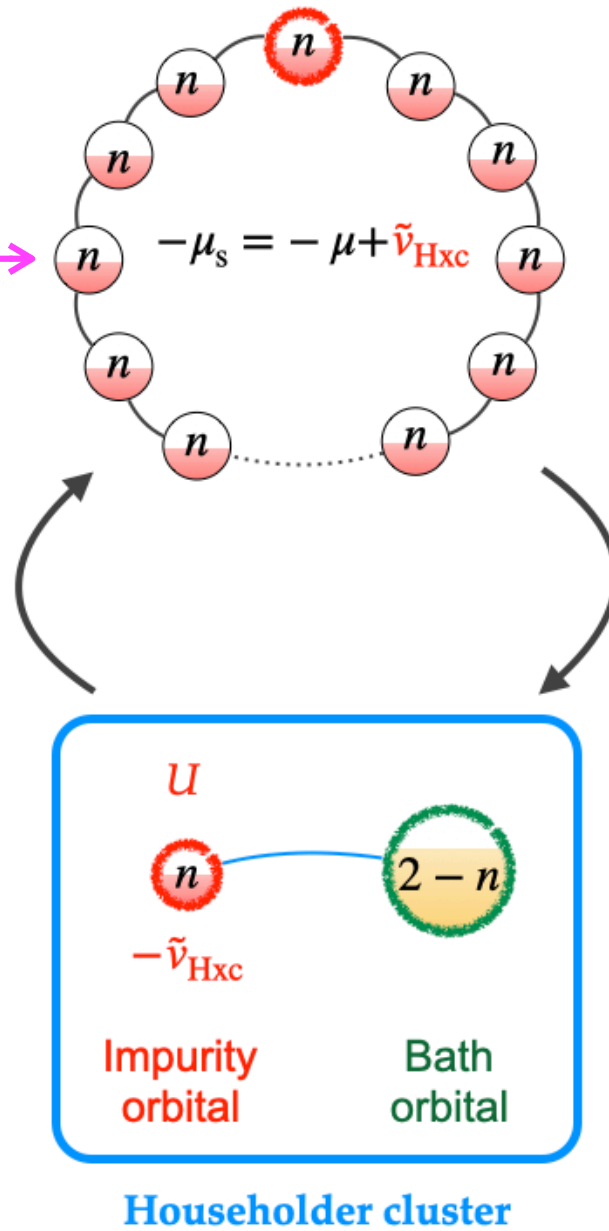


Householder cluster

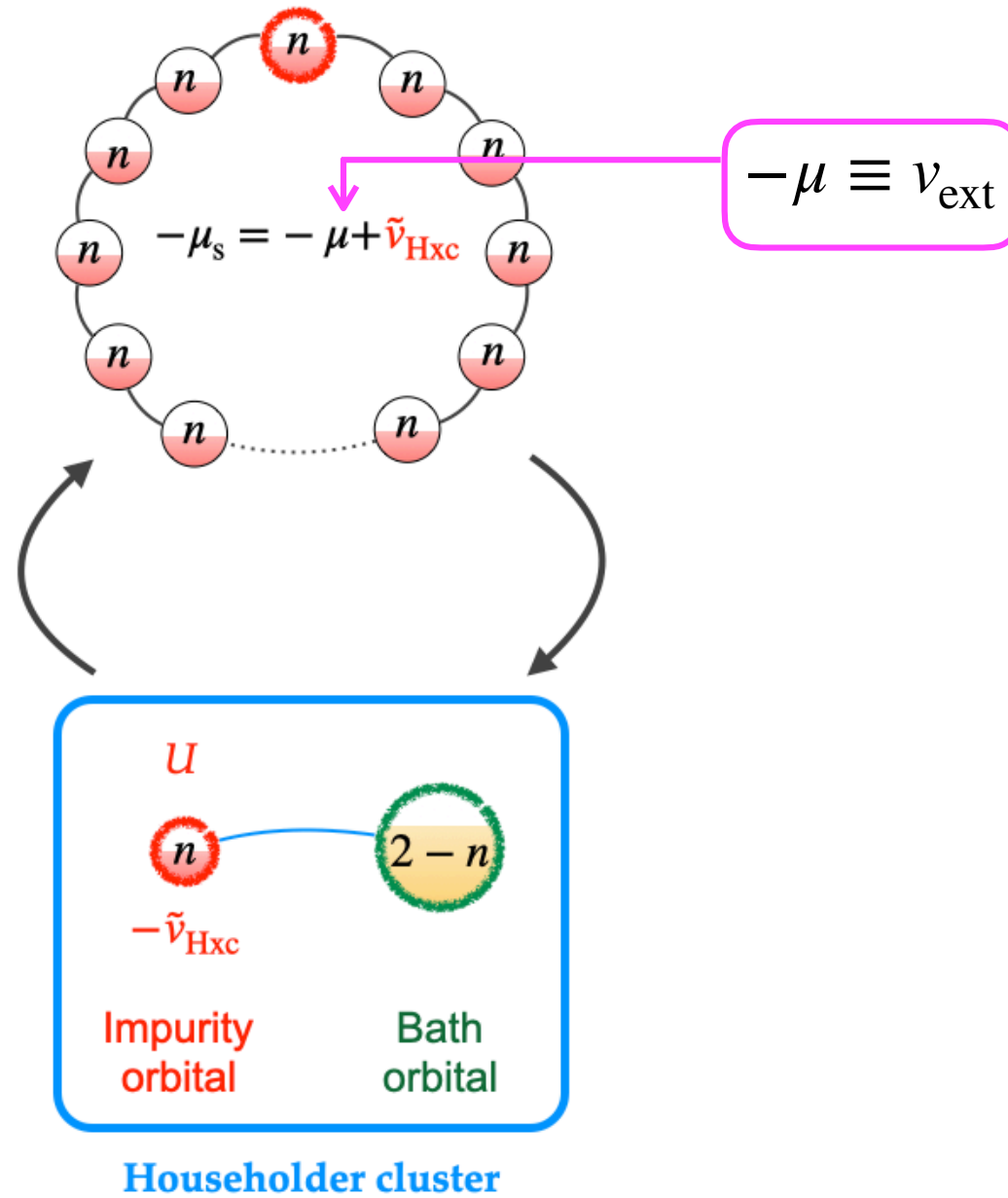


Kohn-Sham lattice

Uniform occupation n
of the atomic orbitals

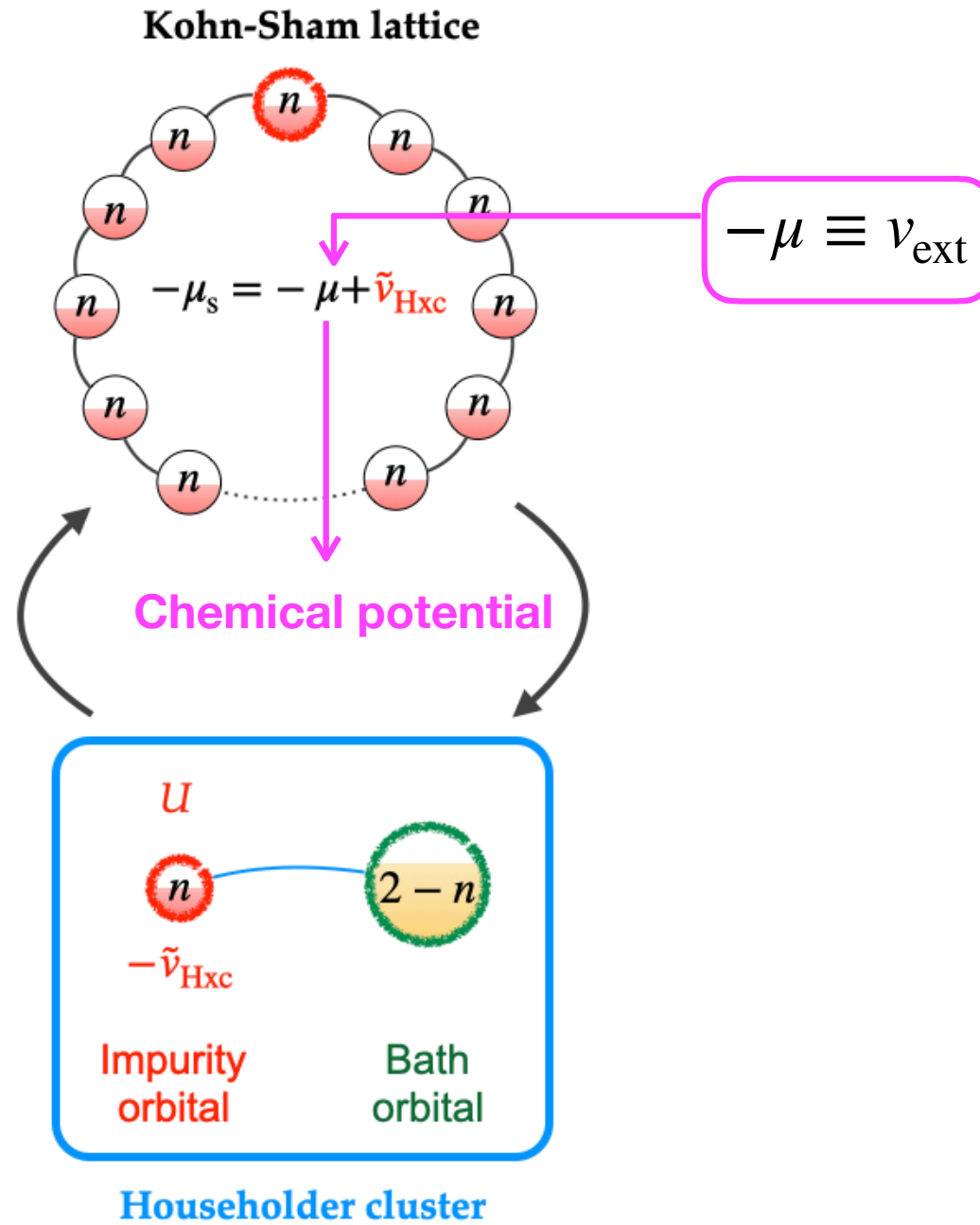


Kohn-Sham lattice

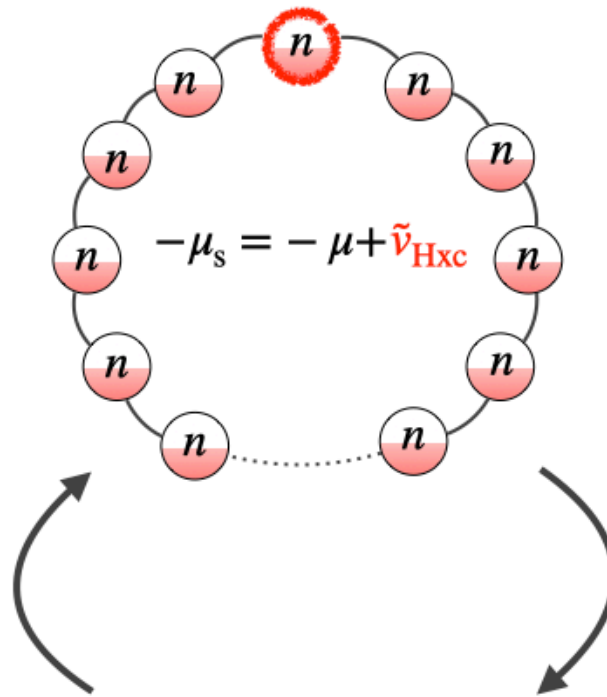


$$n = \frac{N}{L}$$

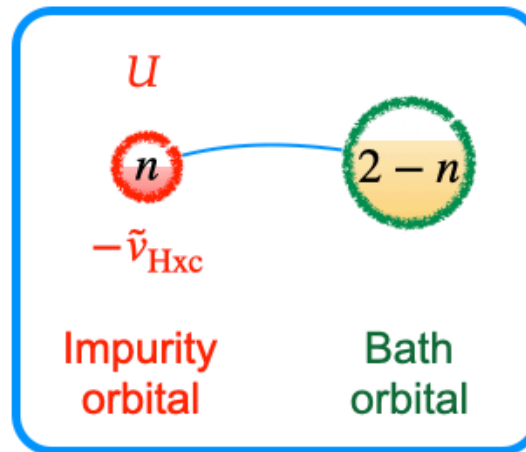
Fixed number of atomic orbitals



Kohn-Sham lattice

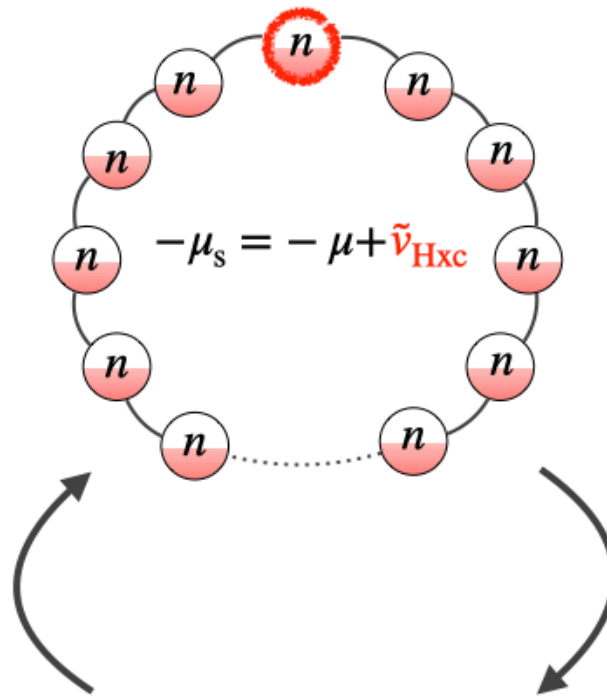


Embedding cluster



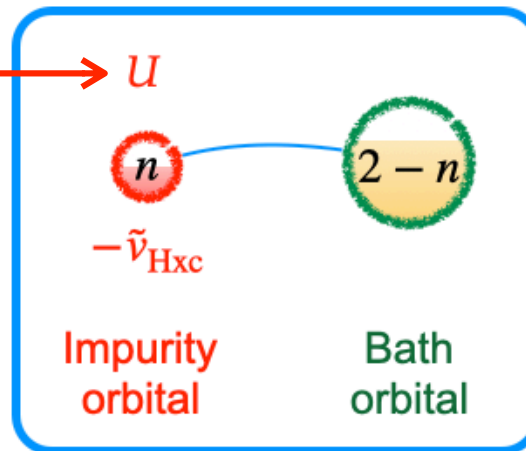
Householder cluster

Kohn-Sham lattice



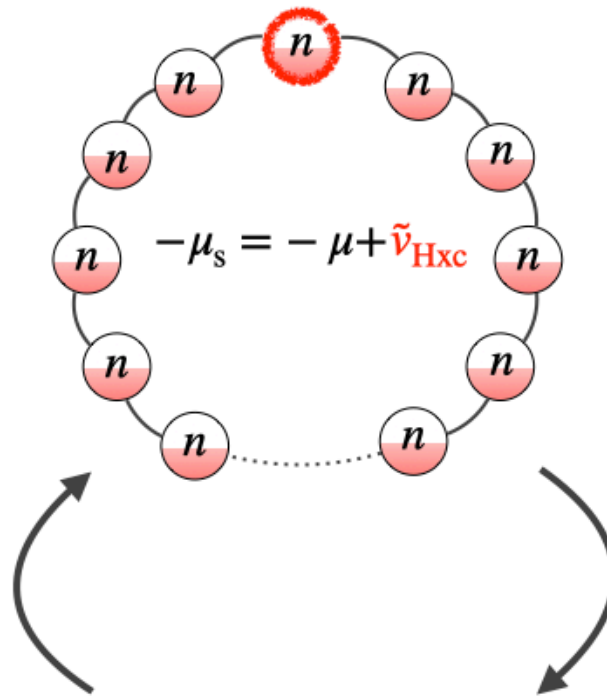
Two-electron repulsion

→ U

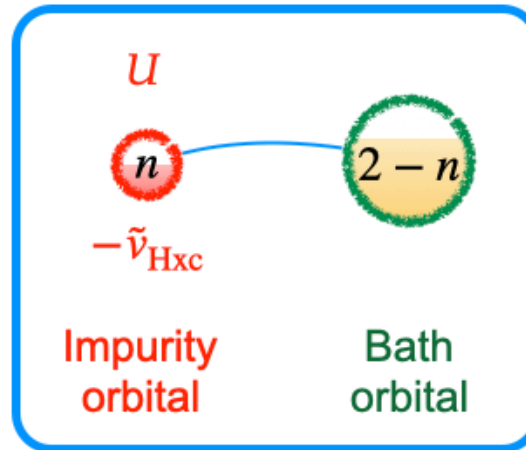


Householder cluster

Kohn-Sham lattice



Embedding cluster

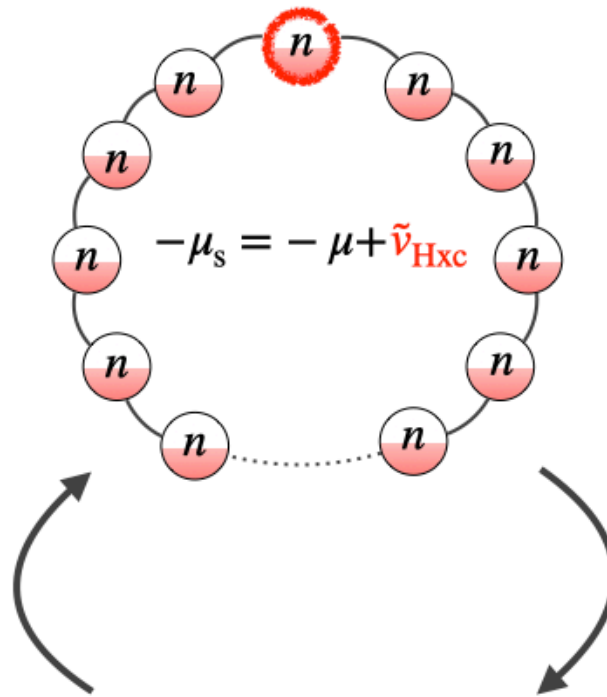


Householder cluster

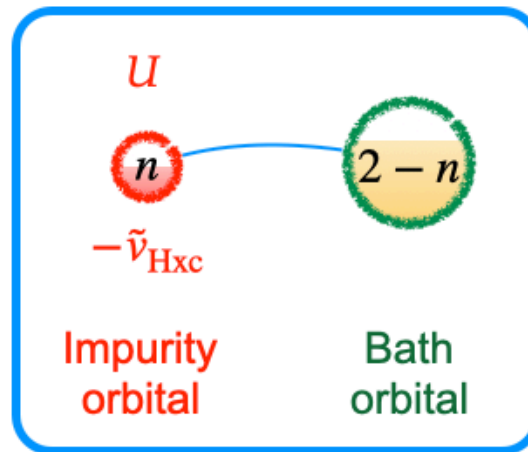
Reduced-in-size version of the full **interacting** lattice

$$\Psi_0 \rightarrow \Psi_0^{cluster}$$

Kohn-Sham lattice

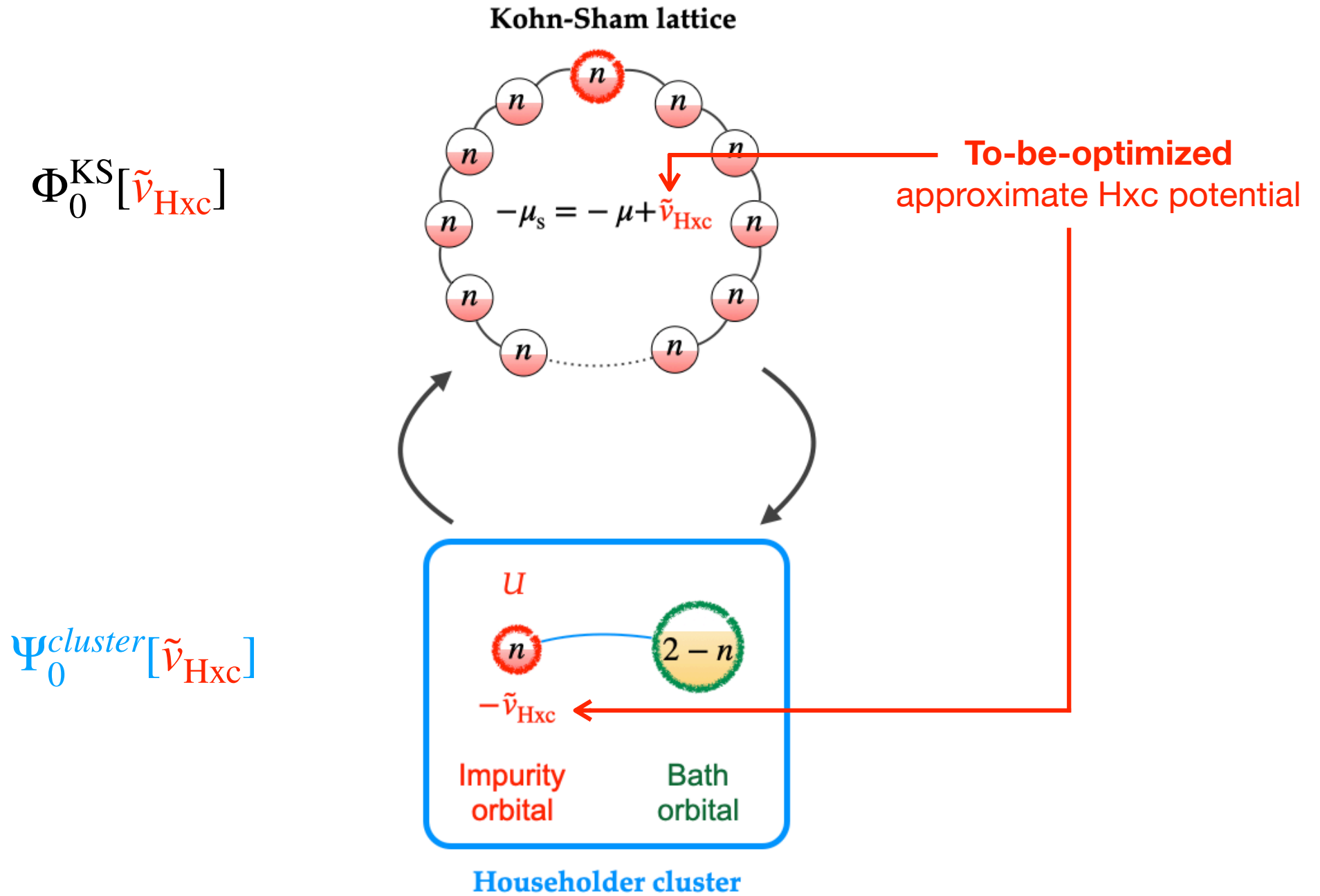


Embedding cluster

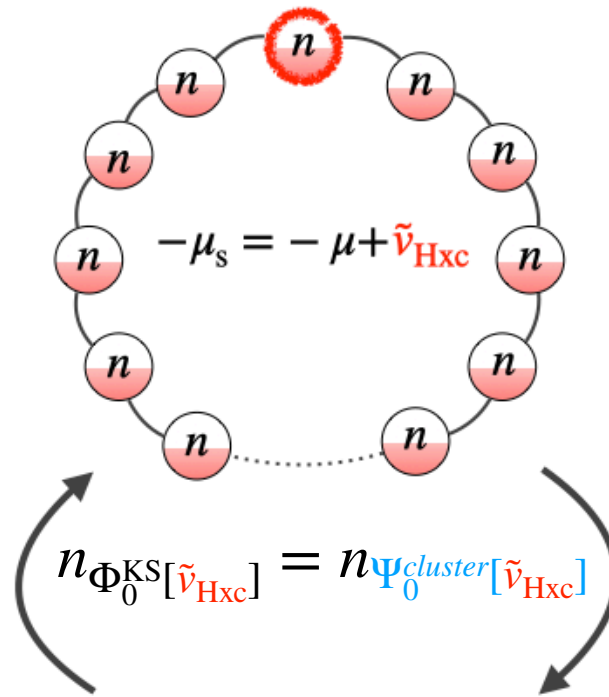


Householder cluster

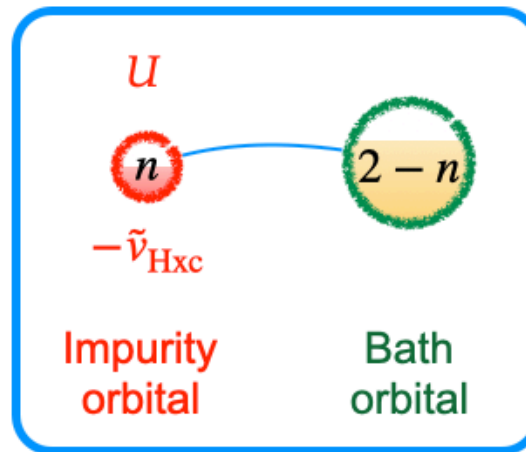
Can be solved exactly!



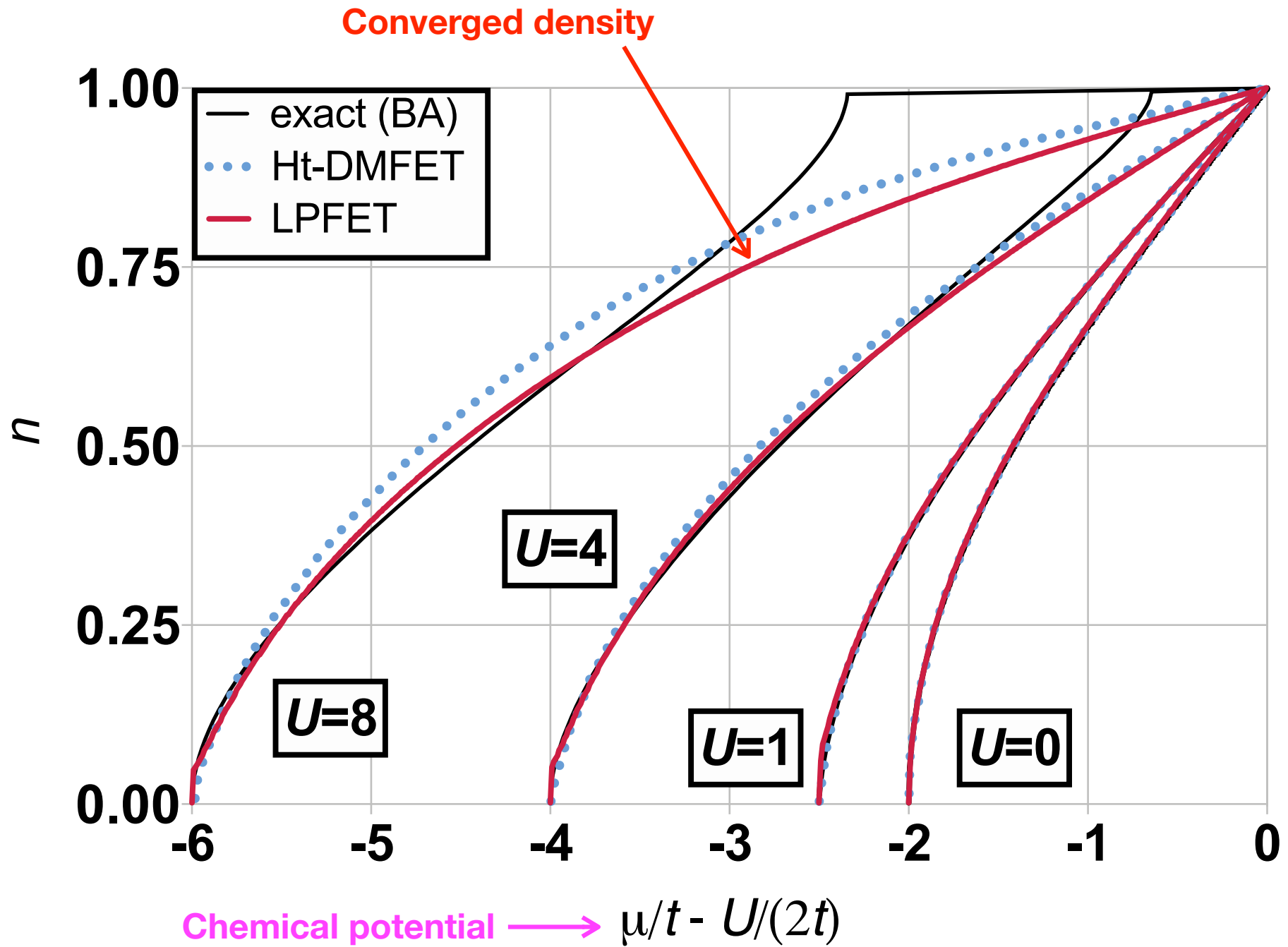
Kohn-Sham lattice

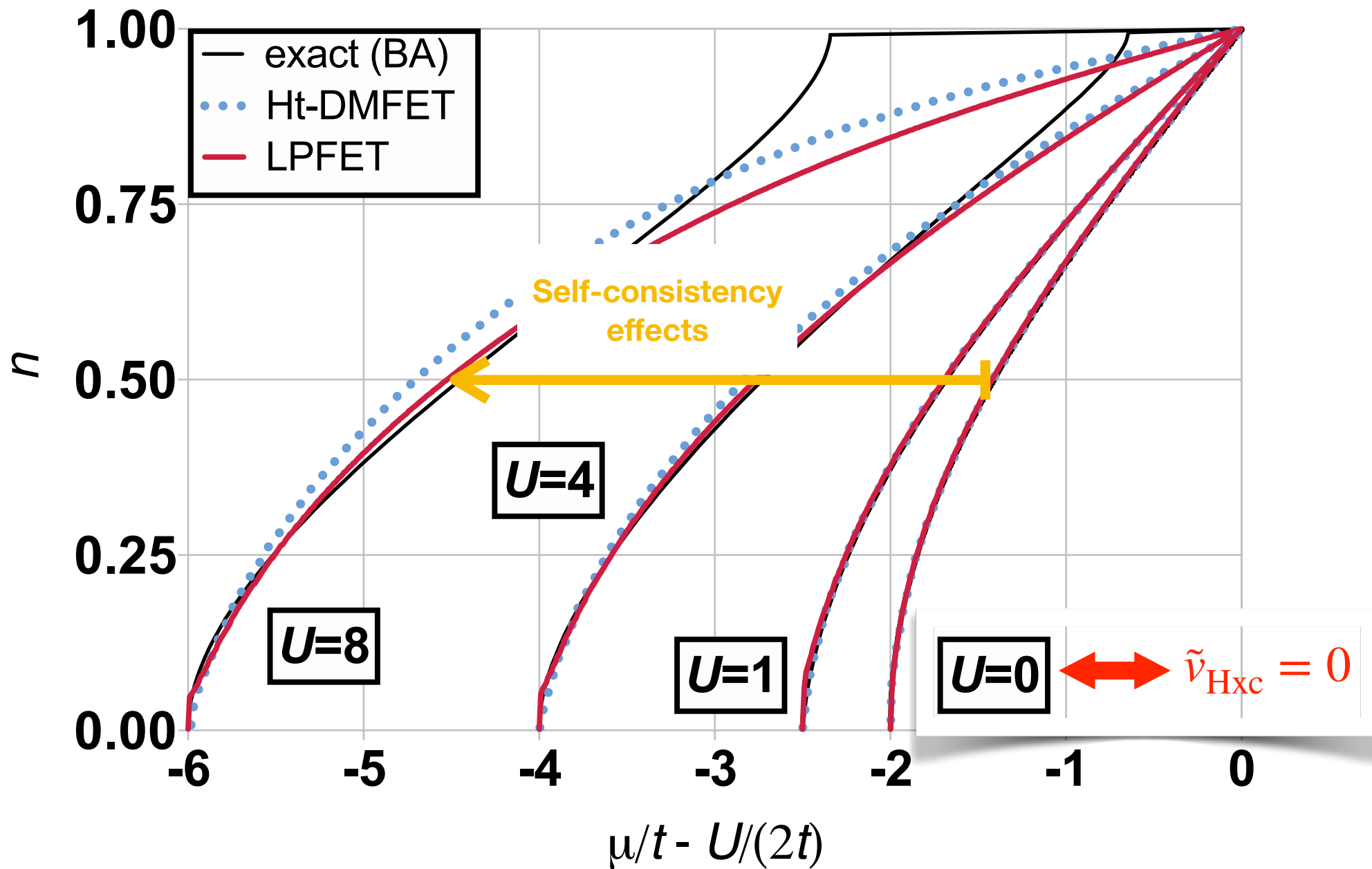


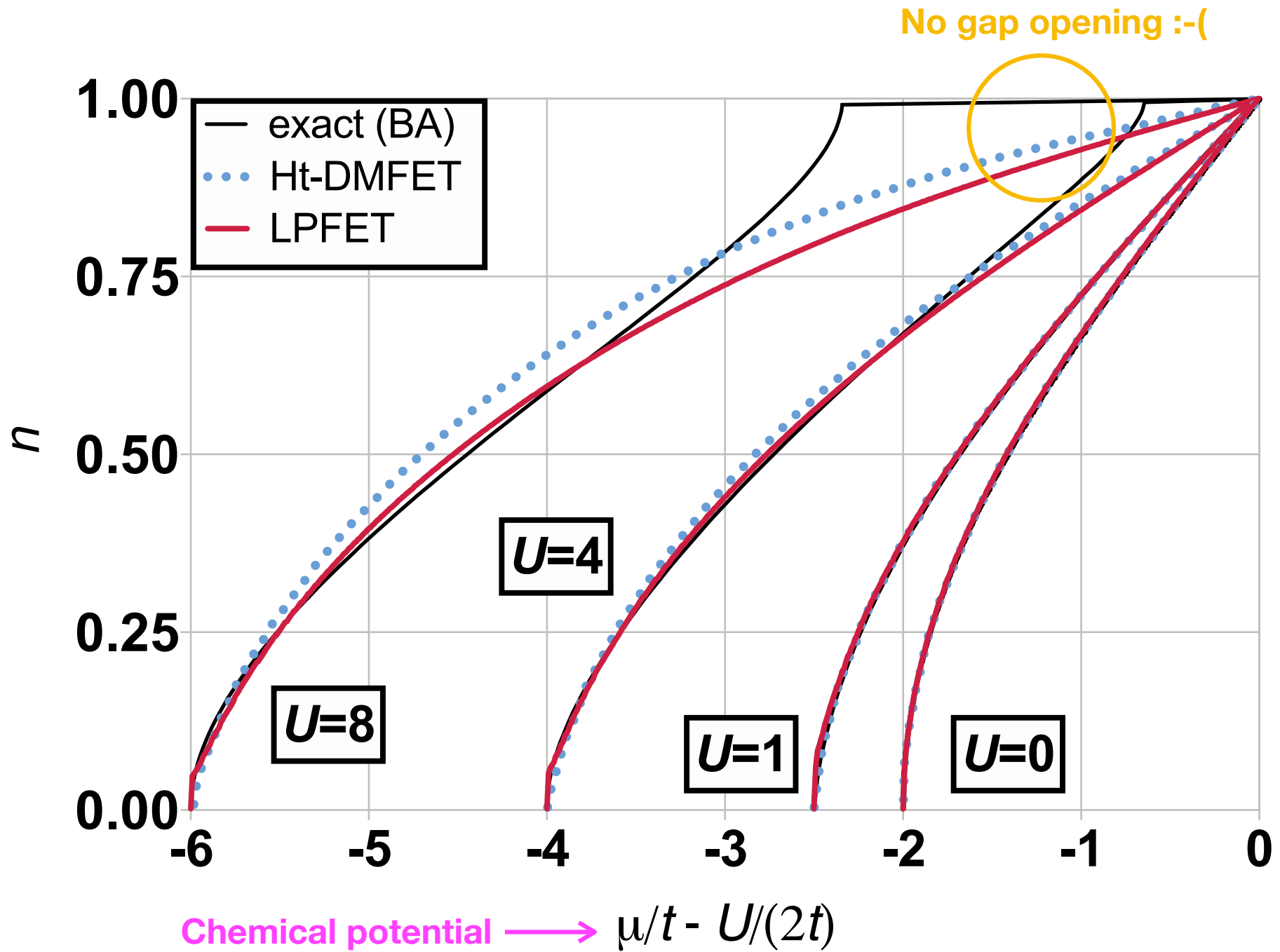
Self-consistency loop



Householder cluster

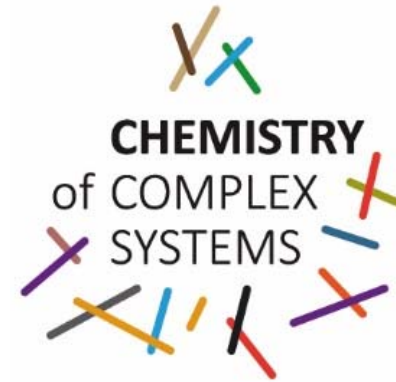






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