

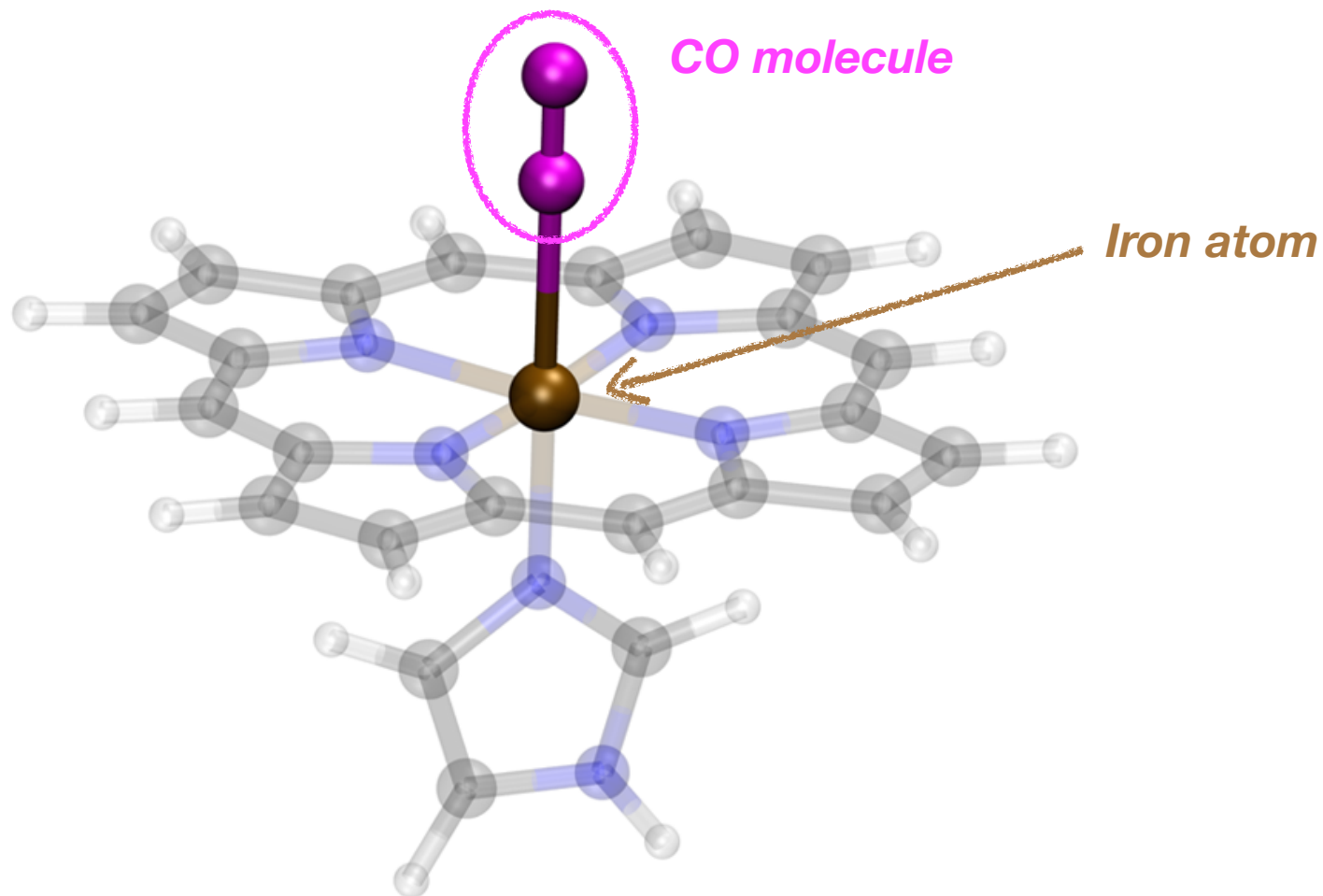
Quantum embedding in electronic structure theory

Part 2: Concept of embedding and its various formulations

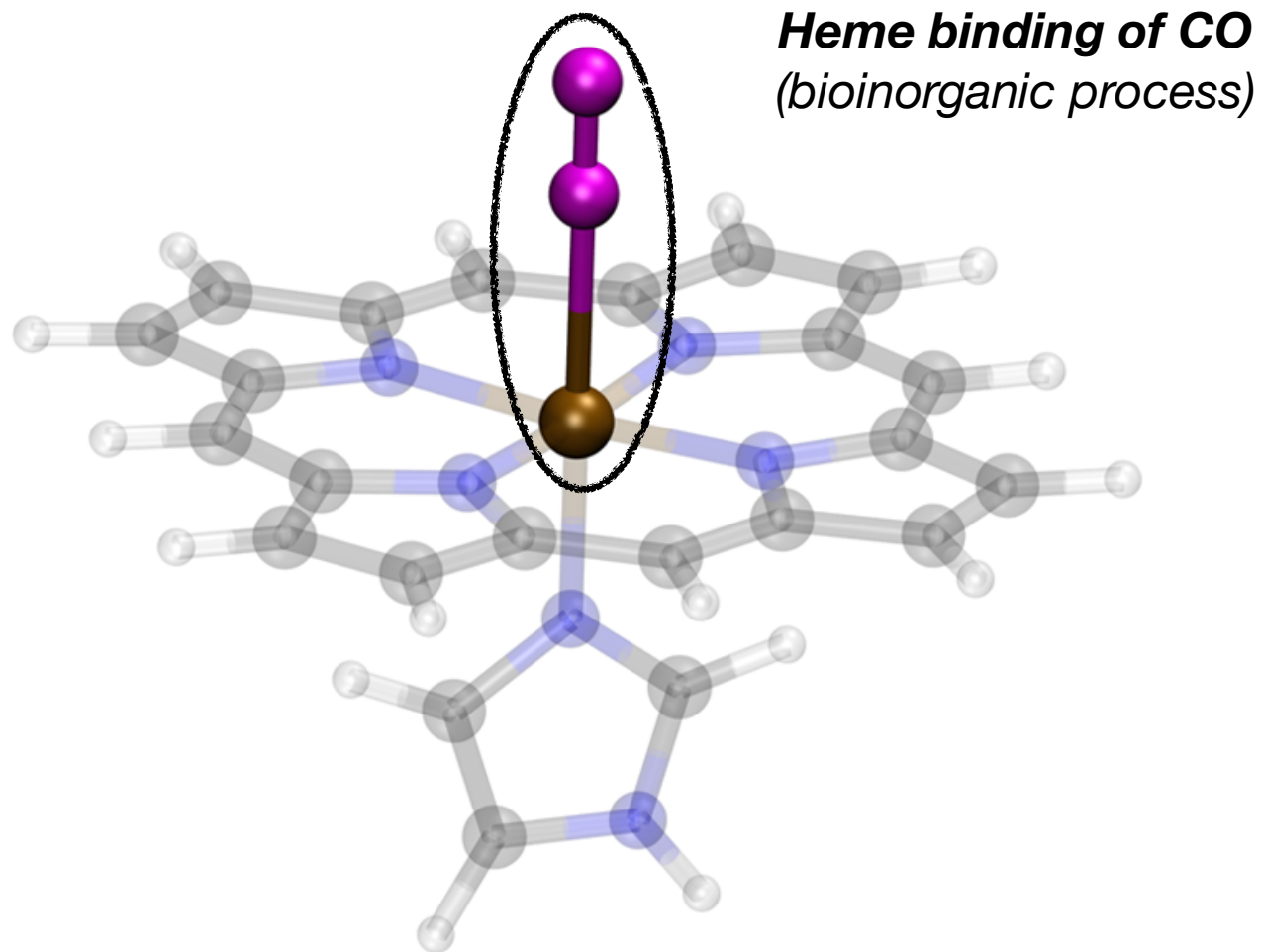
Emmanuel Fromager

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Université de Strasbourg, Strasbourg, France.*

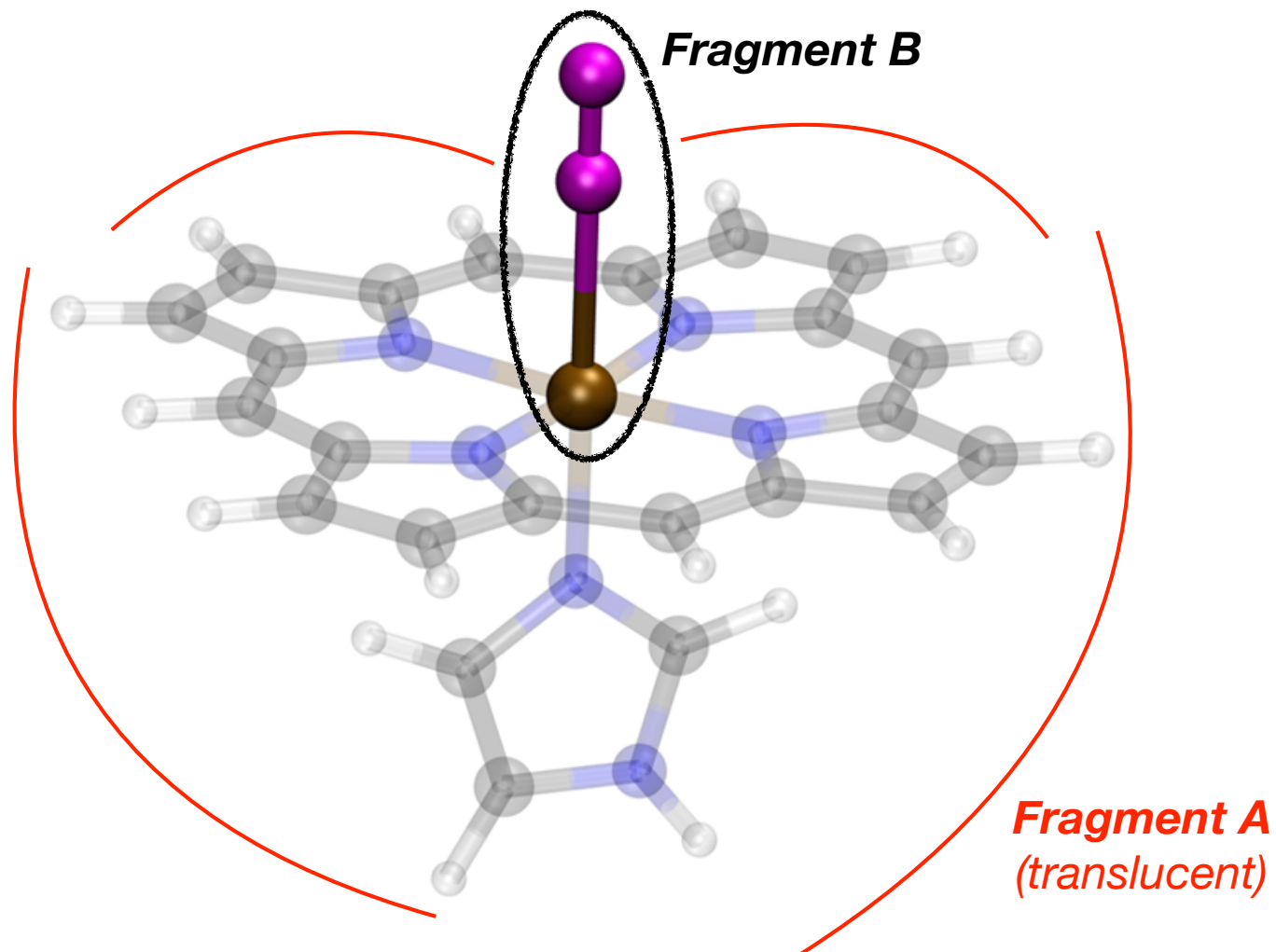
CO Binding in an Iron–Porphyrin Complex



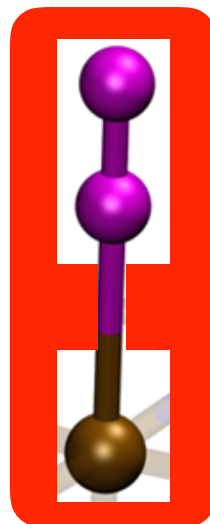
CO Binding in an Iron–Porphyrin Complex



CO Binding in an Iron–Porphyrin Complex

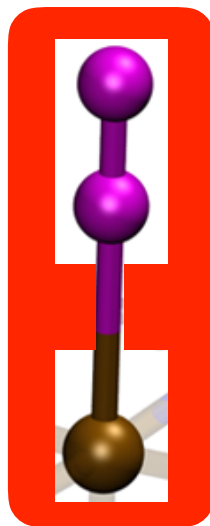


CO Binding in an Iron–Porphyrin Complex



To-be-described *embedded*
fragment B

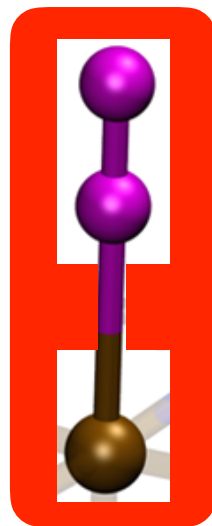
CO Binding in an Iron–Porphyrin Complex



To-be-described **embedded**
fragment **B**

$$\hat{H}^{emb} \Psi_B = \mathcal{E}_B \Psi_B$$

CO Binding in an Iron–Porphyrin Complex



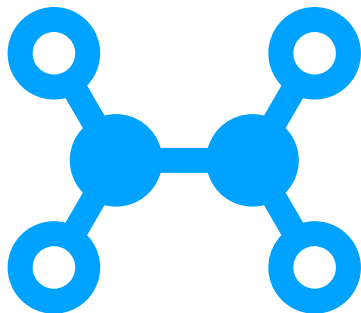
To-be-described **embedded**
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$$\hat{H}^{emb}\Psi_B = \mathcal{E}_B\Psi_B$$



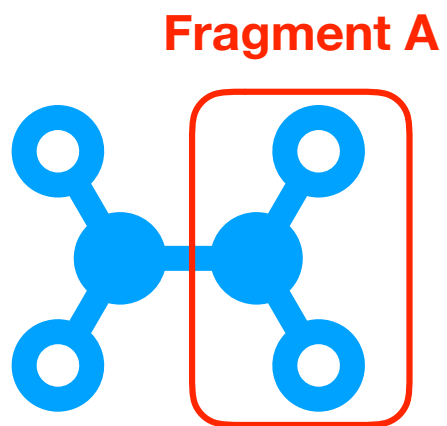
Strategy 1: DFT-based quantum embedding approach

Frozen density embedding theory



$$n_0(\mathbf{r}) = ?$$

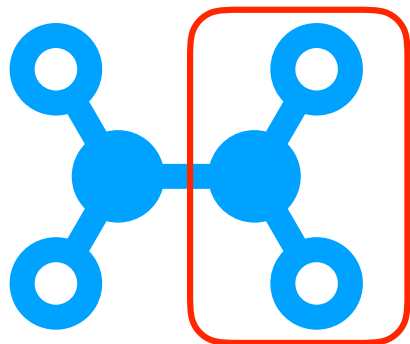
Frozen density embedding theory



$n_A(\mathbf{r})$ ← Density giving an approximate description of a molecular *fragment*

Frozen density embedding theory

Fragment A



$$n_A(\mathbf{r})$$

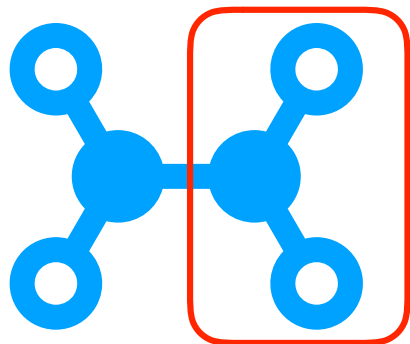
Density giving an approximate description of a molecular *fragment*

Total number of electrons in the molecule

$$\int d\mathbf{r} n_A(\mathbf{r}) = N_A < N$$

Frozen density embedding theory

Fragment A



$n_A(\mathbf{r})$

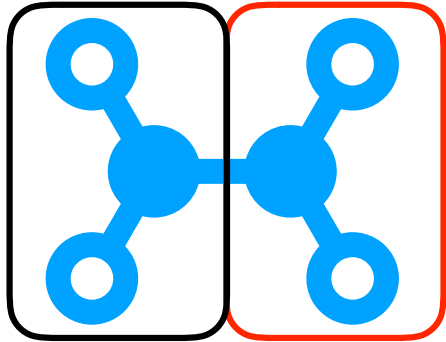
Density giving an
approximate description
of a molecular *fragment*



Fixed in the theory
(hence the name “*frozen density*”)

Frozen density embedding theory

Fragment B **Fragment A**

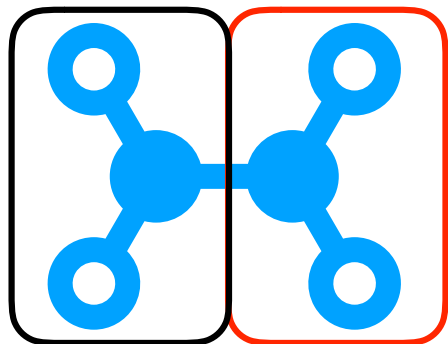


$n_A(\mathbf{r})$ *known*

$$n_B(\mathbf{r}) = n_0(\mathbf{r}) - n_A(\mathbf{r}) = ?$$

Frozen density embedding theory

Fragment B **Fragment A**

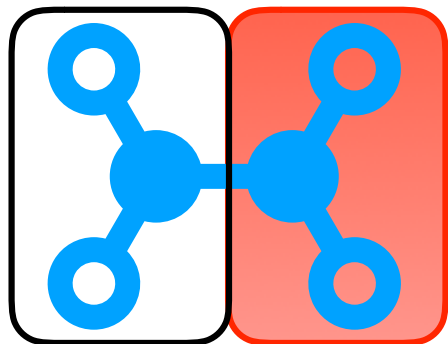


$n_A(\mathbf{r})$ *known*

$$n_B(\mathbf{r}) = n_0(\mathbf{r}) - n_A(\mathbf{r}) \equiv n_{\Psi_B} = ?$$

Frozen density embedding theory

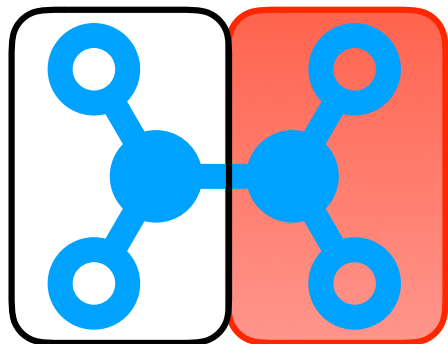
Fragment B



unknown $\Psi_B = ?$ $n_A(\mathbf{r})$ *known*

Frozen density embedding theory

Fragment B



unknown $\Psi_B = ?$ $n_A(\mathbf{r})$ *known*


Fragment B is **embedded** into the density-functional fragment A

Variational principle for fragment B

$$E_0 = \min_n \left\{ F[n] + \int d\mathbf{r} v_{\text{ext}}(\mathbf{r})n(\mathbf{r}) \right\}$$

Variational principle for fragment B

$$E_0 = \min_n \left\{ F[n] + \int d\mathbf{r} v_{\text{ext}}(\mathbf{r})n(\mathbf{r}) \right\}$$
$$= \min_{n_B \rightarrow N - N_A} \left\{ F[n_B + n_A] + \int d\mathbf{r} v_{\text{ext}}(\mathbf{r})n_B(\mathbf{r}) \right\} + \int d\mathbf{r} v_{\text{ext}}(\mathbf{r})n_A(\mathbf{r})$$

**Frozen (i.e., fixed)**

Variational principle for fragment B

$$\begin{aligned} E_0 &= \min_n \left\{ F[n] + \int d\mathbf{r} v_{\text{ext}}(\mathbf{r})n(\mathbf{r}) \right\} \\ &= \min_{n_B \rightarrow N - N_A} \left\{ F[n_B + n_A] + \int d\mathbf{r} v_{\text{ext}}(\mathbf{r})n_B(\mathbf{r}) \right\} + \int d\mathbf{r} v_{\text{ext}}(\mathbf{r})n_A(\mathbf{r}) \\ &= \min_{n_B \rightarrow N - N_A} \left\{ F[n_B] + \int d\mathbf{r} v_{\text{ext}}(\mathbf{r})n_B(\mathbf{r}) + \Delta F[n_B, n_A] \right\} + \int d\mathbf{r} v_{\text{ext}}(\mathbf{r})n_A(\mathbf{r}) + F[n_A] \end{aligned}$$

$$F[n_B + n_A] = F[n_B] + F[n_A] + \Delta F[n_B, n_A]$$

Formal decomposition

Variational principle for fragment B

$$\begin{aligned} E_0 &= \min_n \left\{ F[n] + \int d\mathbf{r} v_{\text{ext}}(\mathbf{r})n(\mathbf{r}) \right\} \\ &= \min_{n_B \rightarrow N - N_A} \left\{ F[n_B + n_A] + \int d\mathbf{r} v_{\text{ext}}(\mathbf{r})n_B(\mathbf{r}) \right\} + \int d\mathbf{r} v_{\text{ext}}(\mathbf{r})n_A(\mathbf{r}) \\ &= \min_{n_B \rightarrow N - N_A} \left\{ F[n_B] + \int d\mathbf{r} v_{\text{ext}}(\mathbf{r})n_B(\mathbf{r}) + \Delta F[n_B, n_A] \right\} + \int d\mathbf{r} v_{\text{ext}}(\mathbf{r})n_A(\mathbf{r}) + F[n_A] \end{aligned}$$

*Bifunctional
describing the **coupling**
between the two fragments*

$$F[n_B + n_A] = F[n_B] + F[n_A] + \Delta F[n_B, n_A]$$

Variational principle for fragment B

$$\begin{aligned}
 E_0 &= \min_n \left\{ F[n] + \int d\mathbf{r} v_{\text{ext}}(\mathbf{r})n(\mathbf{r}) \right\} \\
 &= \min_{n_B \rightarrow N - N_A} \left\{ F[n_B + n_A] + \int d\mathbf{r} v_{\text{ext}}(\mathbf{r})n_B(\mathbf{r}) \right\} + \int d\mathbf{r} v_{\text{ext}}(\mathbf{r})n_A(\mathbf{r}) \\
 &= \min_{n_B \rightarrow N - N_A} \left\{ F[n_B] + \int d\mathbf{r} v_{\text{ext}}(\mathbf{r})n_B(\mathbf{r}) + \Delta F[n_B, n_A] \right\} + \int d\mathbf{r} v_{\text{ext}}(\mathbf{r})n_A(\mathbf{r}) + F[n_A] \\
 &\quad \downarrow \\
 &= \min_{n_B \rightarrow N - N_A} \left\{ \min_{\Psi \rightarrow n_B} \left\{ \langle \Psi | \hat{T} + \hat{W}_{\text{ee}} | \Psi \rangle \right\} + \int d\mathbf{r} v_{\text{ext}}(\mathbf{r})n_B(\mathbf{r}) + \Delta F[n_B, n_A] \right\}
 \end{aligned}$$

Variational principle for fragment B

$$\begin{aligned}
 E_0 &= \min_n \left\{ F[n] + \int d\mathbf{r} v_{\text{ext}}(\mathbf{r})n(\mathbf{r}) \right\} \\
 &= \min_{n_B \rightarrow N - N_A} \left\{ F[n_B + n_A] + \int d\mathbf{r} v_{\text{ext}}(\mathbf{r})n_B(\mathbf{r}) \right\} + \int d\mathbf{r} v_{\text{ext}}(\mathbf{r})n_A(\mathbf{r}) \\
 &= \min_{n_B \rightarrow N - N_A} \left\{ F[n_B] + \int d\mathbf{r} v_{\text{ext}}(\mathbf{r})n_B(\mathbf{r}) + \Delta F[n_B, n_A] \right\} + \int d\mathbf{r} v_{\text{ext}}(\mathbf{r})n_A(\mathbf{r}) + F[n_A] \\
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 &= \min_{n_B \rightarrow N - N_A} \left\{ \min_{\Psi \rightarrow n_B} \left\{ \langle \Psi | \hat{T} + \hat{W}_{\text{ee}} | \Psi \rangle + \int d\mathbf{r} v_{\text{ext}}(\mathbf{r})n_{\Psi}(\mathbf{r}) + \Delta F[n_{\Psi}, n_A] \right\} \right\}
 \end{aligned}$$

Variational principle for fragment B

$$E_0 - \int d\mathbf{r} v_{\text{ext}}(\mathbf{r}) n_A(\mathbf{r}) - F[n_A] = \min_{\Psi \rightarrow N - N_A} \left\{ \langle \Psi | \hat{H} | \Psi \rangle + \Delta F[n_\Psi, n_A] \right\}$$

Variational principle for fragment B

$$E_0 - \int d\mathbf{r} v_{\text{ext}}(\mathbf{r}) n_A(\mathbf{r}) - F[n_A] = \min_{\Psi \rightarrow N - N_A} \left\{ \langle \Psi | \hat{H} | \Psi \rangle + \Delta F[n_\Psi, n_A] \right\}$$



Analogous to solving
the Schrödinger equation for **fragment B**
in the presence of **fragment A**

Variational principle for fragment B

$$E_0 - \int d\mathbf{r} v_{\text{ext}}(\mathbf{r}) n_A(\mathbf{r}) - F[n_A] = \min_{\Psi \rightarrow N - N_A} \left\{ \langle \Psi | \hat{H} | \Psi \rangle + \Delta F[n_\Psi, n_A] \right\}$$

*Embedding
density functional*

$$\Delta F[n_B, n_A] = F[n_B + n_A] - F[n_B] - F[n_A]$$

Variational principle for fragment B

$$E_0 - \int d\mathbf{r} v_{\text{ext}}(\mathbf{r}) n_A(\mathbf{r}) - F[n_A] = \min_{\Psi \rightarrow N - N_A} \left\{ \langle \Psi | \hat{H} | \Psi \rangle + \Delta F[n_\Psi, n_A] \right\}$$

$$\Delta F[n_B, n_A] = F[n_B + n_A] - F[n_B] - F[n_A]$$

$$= E_{\text{Hxc}}[n_B + n_A] - E_{\text{Hxc}}[n_B] - E_{\text{Hxc}}[n_A]$$

$$+ T_s[n_B + n_A] - T_s[n_B] - T_s[n_A]$$

KS decompositions

$$F[n] = T_s[n] + E_{\text{Hxc}}[n]$$

Variational principle for fragment B

$$E_0 - \int d\mathbf{r} v_{\text{ext}}(\mathbf{r}) n_A(\mathbf{r}) - F[n_A] = \min_{\Psi \rightarrow N - N_A} \left\{ \langle \Psi | \hat{H} | \Psi \rangle + \Delta F[n_\Psi, n_A] \right\}$$

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*Evaluated from
xc density-functional
approximations (DFAs)*

Variational principle for fragment B

$$E_0 - \int d\mathbf{r} v_{\text{ext}}(\mathbf{r}) n_A(\mathbf{r}) - F[n_A] = \min_{\Psi \rightarrow N - N_A} \left\{ \langle \Psi | \hat{H} | \Psi \rangle + \Delta F[n_\Psi, n_A] \right\}$$

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$$+ T_s[n_B + n_A] - T_s[n_B] - T_s[n_A] \quad \text{Described with KS orbitals} \\ \text{in KS-DFT}$$

Variational principle for fragment B

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$$+ T_s[n_B + n_A] - T_s[n_B] - T_s[n_A]$$

*Implicit functional
of the density*



*Described with KS orbitals
in KS-DFT*

Variational principle for fragment B

$$E_0 - \int d\mathbf{r} v_{\text{ext}}(\mathbf{r}) n_A(\mathbf{r}) - F[n_A] = \min_{\Psi \rightarrow N - N_A} \left\{ \langle \Psi | \hat{H} | \Psi \rangle + \Delta F[n_\Psi, n_A] \right\}$$

$$\Delta F[n_B, n_A] = F[n_B + n_A] - F[n_B] - F[n_A]$$

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$$+ T_s[n_B + n_A] - T_s[n_B] - T_s[n_A]$$

*More difficult
to approximate with
density functionals*

Variational principle for fragment B

$$E_0 - \int d\mathbf{r} v_{\text{ext}}(\mathbf{r}) n_A(\mathbf{r}) - F[n_A] = \min_{\Psi \rightarrow N - N_A} \left\{ \langle \Psi | \hat{H} | \Psi \rangle + \Delta F[n_\Psi, n_A] \right\}$$

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$$+ T_s[n_B + n_A] - T_s[n_B] - T_s[n_A]$$

challenging task!



*More difficult
to approximate with
density functionals*

Strategy 2: Quantum embedding in the N -electron space

Strategy 2: Quantum embedding in the N -electron space

*We need to write the Hamiltonian in **second quantization** for that purpose*

Second-quantized encoding of many-electron wave functions

Hydrogen molecule in the dissociation limit:

$$\sqrt{2}\Psi_0(\mathbf{r}_1, \mathbf{r}_2) = \varphi_{1\sigma_g}(\mathbf{r}_1)\varphi_{1\sigma_g}(\mathbf{r}_2) - \varphi_{1\sigma_u}(\mathbf{r}_1)\varphi_{1\sigma_u}(\mathbf{r}_2)$$

Second-quantized encoding of many-electron wave functions

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$$\sqrt{2}|\Psi_0\rangle = \left| \begin{array}{c} \text{---} \varphi_{1\sigma_u} \\ \uparrow\downarrow \varphi_{1\sigma_g} \end{array} \right\rangle - \left| \begin{array}{c} \uparrow\downarrow \varphi_{1\sigma_u} \\ \text{---} \varphi_{1\sigma_g} \end{array} \right\rangle$$

Dirac notation

Second-quantized encoding of many-electron wave functions

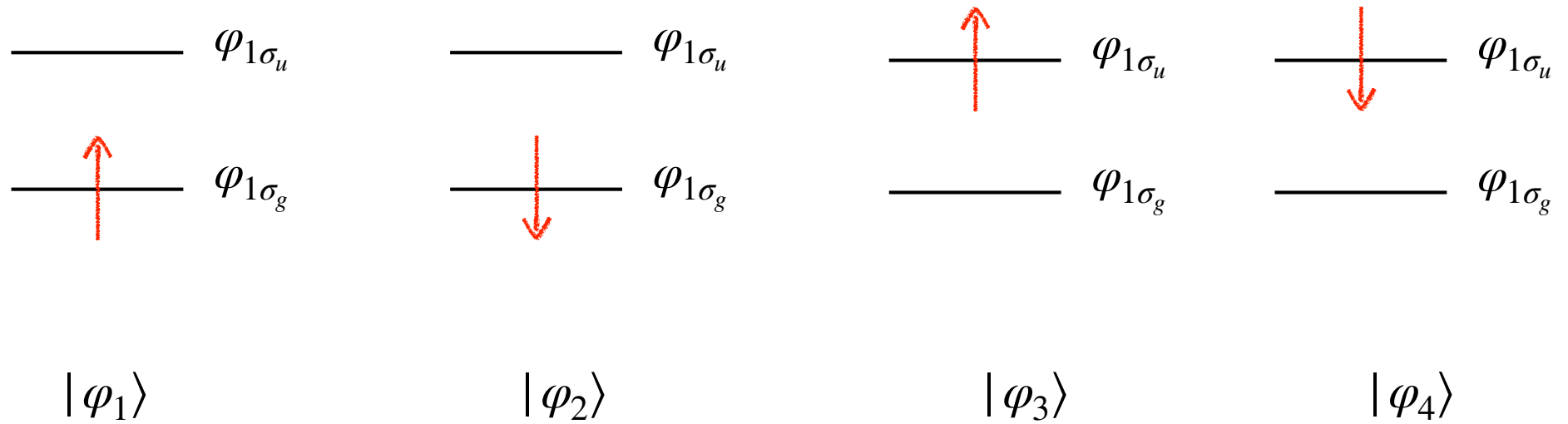
———— $\varphi_{1\sigma_u}$

Zero-electron (vacuum) quantum state

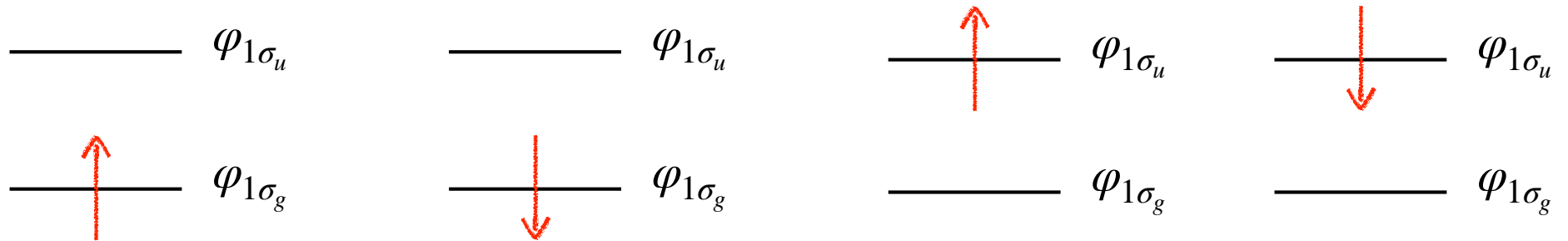
———— $\varphi_{1\sigma_g}$

$|\text{vac}\rangle$

One-electron quantum states



One-electron quantum states



$|\varphi_1\rangle$

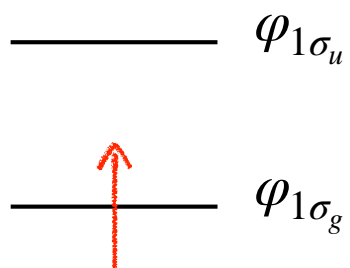
$|\varphi_2\rangle$

$|\varphi_3\rangle$

$|\varphi_4\rangle$

Spin-orbitals

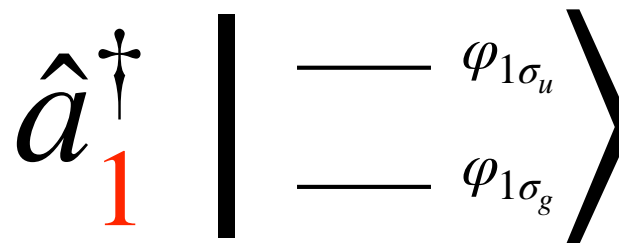
One-electron quantum states



$|\varphi_1\rangle$



Second quantization
notation



One-electron quantum states



$|\varphi_1\rangle$

Quantum operator that **creates**
an electron and put it in the spin-orbital φ_1

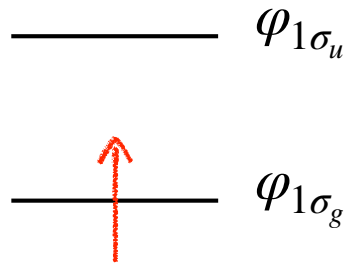
One-electron quantum states



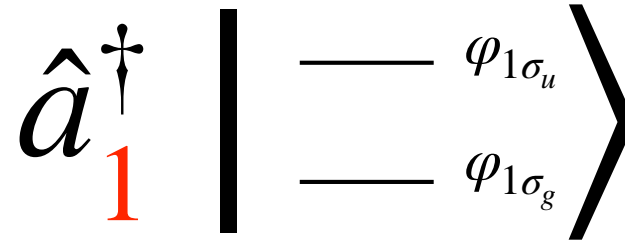
$|\varphi_1\rangle$

Creation operator

One-electron quantum states



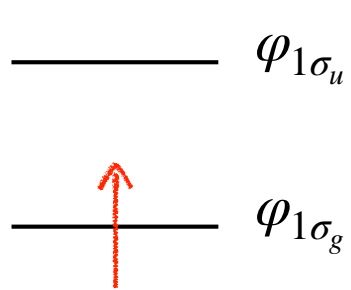
$|\varphi_1\rangle$



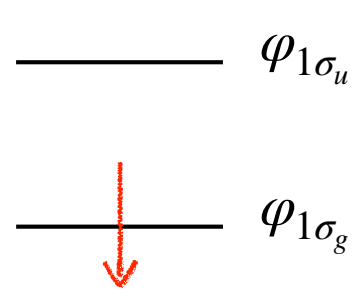
$$|\varphi_1\rangle \equiv \hat{a}_1^\dagger |\text{vac}\rangle$$

*Second-quantized
notation*

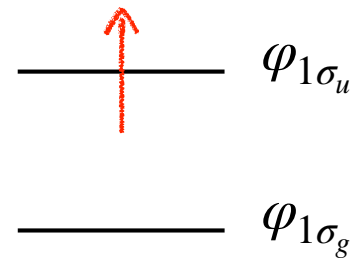
One-electron quantum states



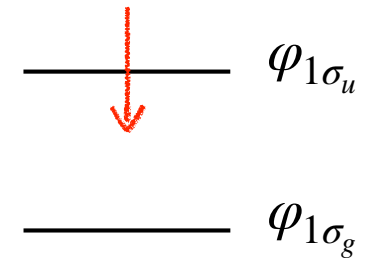
$$|\varphi_1\rangle \equiv \hat{a}_1^\dagger |\text{vac}\rangle$$



$$|\varphi_2\rangle \equiv \hat{a}_2^\dagger |\text{vac}\rangle$$

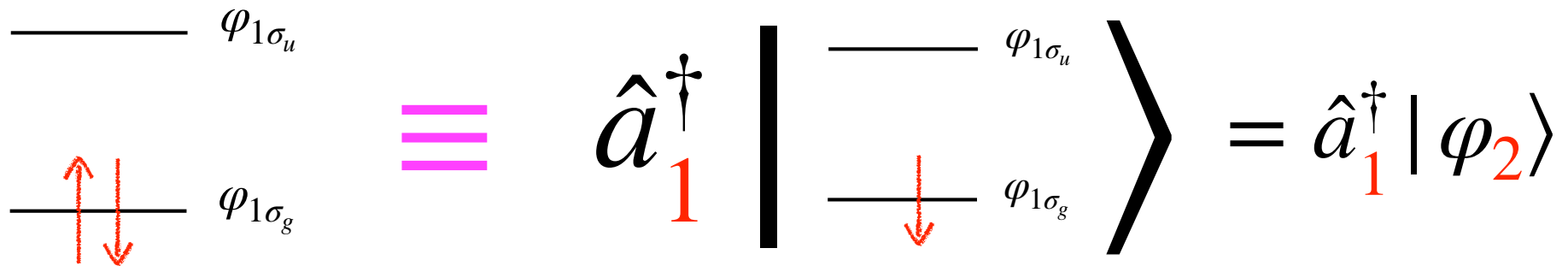


$$|\varphi_3\rangle \equiv \hat{a}_3^\dagger |\text{vac}\rangle$$

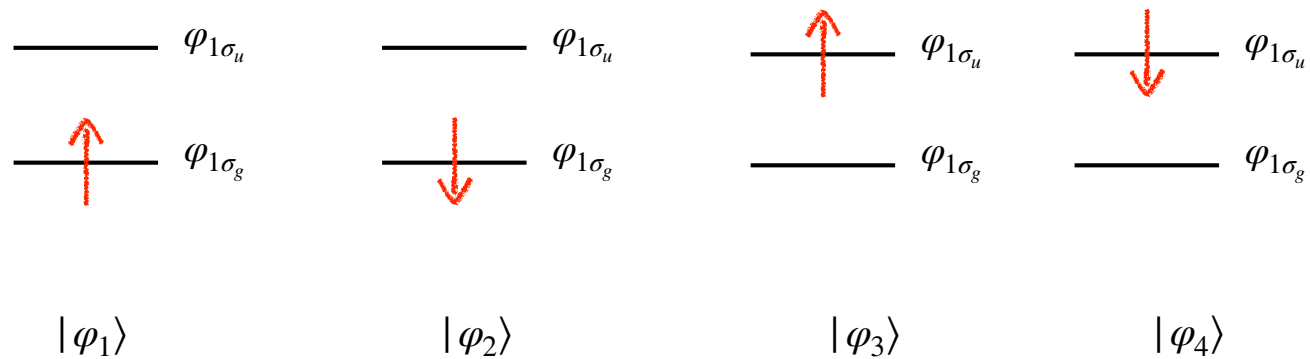


$$|\varphi_4\rangle \equiv \hat{a}_4^\dagger |\text{vac}\rangle$$

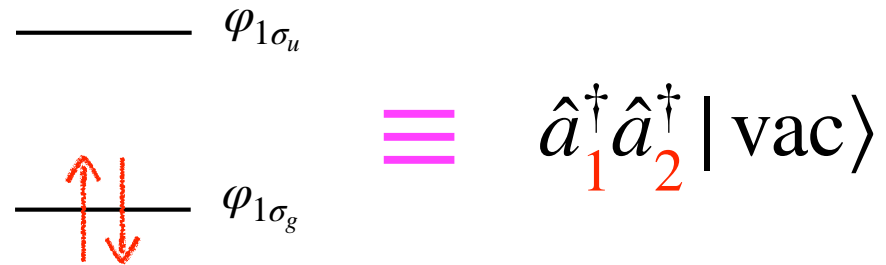
Two-electron quantum states



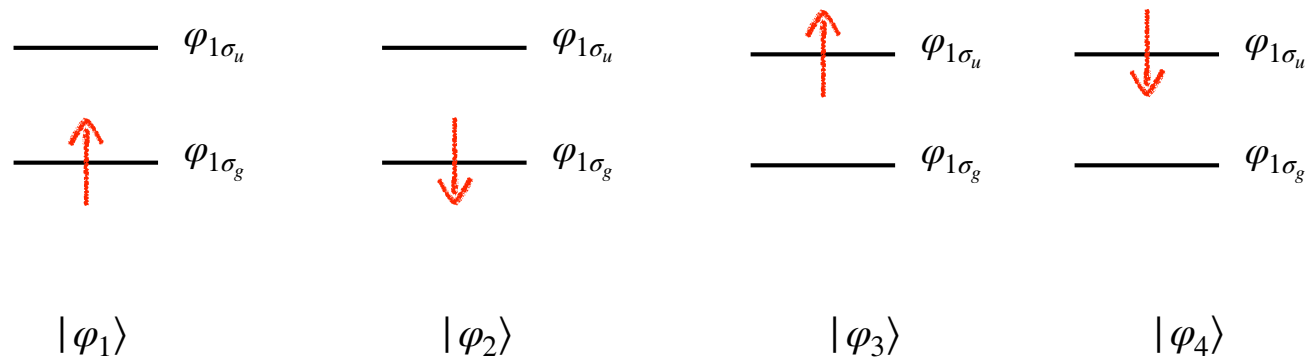
One-electron quantum states



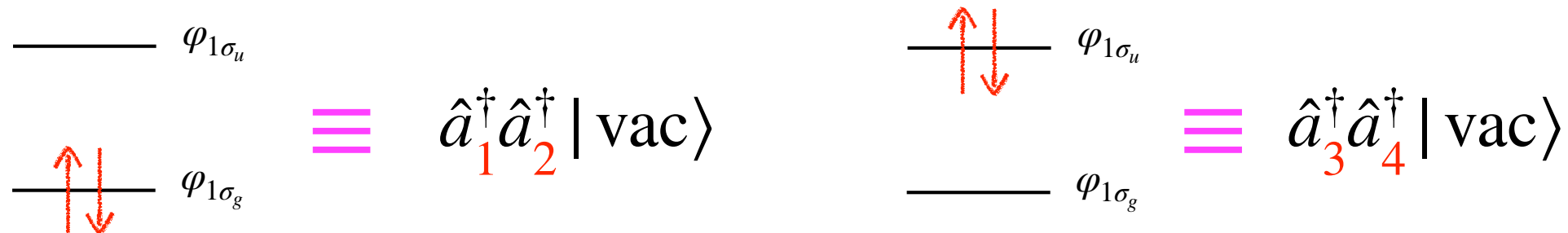
Two-electron quantum states



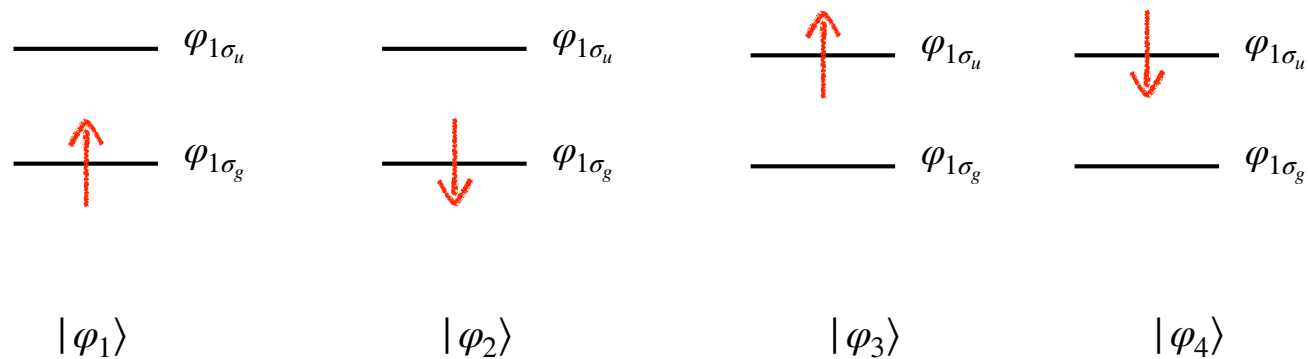
One-electron quantum states



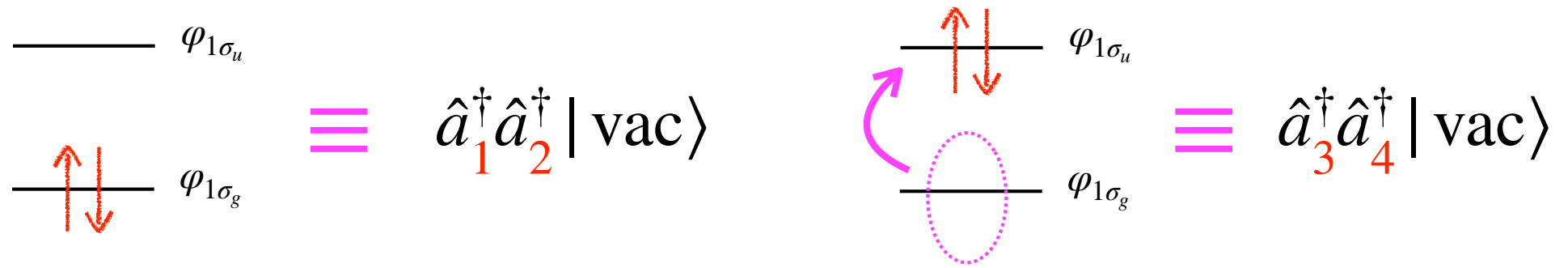
Two-electron quantum states



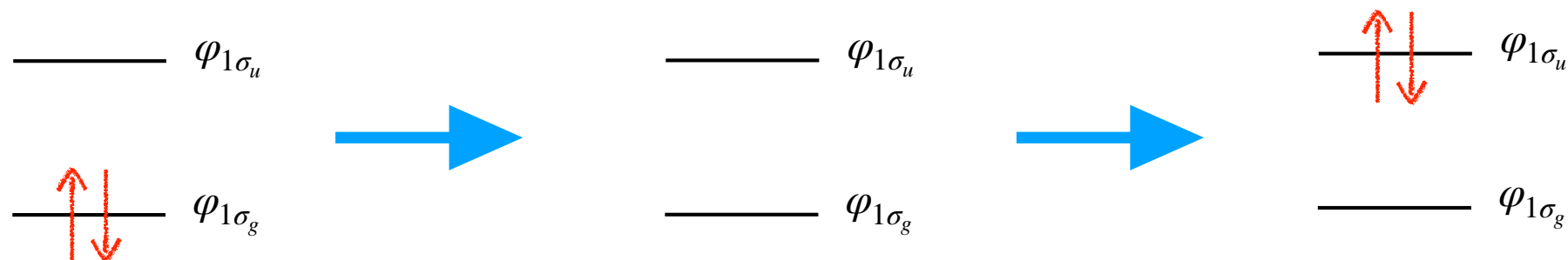
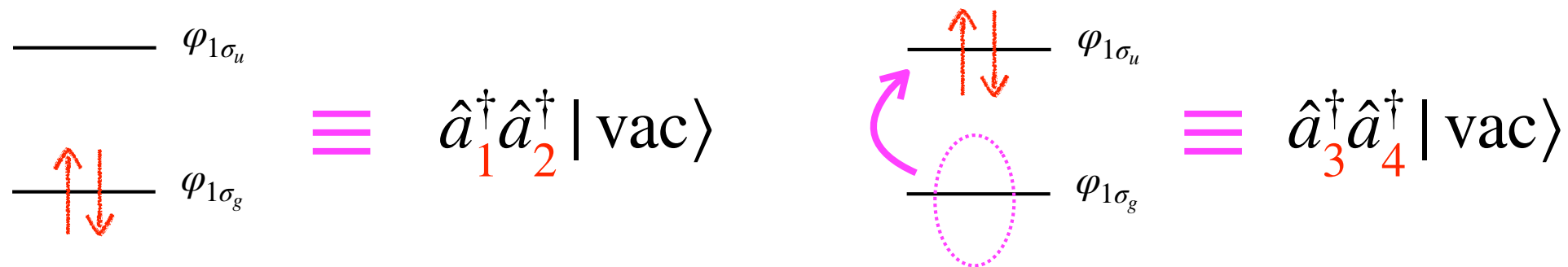
One-electron quantum states



Two-electron quantum states



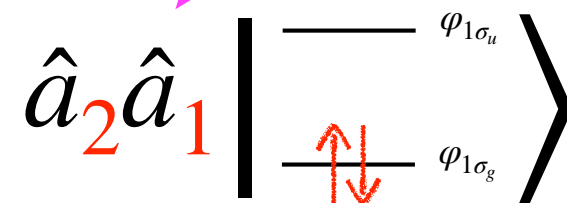
Two-electron quantum states



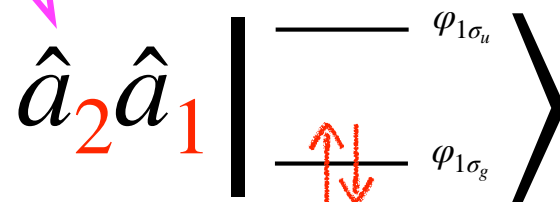
|||

$$\hat{a}_2 \hat{a}_1 \left| \begin{array}{c} \text{---} \varphi_{1\sigma_u} \\ \uparrow\downarrow \text{---} \varphi_{1\sigma_g} \end{array} \right\rangle = |\text{vac}\rangle$$

Quantum operator that **removes**
the electron occupying the spin-orbital φ_1



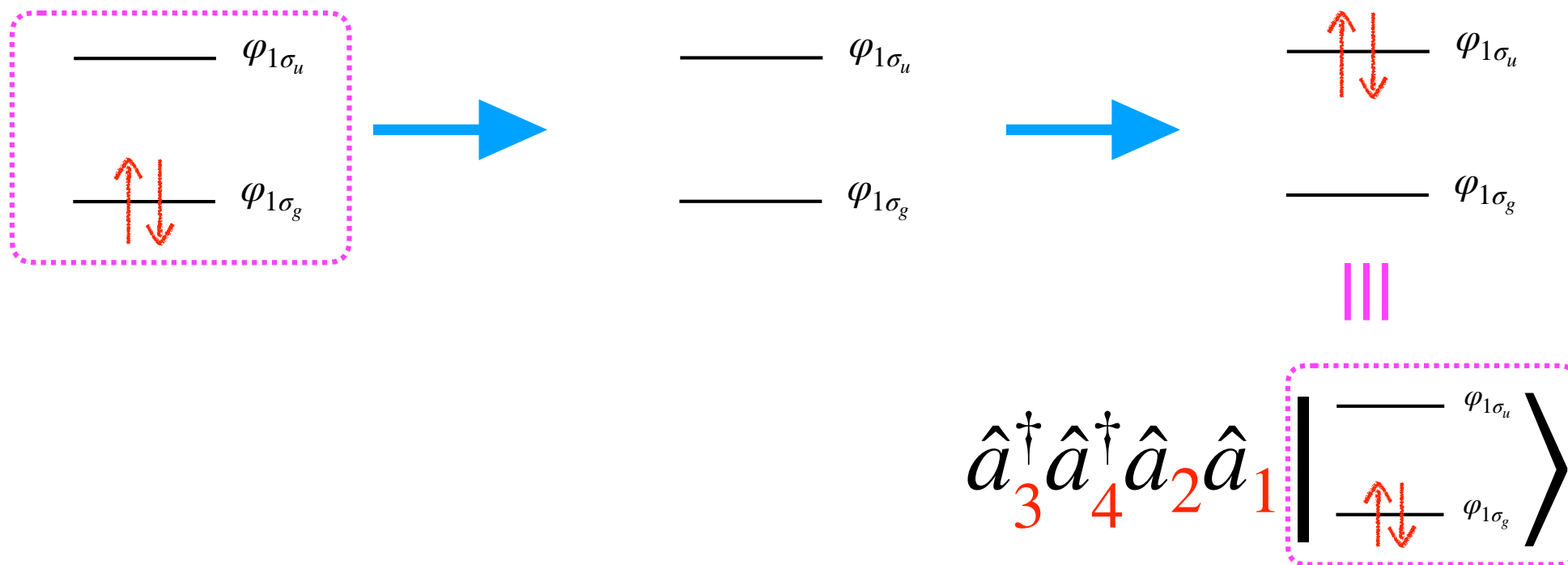
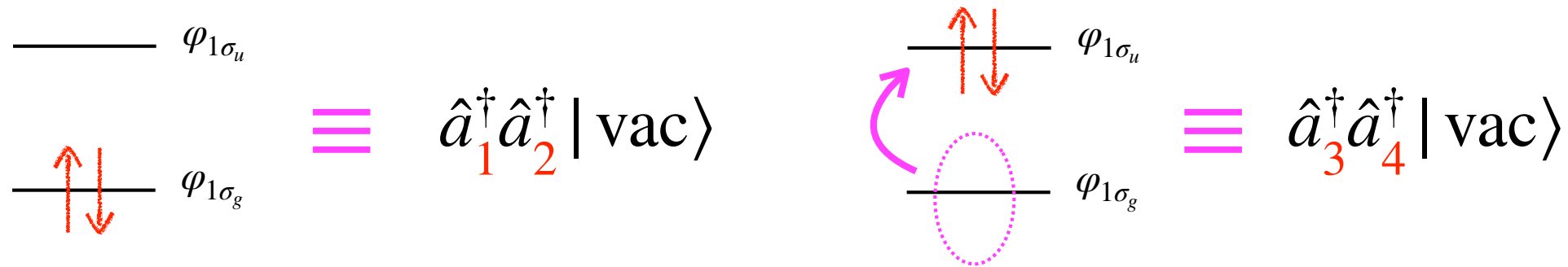
Quantum operator that **removes**
the electron occupying the spin-orbital φ_2



Annihilation operators

$$\hat{a}_2 \hat{a}_1 \left| \begin{array}{c} \text{---} \varphi_{1\sigma_u} \\ \uparrow\downarrow \text{---} \varphi_{1\sigma_g} \end{array} \right\rangle$$

Two-electron quantum states



Double excitation operator

$$\left| \begin{array}{c} \uparrow\downarrow \\ \varphi_{1\sigma_u} \\ \varphi_{1\sigma_g} \end{array} \right\rangle \equiv \hat{a}_3^\dagger \hat{a}_4^\dagger \hat{a}_2 \hat{a}_1 \left| \begin{array}{c} \varphi_{1\sigma_u} \\ \uparrow\downarrow \\ \varphi_{1\sigma_g} \end{array} \right\rangle$$

Schrödinger equation in second quantization

$$\hat{H} |\Psi_0\rangle = E_0 |\Psi_0\rangle$$

Step 1: Choose a *one-electron basis* of molecular spin orbitals $\{\varphi_P\}_{P=1,2,3,\dots,\mathcal{M}}$

Schrödinger equation in second quantization

$$\hat{H} |\Psi_0\rangle = E_0 |\Psi_0\rangle$$

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Usually Hartree-Fock (HF) orbitals

Schrödinger equation in second quantization

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Step 1: Choose a *one-electron basis* of molecular spin orbitals $\{\varphi_P\}_{P=1,2,3,\dots,\mathcal{M}}$



Usually Hartree-Fock (HF) orbitals

$$\{\chi_\nu(\mathbf{x})\}_{\nu=A,B,\dots}$$



*Atomic spin-orbitals centered
on the nuclei A, B, ...*

Schrödinger equation in second quantization

$$\hat{H} |\Psi_0\rangle = E_0 |\Psi_0\rangle$$

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Usually Hartree-Fock (HF) spin-orbitals

$$\{\chi_\nu(\mathbf{x})\}_{\nu=A,B,\dots}$$



Atomic spin-orbitals centered
on the nuclei A, B, \dots

$$\varphi_P(\mathbf{x}) = \sum_{\nu}^{\text{nuclei}} C_{\nu P} \chi_\nu(\mathbf{x})$$



Schrödinger equation in second quantization

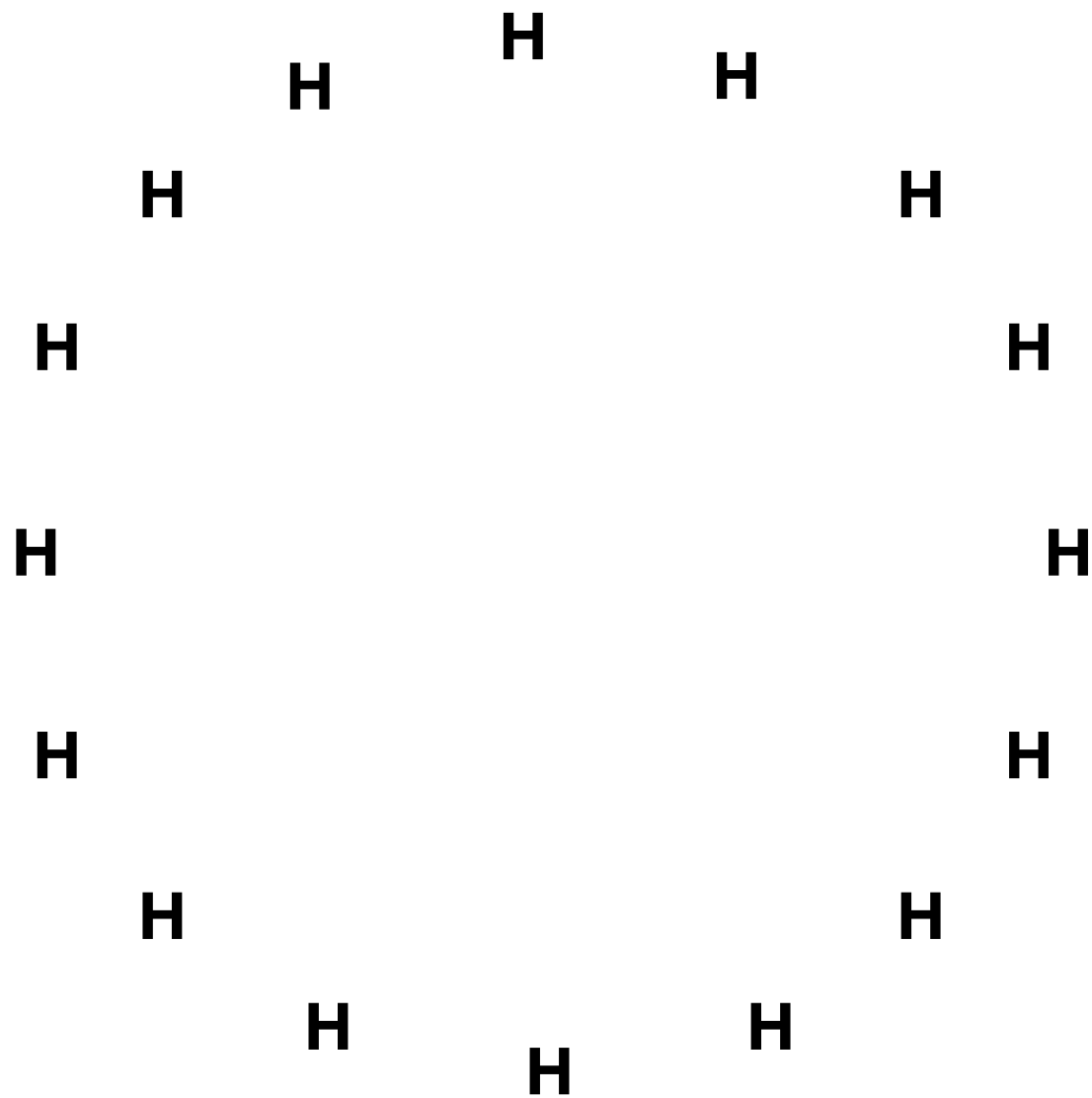
$$\hat{H}|\Psi_0\rangle = E_0|\Psi_0\rangle$$

Step 1: Choose a *one-electron basis* of molecular spin orbitals $\{\varphi_P\}_{P=1,2,3,\dots,\mathcal{M}}$

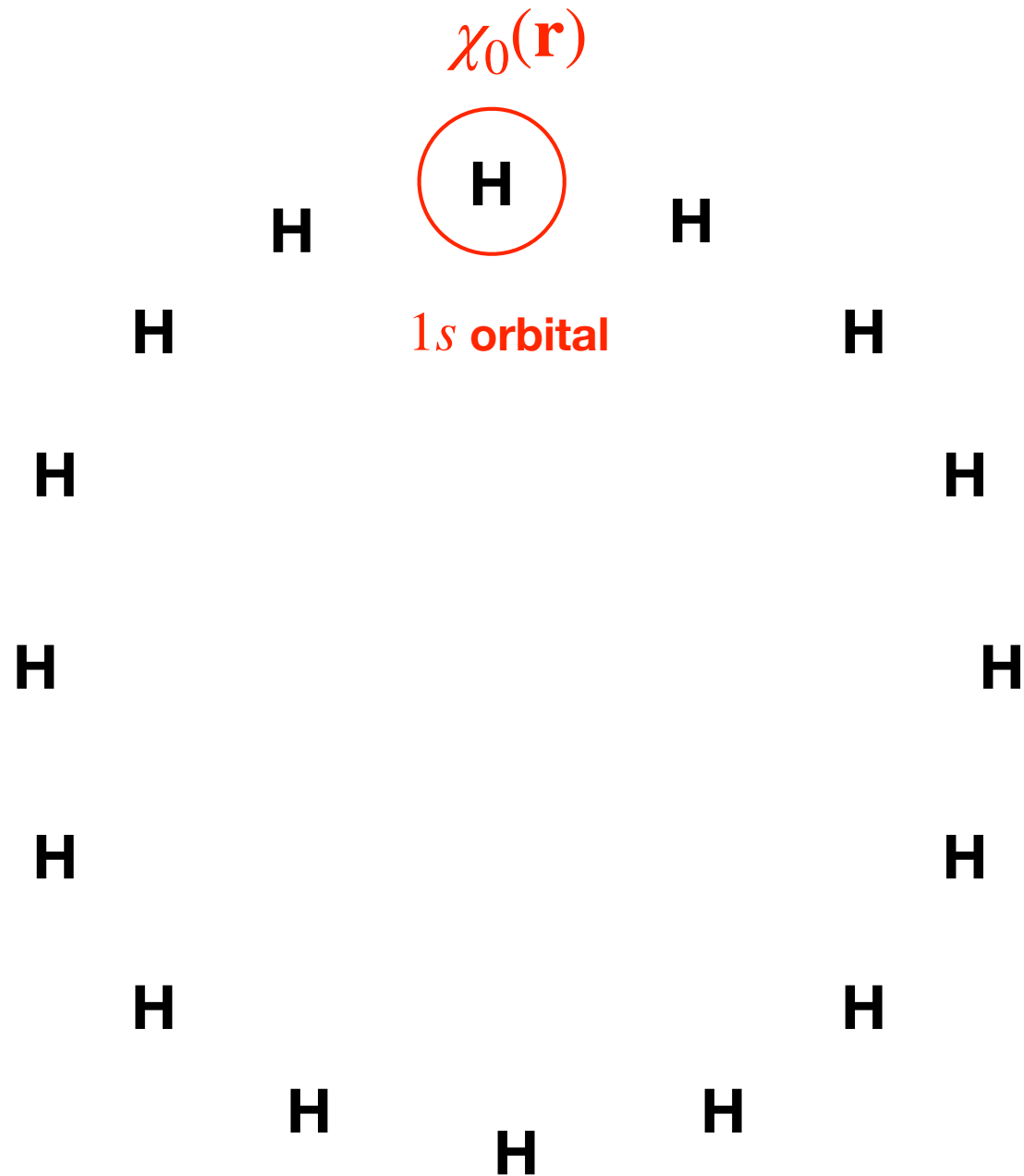
Optimised through
energy minimisation

$$\varphi_P(\mathbf{x}) = \sum_{\nu}^{\text{nuclei}} C_{\nu P} \chi_{\nu}(\mathbf{x})$$

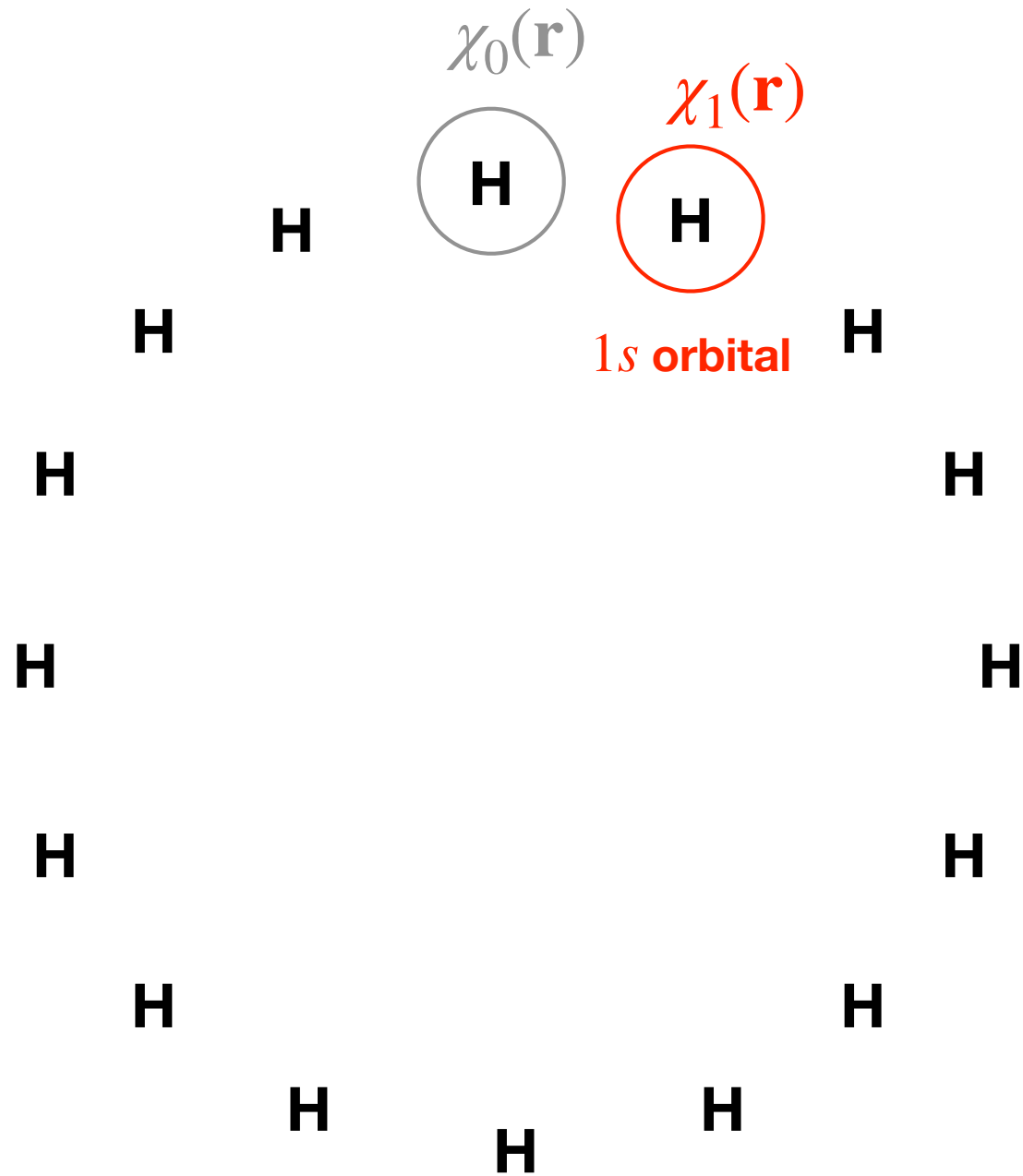
Prototypical ring of $L = 16$ hydrogen atoms



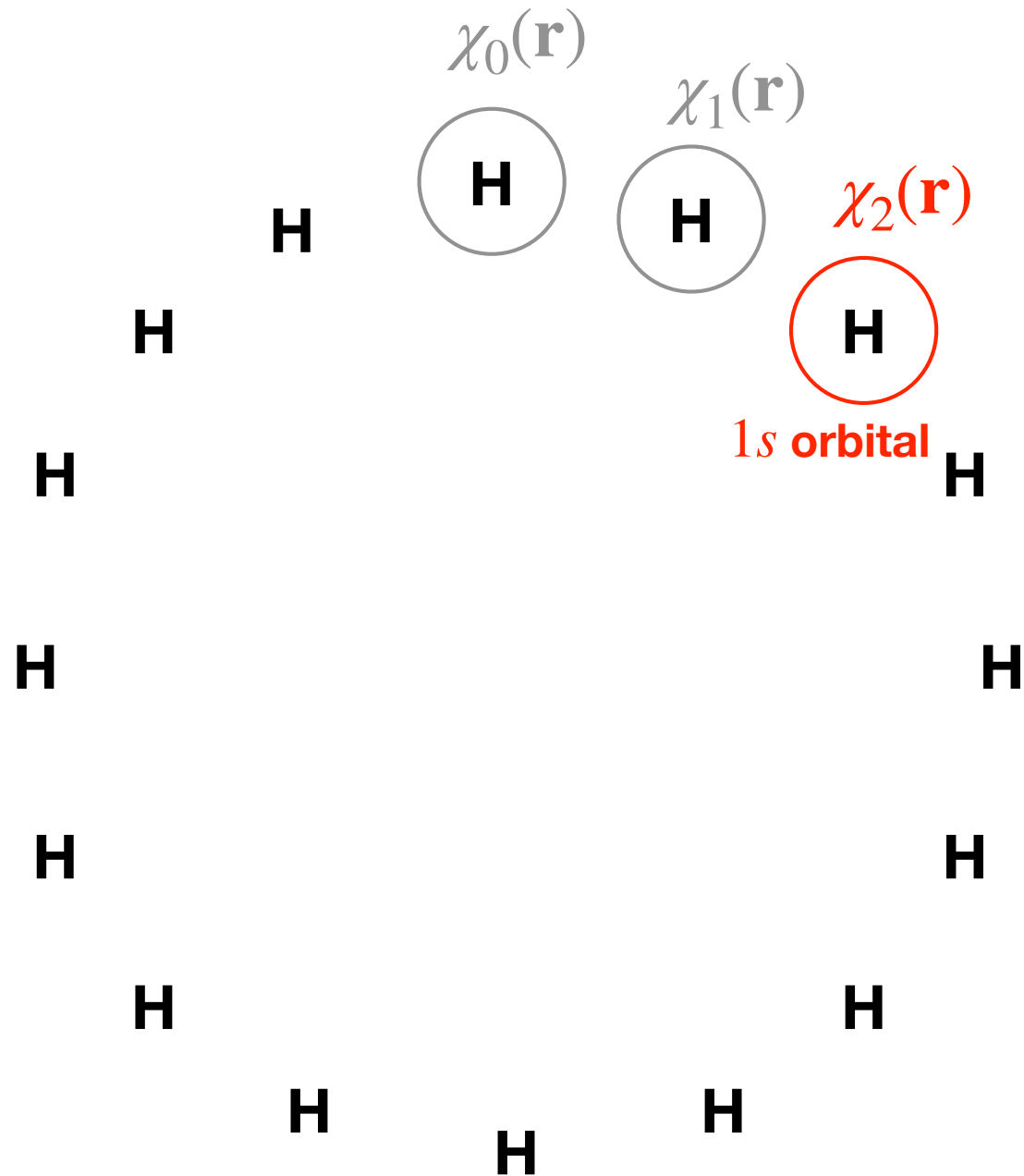
Prototypical ring of $L = 16$ hydrogen atoms



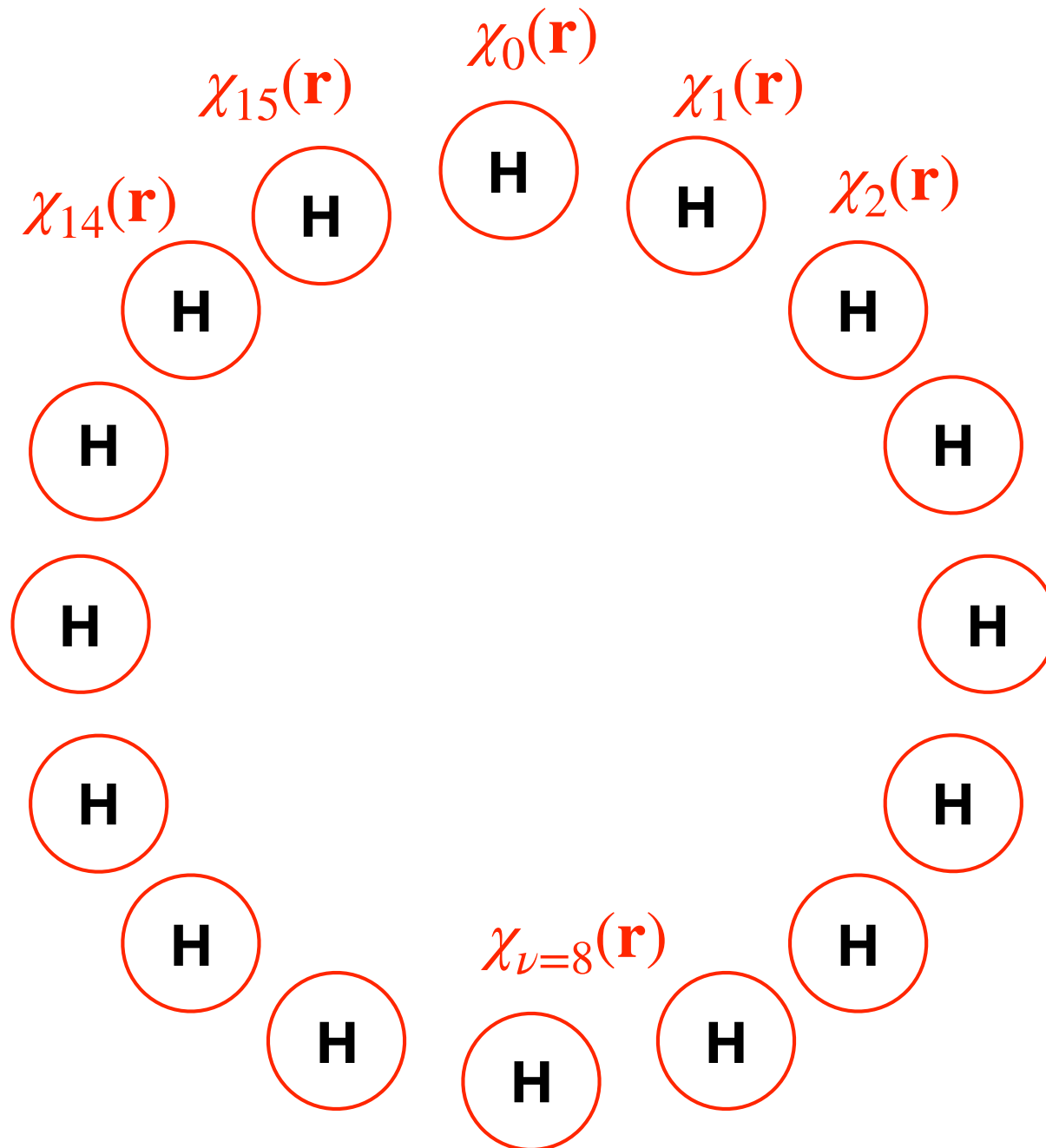
Prototypical ring of $L = 16$ hydrogen atoms



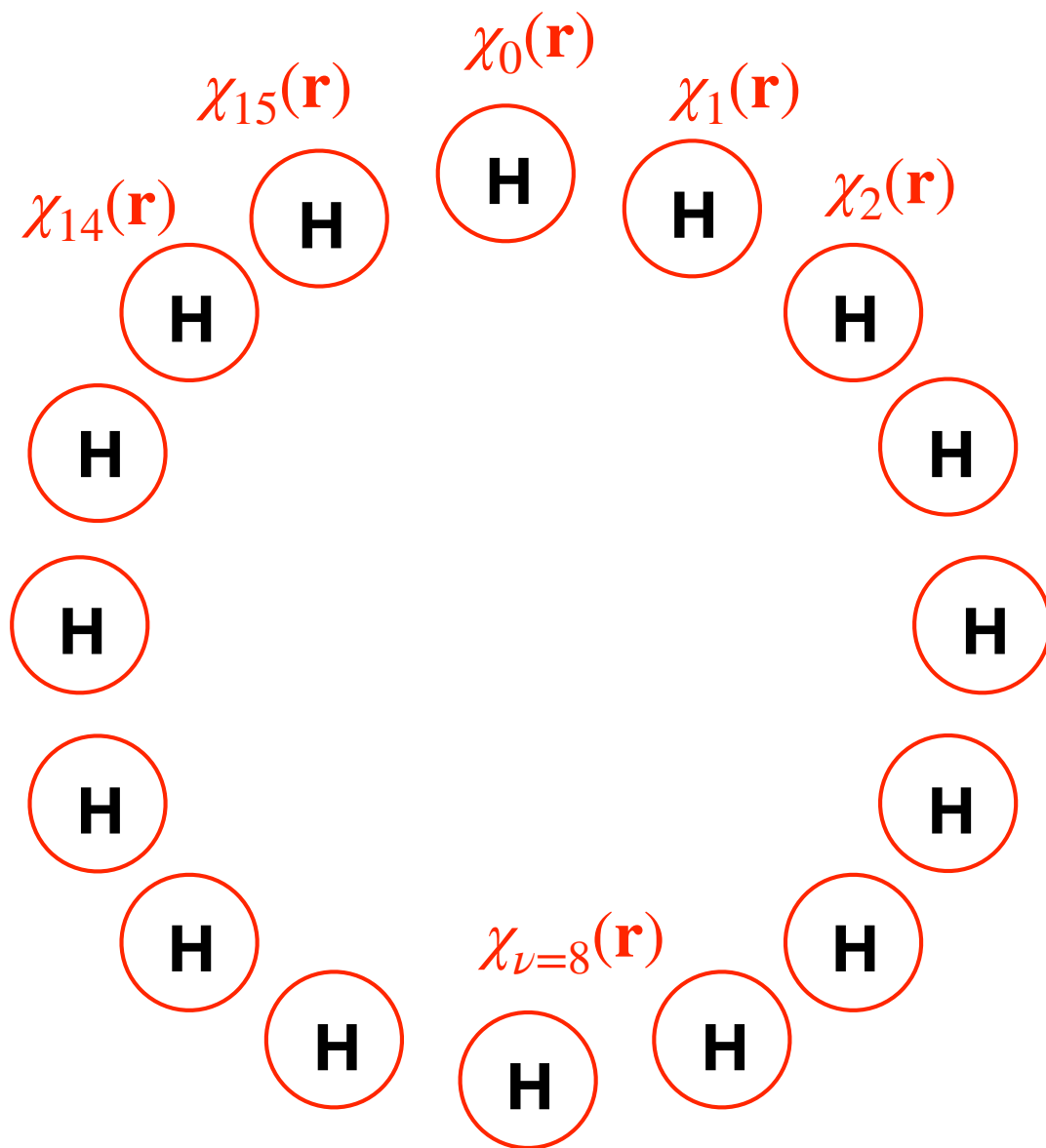
Prototypical ring of $L = 16$ hydrogen atoms



Prototypical ring of $L = 16$ hydrogen atoms



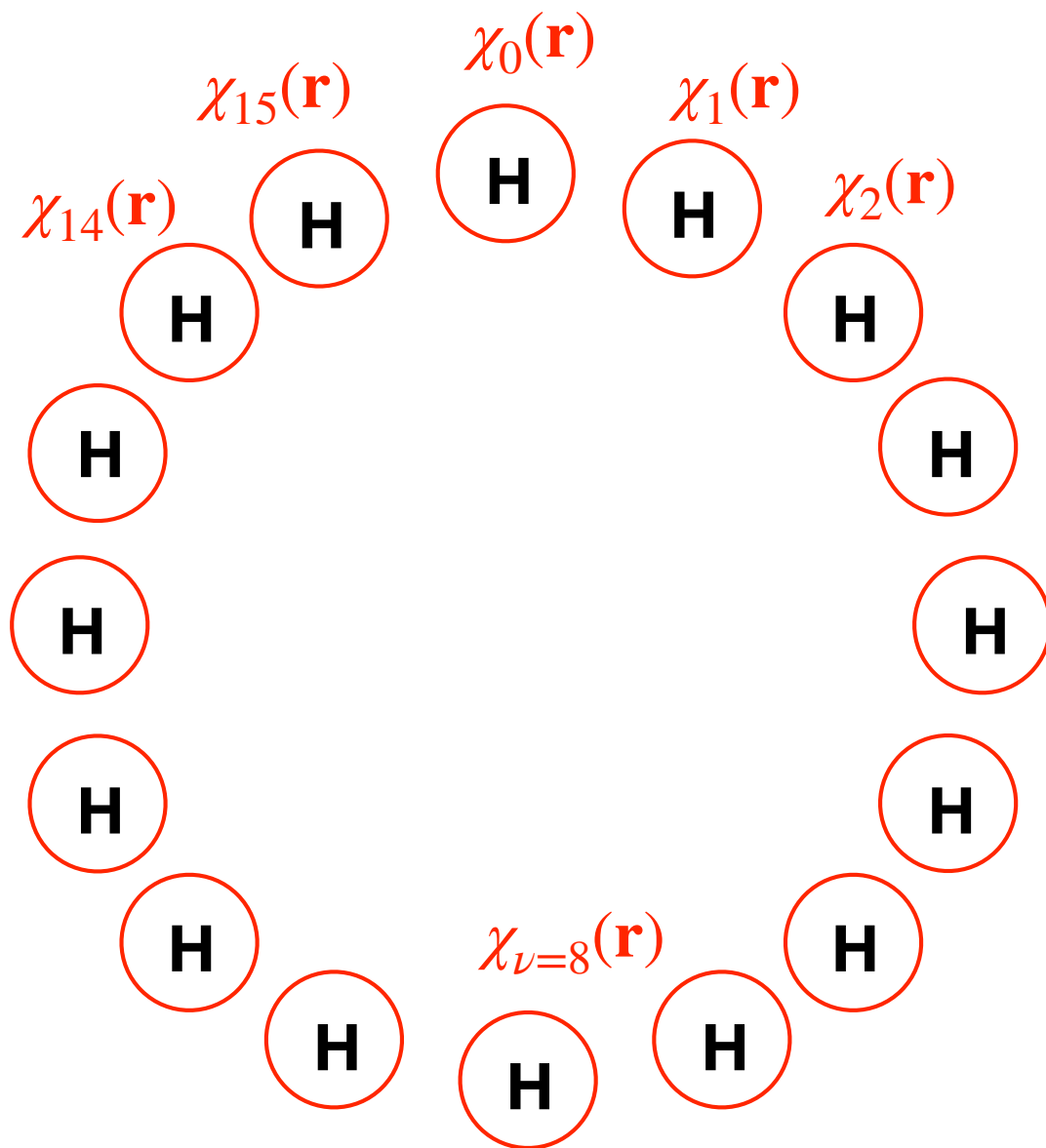
Prototypical ring of $L = 16$ hydrogen atoms



Molecular spin-orbitals

$$\varphi_P(\mathbf{x}) = \sum_{\nu=0}^{15} C_{\nu P} \chi_{\nu}(\mathbf{x})$$

Prototypical ring of $L = 16$ hydrogen atoms



Molecular spin-orbitals

$$\varphi_P(\mathbf{x}) = \sum_{\nu} C_{\nu P} \chi_{\nu}(\mathbf{x})$$

Delocalised over the ring

Schrödinger equation in second quantization

$$\hat{H} |\Psi_0\rangle = E_0 |\Psi_0\rangle$$

Step 1: Choose a *one-electron basis* of molecular spin orbitals $\{\varphi_P\}_{P=1,2,3,\dots,\mathcal{M}}$

Step 2: Implement the Hamiltonian in second quantization *in that basis*

$$\hat{H} \equiv \sum_{PQ} \langle \varphi_P | \hat{h} | \varphi_Q \rangle \hat{a}_P^\dagger \hat{a}_Q + \frac{1}{2} \sum_{PQRS} \langle \varphi_P \varphi_Q | \hat{g} | \varphi_R \varphi_S \rangle \hat{a}_P^\dagger \hat{a}_Q^\dagger \hat{a}_S \hat{a}_R$$

See the video* for further explanations

Schrödinger equation in second quantization

$$\hat{H} |\Psi_0\rangle = E_0 |\Psi_0\rangle$$

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$$\int d\mathbf{x} \varphi_P(\mathbf{x}) \left(-\frac{1}{2} \nabla_{\mathbf{r}}^2 + v_{\text{ext}}(\mathbf{r}) \right) \varphi_Q(\mathbf{x})$$

One-electron integrals

Schrödinger equation in second quantization

$$\hat{H} |\Psi_0\rangle = E_0 |\Psi_0\rangle$$

Step 1: Choose a *one-electron basis* of molecular spin orbitals $\{\varphi_P\}_{P=1,2,3,\dots,\mathcal{M}}$

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Two-electron integrals

$$\int d\mathbf{x}_1 \int d\mathbf{x}_2 \varphi_P(\mathbf{x}_1) \varphi_Q(\mathbf{x}_2) \frac{1}{|\mathbf{r}_1 - \mathbf{r}_2|} \varphi_R(\mathbf{x}_1) \varphi_S(\mathbf{x}_2)$$

Schrödinger equation in second quantization

$$\hat{H} |\Psi_0\rangle = E_0 |\Psi_0\rangle$$

Step 1: Choose a *one-electron basis* of molecular spin orbitals $\{\varphi_P\}_{P=1,2,3,\dots,\mathcal{M}}$

Step 2: Implement the Hamiltonian in second quantization *in that basis*

Step 3: Determine variationally the ground-state wave function expansion *in the corresponding N -electron basis*

$$|\Psi_0\rangle \equiv \sum_{P_1 < P_2 < \dots < P_N} C_{P_1 P_2 \dots P_N} \hat{a}_{P_1}^\dagger \hat{a}_{P_2}^\dagger \dots \hat{a}_{P_{N-1}}^\dagger \hat{a}_{P_N}^\dagger |\text{vac}\rangle$$

Schrödinger equation in second quantization

$$\hat{H} |\Psi_0\rangle = E_0 |\Psi_0\rangle$$

Step 1: Choose a *one-electron basis* of molecular spin orbitals $\{\varphi_P\}_{P=1,2,3,\dots,\mathcal{M}}$

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Distribute N electrons in \mathcal{M} spin-orbitals!

Schrödinger equation in second quantization

$$\hat{H} |\Psi_0\rangle = E_0 |\Psi_0\rangle$$

Step 1: Choose a *one-electron basis* of molecular spin orbitals $\{\varphi_P\}_{P=1,2,3,\dots,\mathcal{M}}$

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$$|\Psi_0\rangle \equiv \sum_{P_1 < P_2 < \dots < P_N} C_{P_1 P_2 \dots P_N} \hat{a}_{P_1}^\dagger \hat{a}_{P_2}^\dagger \dots \hat{a}_{P_{N-1}}^\dagger \hat{a}_{P_N}^\dagger |\text{vac}\rangle$$

Slater determinant

Schrödinger equation in second quantization

$$\hat{H} |\Psi_0\rangle = E_0 |\Psi_0\rangle$$

Step 1: Choose a *one-electron basis* of molecular spin orbitals $\{\varphi_P\}_{P=1,2,3,\dots,\mathcal{M}}$

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Configuration Interaction (CI) coefficient

Configuration interaction (CI) method

$$|\Psi_0\rangle \equiv \sum_{P_1 < P_2 < \dots < P_N} C_{P_1 P_2 \dots P_N} \hat{a}_{P_1}^\dagger \hat{a}_{P_2}^\dagger \dots \hat{a}_{P_{N-1}}^\dagger \hat{a}_{P_N}^\dagger |\text{vac}\rangle$$

CI coefficient

Slater determinant

Configuration interaction (CI) method

$$|\Psi_0\rangle \equiv \sum_{P_1 < P_2 < \dots < P_N} C_{P_1 P_2 \dots P_N} \hat{a}_{P_1}^\dagger \hat{a}_{P_2}^\dagger \dots \hat{a}_{P_{N-1}}^\dagger \hat{a}_{P_N}^\dagger |\text{vac}\rangle$$

CI coefficient

determinants

$$\equiv \sum_{\mathcal{J}} C_{\mathcal{J}} |\text{det}_{\mathcal{J}}\rangle$$

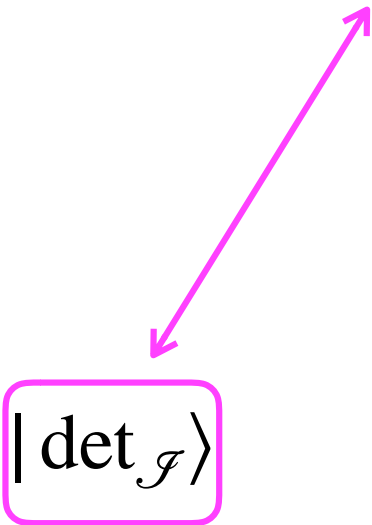
Configuration interaction (CI) method

$$|\Psi_0\rangle \equiv \sum_{P_1 < P_2 < \dots < P_N} C_{P_1 P_2 \dots P_N} \hat{a}_{P_1}^\dagger \hat{a}_{P_2}^\dagger \dots \hat{a}_{P_{N-1}}^\dagger \hat{a}_{P_N}^\dagger |\text{vac}\rangle$$

determinants

$$\equiv \sum_{\mathcal{J}} C_{\mathcal{J}} |\text{det}_{\mathcal{J}}\rangle$$

Known



Configuration interaction (CI) method

$$|\Psi_0\rangle \equiv \sum_{P_1 < P_2 < \dots < P_N} C_{P_1 P_2 \dots P_N} \hat{a}_{P_1}^\dagger \hat{a}_{P_2}^\dagger \dots \hat{a}_{P_{N-1}}^\dagger \hat{a}_{P_N}^\dagger |\text{vac}\rangle$$

$$\equiv \sum_{\mathcal{J}}^{\text{determinants}} C_{\mathcal{J}} |\text{det}_{\mathcal{J}}\rangle$$

Orthonormal basis
of N -electron states

$$\langle \text{det}_{\mathcal{J}} | \text{det}_{\mathcal{L}} \rangle = \delta_{\mathcal{J}\mathcal{L}}$$

Configuration interaction (CI) method

$$|\Psi_0\rangle \equiv \sum_{P_1 < P_2 < \dots < P_N} C_{P_1 P_2 \dots P_N} \hat{a}_{P_1}^\dagger \hat{a}_{P_2}^\dagger \dots \hat{a}_{P_{N-1}}^\dagger \hat{a}_{P_N}^\dagger |\text{vac}\rangle$$

$$\equiv \sum_{\mathcal{J}}^{\text{determinants}} C_{\mathcal{J}} |\text{det}_{\mathcal{J}}\rangle$$

Unknown

Configuration interaction (CI) method

CI energy

$$E_{\text{CI}} = \min_{\{C_{\mathcal{J}}\}} \frac{\langle \Psi(\{C_{\mathcal{J}}\}) | \hat{H} | \Psi(\{C_{\mathcal{J}}\}) \rangle}{\langle \Psi(\{C_{\mathcal{J}}\}) | \Psi(\{C_{\mathcal{J}}\}) \rangle}$$

CI wave function

$$\Psi(\{C_{\mathcal{J}}\}) = \sum_{\mathcal{J}}^{\text{determinants}} C_{\mathcal{J}} | \text{det}_{\mathcal{J}} \rangle$$

Configuration interaction (CI) method

CI energy

$$E_{\text{CI}} = \min_{\{C_{\mathcal{J}}\}} \frac{\langle \Psi(\{C_{\mathcal{J}}\}) | \hat{H} | \Psi(\{C_{\mathcal{J}}\}) \rangle}{\langle \Psi(\{C_{\mathcal{J}}\}) | \Psi(\{C_{\mathcal{J}}\}) \rangle}$$

The minimising CI coefficients are obtained by **diagonalising** the (so-called CI) **Hamiltonian matrix**

$$\mathbf{H} \equiv \left\{ H_{\mathcal{L}\mathcal{J}} = \langle \det_{\mathcal{L}} | \hat{H} | \det_{\mathcal{J}} \rangle \right\}$$

Proof:

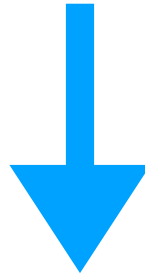
$$E_{\text{CI}} = \min_{\{C_{\mathcal{J}}\}} \frac{\langle \Psi(\{C_{\mathcal{J}}\}) | \hat{H} | \Psi(\{C_{\mathcal{J}}\}) \rangle}{\langle \Psi(\{C_{\mathcal{J}}\}) | \Psi(\{C_{\mathcal{J}}\}) \rangle}$$



$$\frac{\partial}{\partial C_{\mathcal{J}}} \frac{\langle \Psi(\{C_{\mathcal{J}}\}) | \hat{H} | \Psi(\{C_{\mathcal{J}}\}) \rangle}{\langle \Psi(\{C_{\mathcal{J}}\}) | \Psi(\{C_{\mathcal{J}}\}) \rangle} = 0$$

Proof:

$$\frac{\partial}{\partial \mathbf{C}_{\mathcal{J}}} \frac{\langle \Psi(\{\mathbf{C}_{\mathcal{J}}\}) | \hat{H} | \Psi(\{\mathbf{C}_{\mathcal{J}}\}) \rangle}{\langle \Psi(\{\mathbf{C}_{\mathcal{J}}\}) | \Psi(\{\mathbf{C}_{\mathcal{J}}\}) \rangle} = 0$$



$$\begin{aligned} & \frac{\partial}{\partial \mathbf{C}_{\mathcal{J}}} \left[\langle \Psi(\{\mathbf{C}_{\mathcal{J}}\}) | \hat{H} | \Psi(\{\mathbf{C}_{\mathcal{J}}\}) \rangle \right] \\ &= \frac{\partial}{\partial \mathbf{C}_{\mathcal{J}}} \left[\langle \Psi(\{\mathbf{C}_{\mathcal{J}}\}) | \Psi(\{\mathbf{C}_{\mathcal{J}}\}) \rangle \times \frac{\langle \Psi(\{\mathbf{C}_{\mathcal{J}}\}) | \hat{H} | \Psi(\{\mathbf{C}_{\mathcal{J}}\}) \rangle}{\langle \Psi(\{\mathbf{C}_{\mathcal{J}}\}) | \Psi(\{\mathbf{C}_{\mathcal{J}}\}) \rangle} \right] \\ &= E_{\text{CI}} \times \frac{\partial}{\partial \mathbf{C}_{\mathcal{J}}} \left[\langle \Psi(\{\mathbf{C}_{\mathcal{J}}\}) | \Psi(\{\mathbf{C}_{\mathcal{J}}\}) \rangle \right] \end{aligned}$$

Proof:

$$\frac{\partial}{\partial \mathbf{C}_{\mathcal{L}}} \left[\langle \Psi(\{\mathbf{C}_{\mathcal{J}}\}) | \hat{H} | \Psi(\{\mathbf{C}_{\mathcal{J}}\}) \rangle \right] = E_{\text{CI}} \frac{\partial}{\partial \mathbf{C}_{\mathcal{L}}} \left[\langle \Psi(\{\mathbf{C}_{\mathcal{J}}\}) | \Psi(\{\mathbf{C}_{\mathcal{J}}\}) \rangle \right]$$

Proof:

$$\frac{\partial}{\partial \mathbf{C}_{\mathcal{L}}} \left[\langle \Psi(\{\mathbf{C}_{\mathcal{J}}\}) | \hat{H} | \Psi(\{\mathbf{C}_{\mathcal{J}}\}) \rangle \right] = E_{\text{CI}} \frac{\partial}{\partial \mathbf{C}_{\mathcal{L}}} \left[\langle \Psi(\{\mathbf{C}_{\mathcal{J}}\}) | \Psi(\{\mathbf{C}_{\mathcal{J}}\}) \rangle \right]$$



$$= 2 \left\langle \frac{\partial \Psi(\{\mathbf{C}_{\mathcal{J}}\})}{\partial \mathbf{C}_{\mathcal{L}}} \middle| \hat{H} \middle| \Psi(\{\mathbf{C}_{\mathcal{J}}\}) \right\rangle$$

$$= 2 \langle \text{det}_{\mathcal{L}} | \hat{H} | \Psi(\{\mathbf{C}_{\mathcal{J}}\}) \rangle$$

Proof:

$$2\langle \text{det}_{\mathcal{L}} | \hat{H} | \Psi(\{C_{\mathcal{J}}\}) \rangle = E_{\text{CI}} \frac{\partial}{\partial C_{\mathcal{L}}} [\langle \Psi(\{C_{\mathcal{J}}\}) | \Psi(\{C_{\mathcal{J}}\}) \rangle]$$



$$2E_{\text{CI}} \times \langle \text{det}_{\mathcal{L}} | \Psi(\{C_{\mathcal{J}}\}) \rangle$$

Proof:

$$\langle \det_{\mathcal{L}} | \hat{H} | \Psi(\{C_{\mathcal{J}}\}) \rangle = E_{\text{CI}} \times \langle \det_{\mathcal{L}} | \Psi(\{C_{\mathcal{J}}\}) \rangle$$

$$\Psi(\{C_{\mathcal{J}}\}) = \sum_{\mathcal{J}}^{\text{determinants}} C_{\mathcal{J}} | \det_{\mathcal{J}} \rangle$$

$$\begin{aligned} & \sum_{\mathcal{J}} C_{\mathcal{J}} \langle \det_{\mathcal{L}} | \hat{H} | \det_{\mathcal{J}} \rangle \\ &= \sum_{\mathcal{J}} H_{\mathcal{L}\mathcal{J}} C_{\mathcal{J}} \end{aligned}$$

Proof:

$$\sum_{\mathcal{J}} H_{\mathcal{L}\mathcal{J}} C_{\mathcal{J}} = E_{\text{CI}} \times \langle \text{det}_{\mathcal{L}} | \Psi(\{C_{\mathcal{J}}\}) \rangle$$

$$\Psi(\{C_{\mathcal{J}}\}) = \sum_{\mathcal{J}}^{\text{determinants}} C_{\mathcal{J}} | \text{det}_{\mathcal{J}} \rangle$$

$C_{\mathcal{L}}$

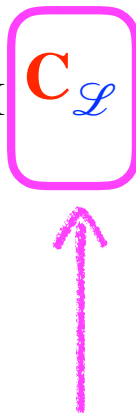
Proof:

$$\sum_{\mathcal{J}} H_{\mathcal{L}\mathcal{J}} C_{\mathcal{J}} = E_{\text{CI}} C_{\mathcal{L}}$$

$$\mathbf{C} = \begin{bmatrix} C_1 \\ C_2 \\ \vdots \\ C_{\mathcal{J}} \\ \vdots \end{bmatrix}$$

$$[\mathbf{HC}]_{\mathcal{L}}$$

Proof:

$$[\mathbf{HC}]_{\mathcal{L}} = E_{\text{CI}} \mathbf{C}_{\mathcal{L}}$$

$$\mathbf{C} = \begin{bmatrix} C_1 \\ C_2 \\ \vdots \\ C_{\mathcal{J}} \\ \vdots \end{bmatrix}$$

Proof:

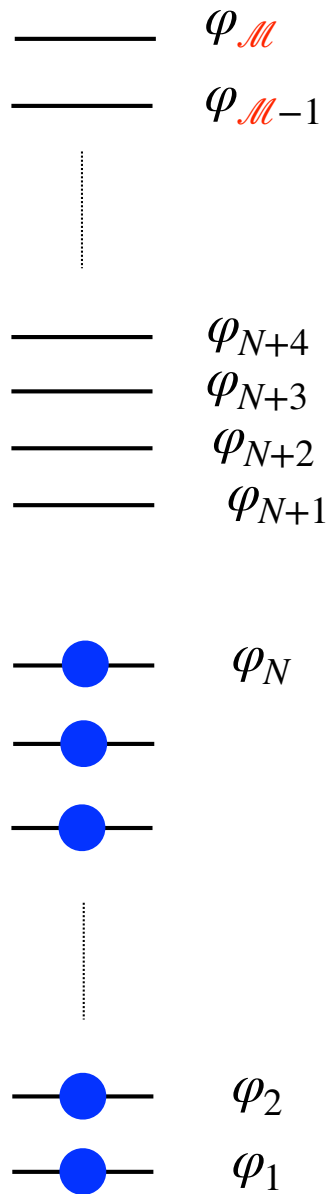
$$[\mathbf{HC}]_{\mathcal{L}} = E_{\text{CI}} \mathbf{C}_{\mathcal{L}}$$

$$\mathbf{C} = \begin{bmatrix} C_1 \\ C_2 \\ \vdots \\ C_{\mathcal{J}} \\ \vdots \end{bmatrix} \quad \Leftrightarrow$$

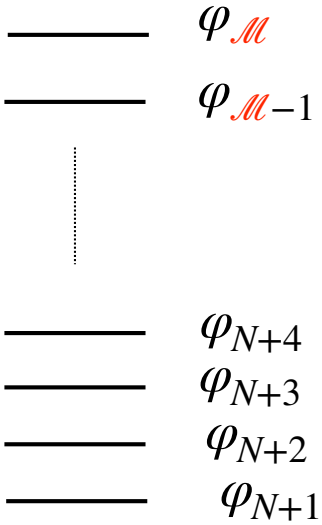
$$\mathbf{HC} = E_{\text{CI}} \mathbf{C}$$



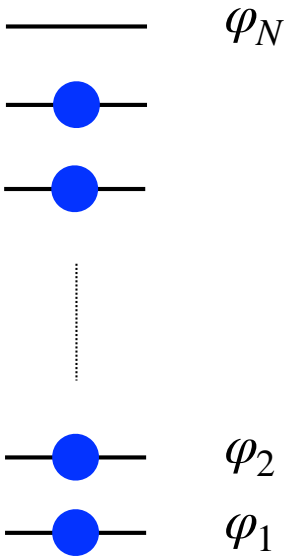
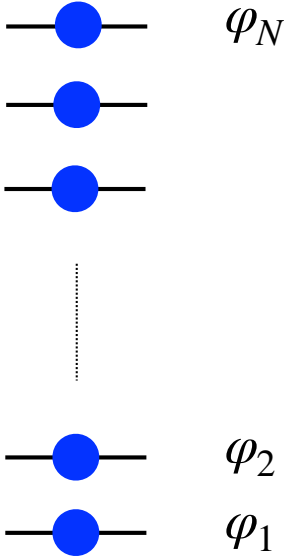
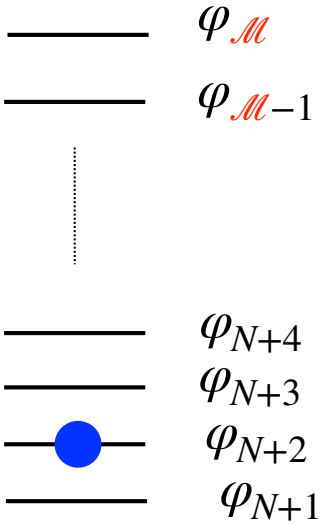
How many determinants in total?



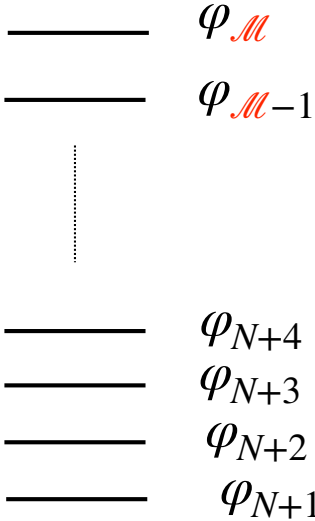
How many determinants in total?



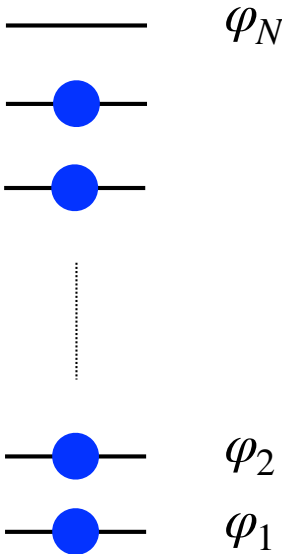
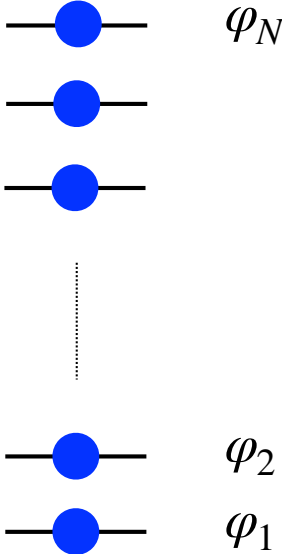
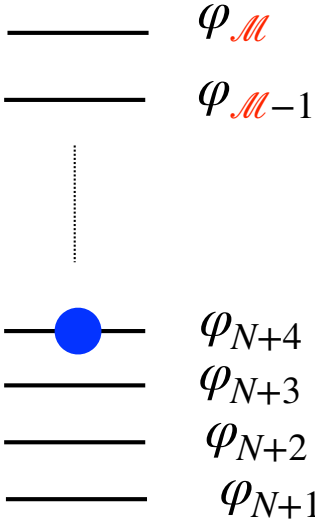
or



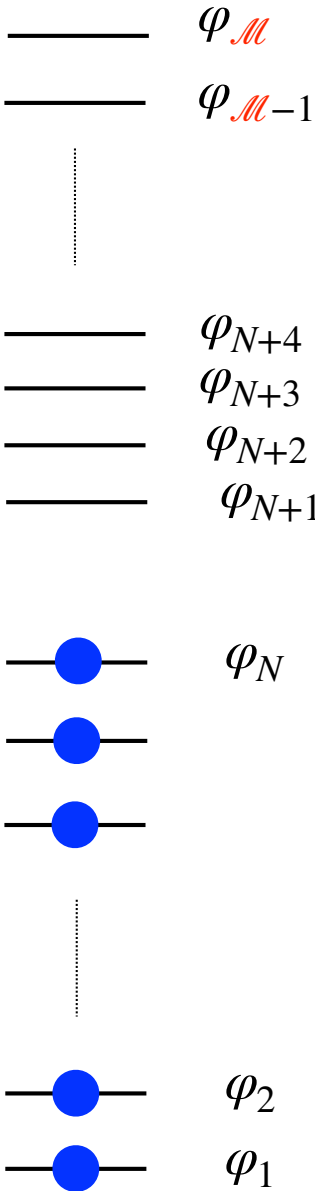
How many determinants in total?



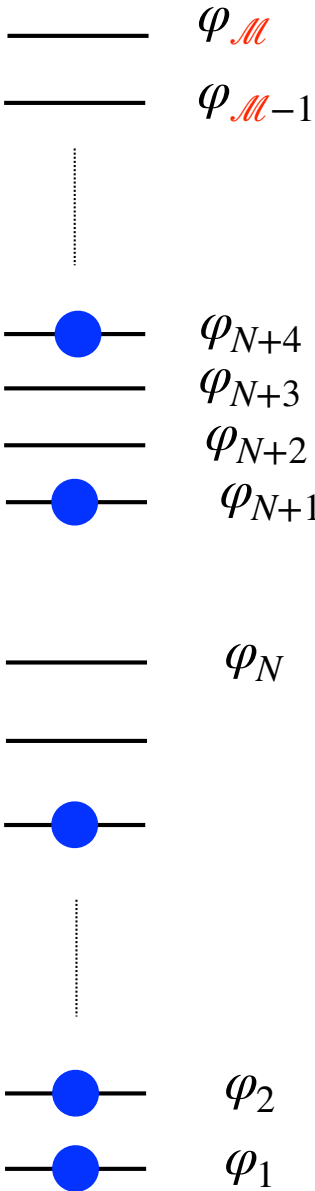
or



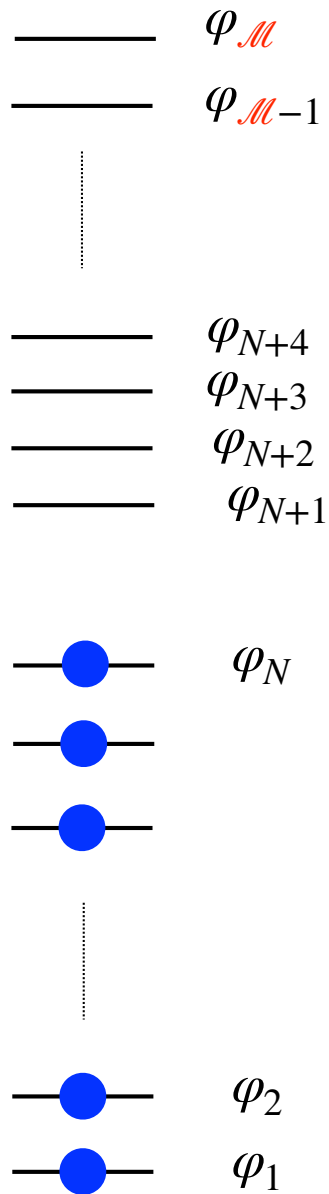
How many determinants in total?



or



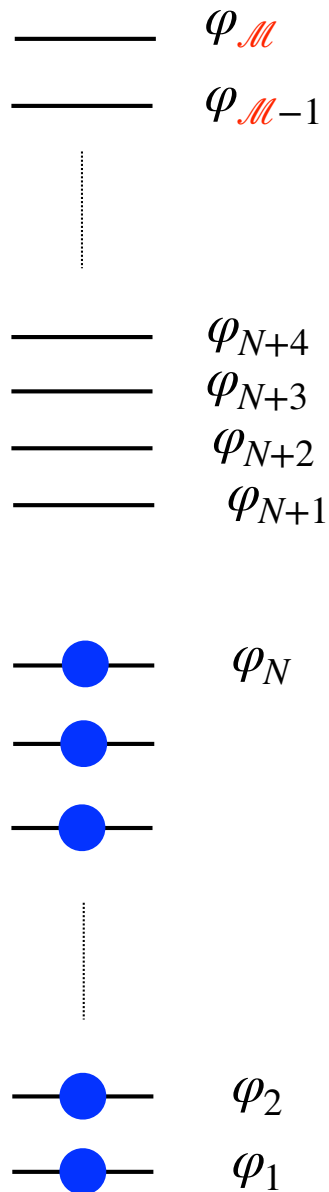
How many determinants in total?



or

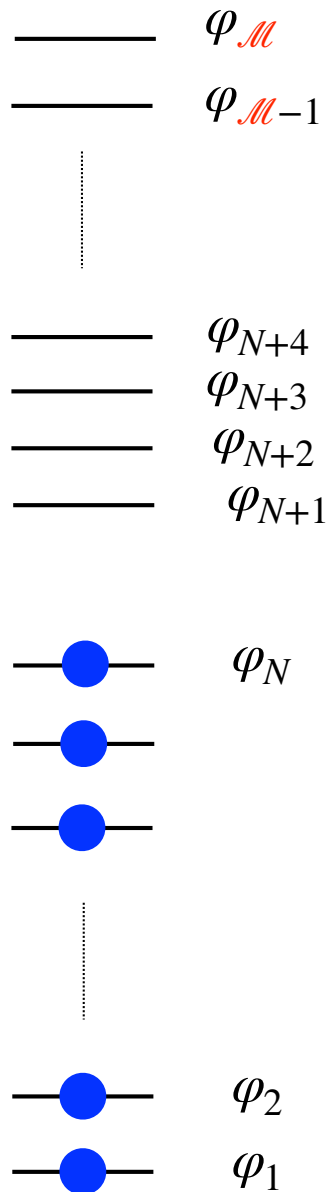
Etc...

How many determinants in total?



We have M spin-orbitals available for N electrons

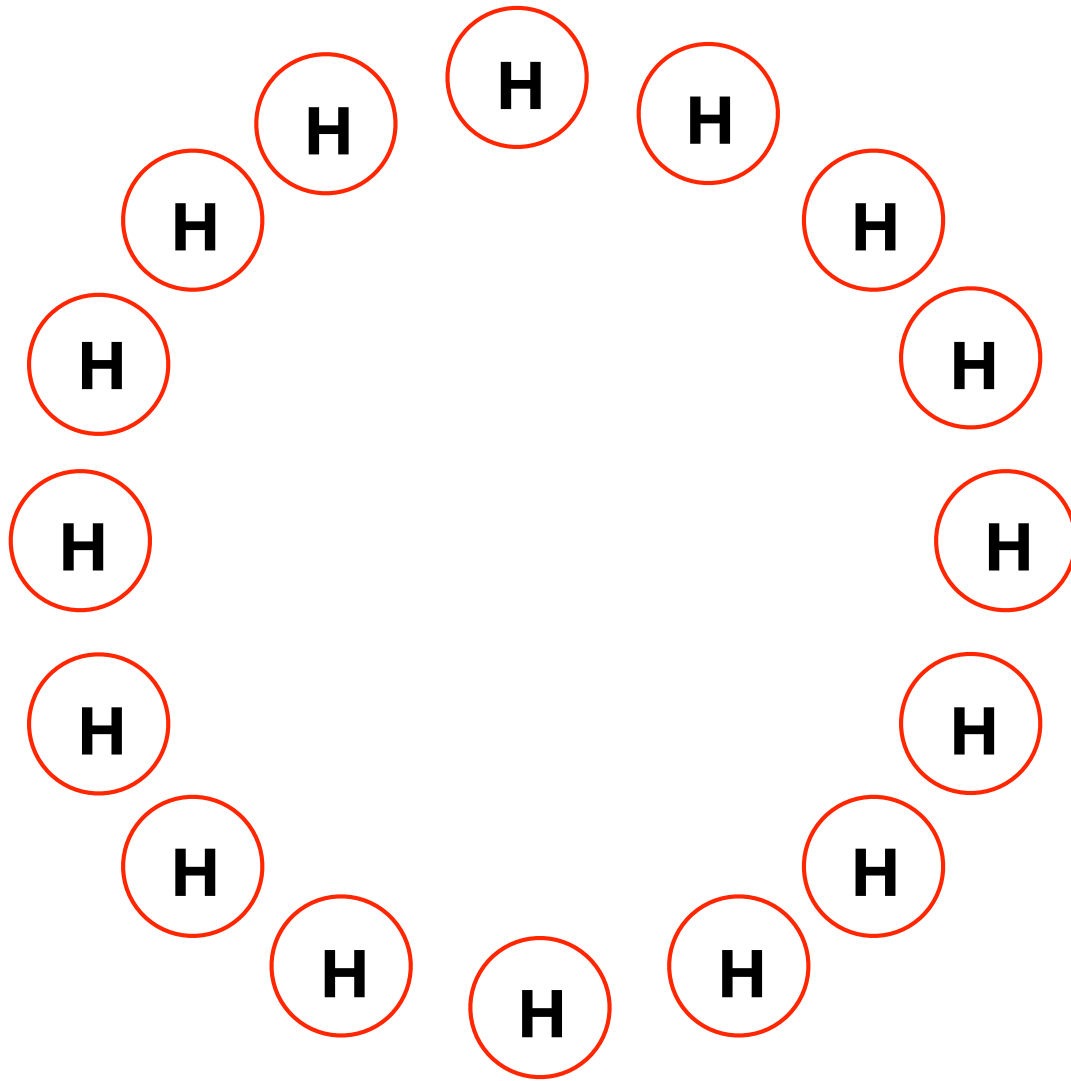
How many determinants in total?



We have M spin-orbitals available for N electrons

$$N_{\text{det.}} = \binom{M}{N} = \frac{M!}{N!(M-N)!}$$

How many determinants in total?



$$\mathcal{M} = 2 \times N$$

↑
Spin

How many determinants in total?

$$\mathcal{M} = 2 \times N$$

$$N_{\text{det.}} = \frac{\mathcal{M}!}{N!(\mathcal{M}-N)!} = \frac{(2N)!}{(N!)^2}$$

How many determinants in total?

$$\mathcal{M} = 2 \times N$$

$$N_{\text{det.}} = \frac{\mathcal{M}!}{N!(\mathcal{M}-N)!} = \frac{(2N)!}{(N!)^2}$$

$$N! \approx \sqrt{2\pi N} \left(\frac{N}{e}\right)^N$$

Stirling formula for large N values

$$\approx \frac{2^{2N}}{\sqrt{\pi N}} = \frac{e^{2N \ln 2}}{\sqrt{\pi N}}$$

How many determinants in total?

$$N_{\text{det.}} \approx \frac{e^{2N \ln 2}}{\sqrt{\pi N}}$$

“Exponential wall”

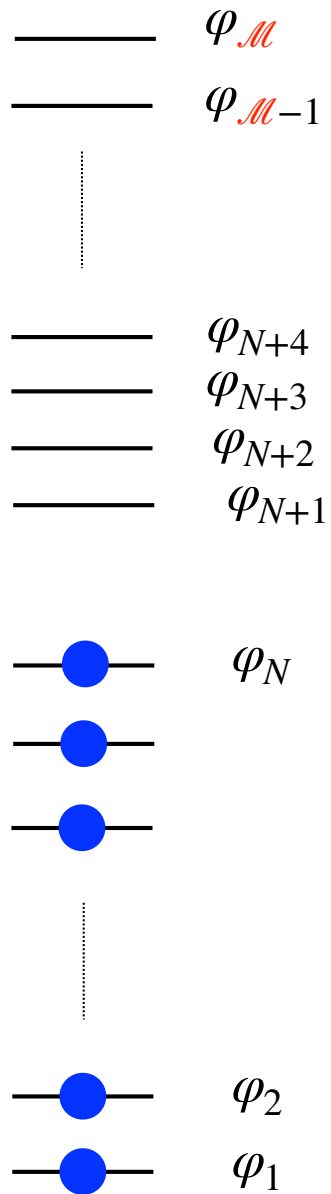
How many determinants in total?

$$N_{\text{det.}} \approx \frac{e^{2N \ln 2}}{\sqrt{\pi N}} \quad \stackrel{N=50}{\approx} 10^{29}$$

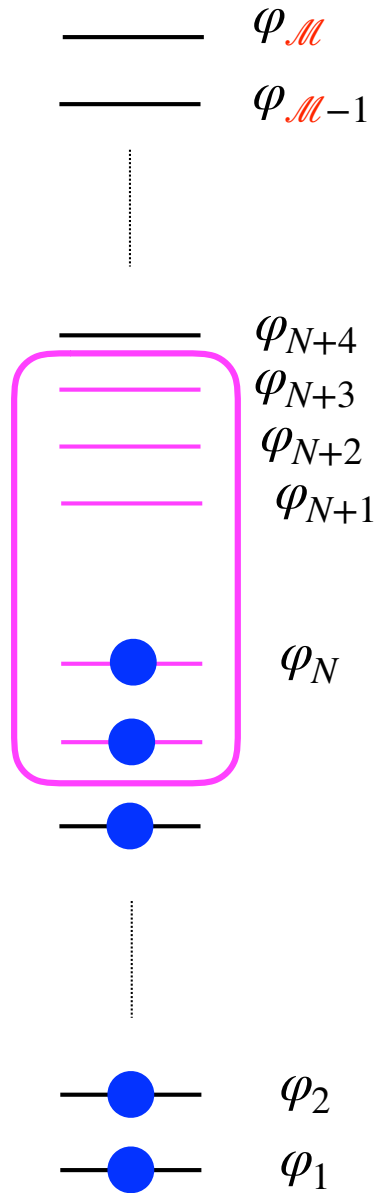
How many determinants in total?

$$N_{\text{det.}} \approx \frac{e^{2N \ln 2}}{\sqrt{\pi N}} \quad \stackrel{N=400}{\approx} 1.88 \times 10^{239}$$

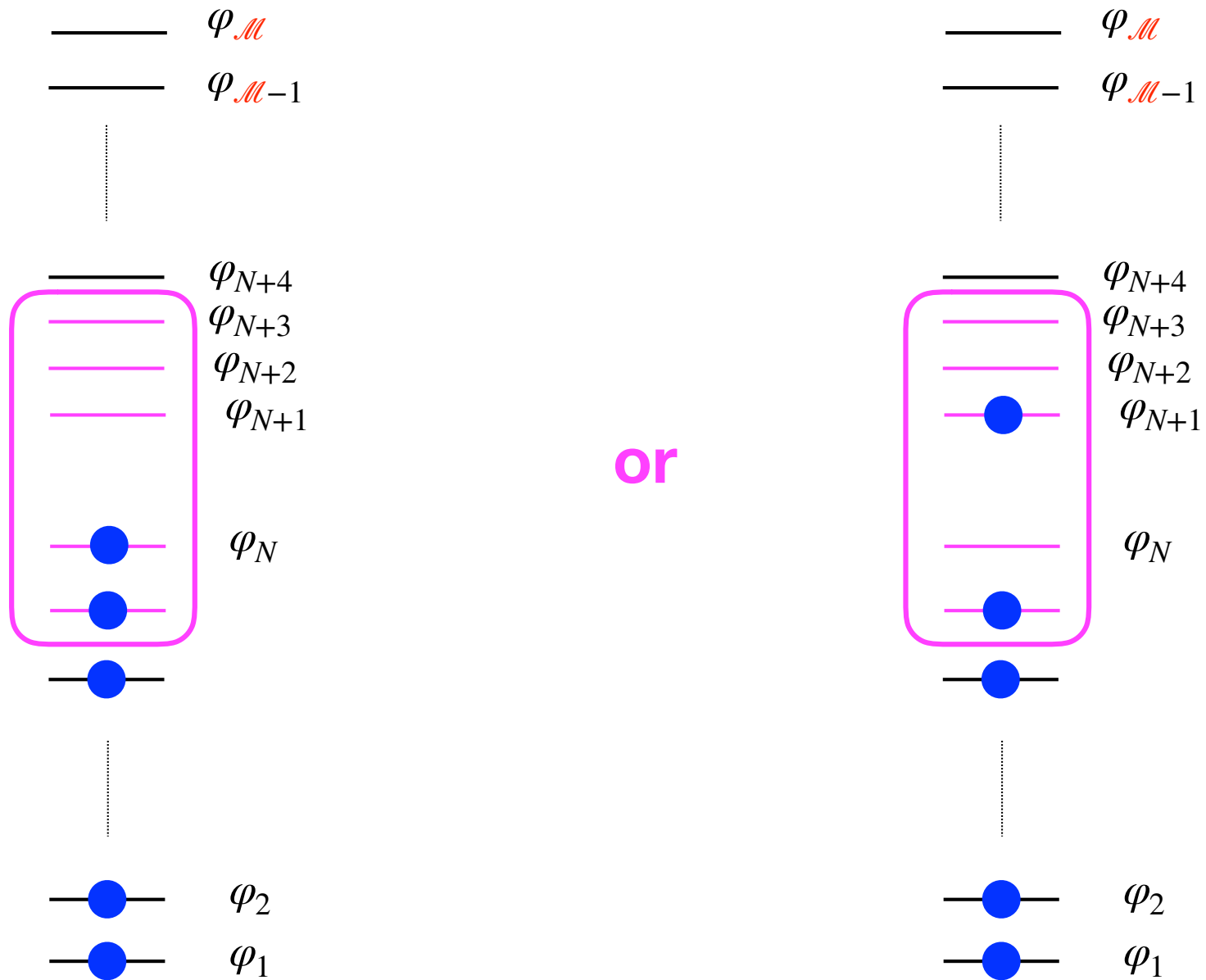
Complete active space CI (CAS-CI) method



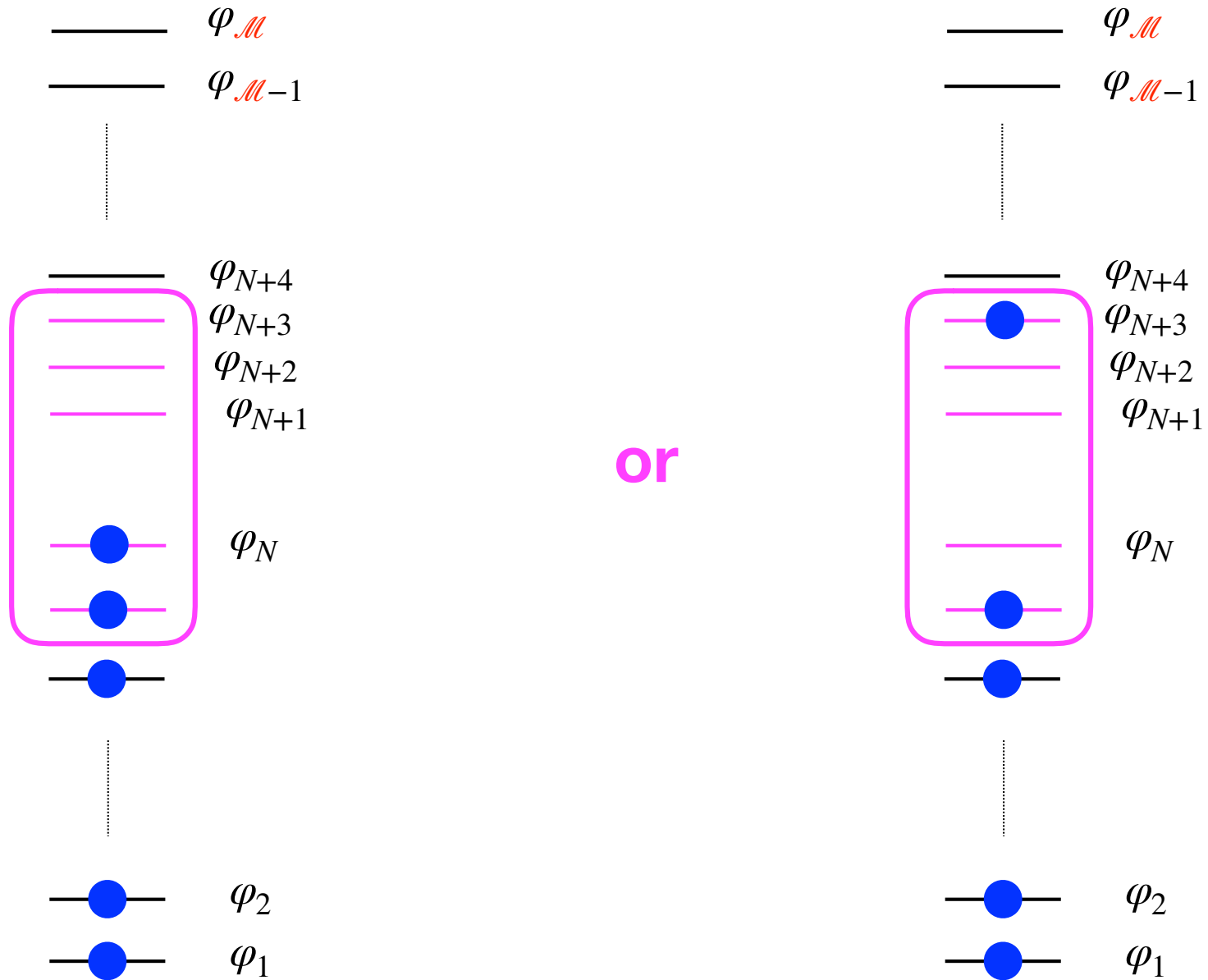
Complete active space CI (CAS-CI) method



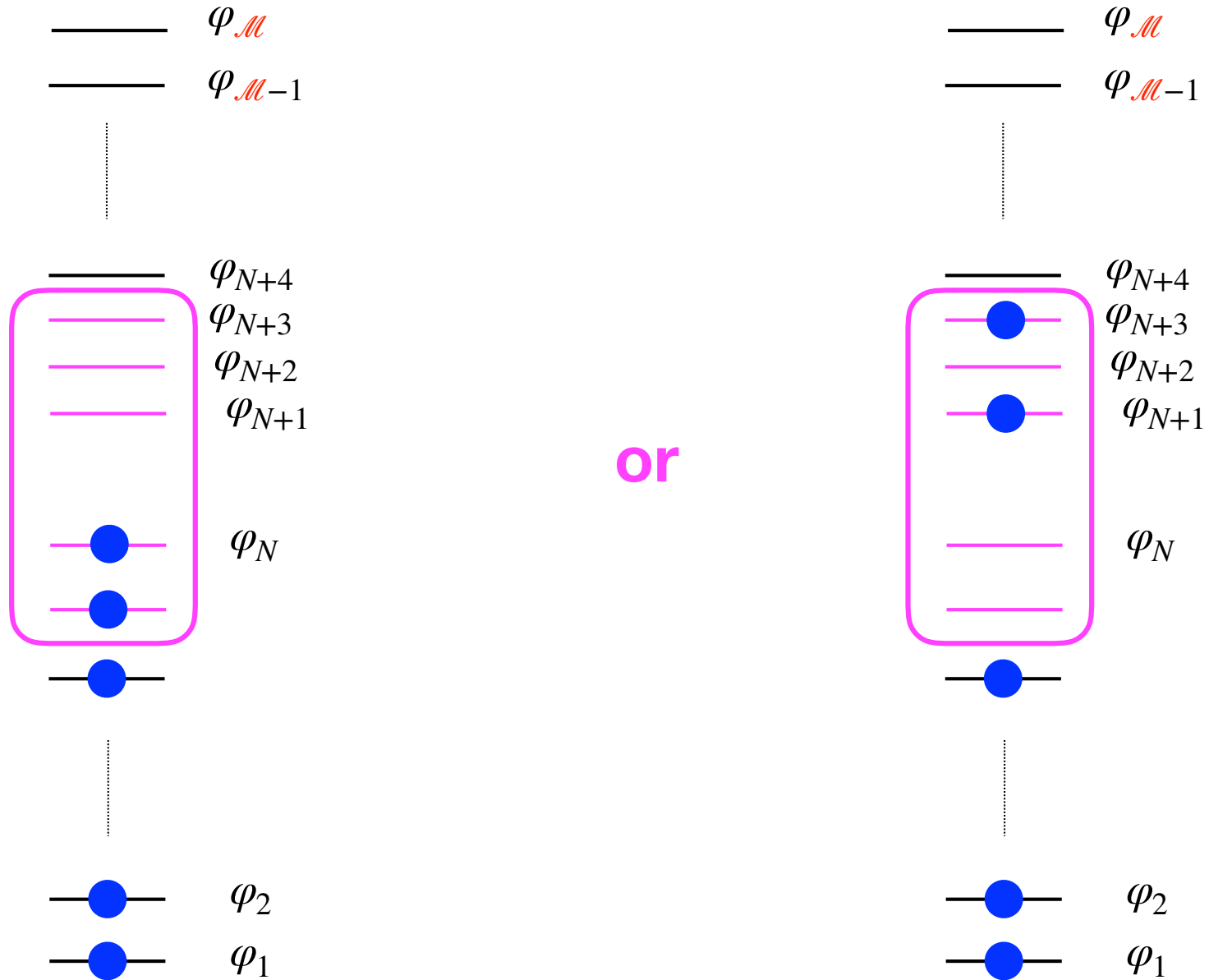
Complete active space CI (CAS-CI) method



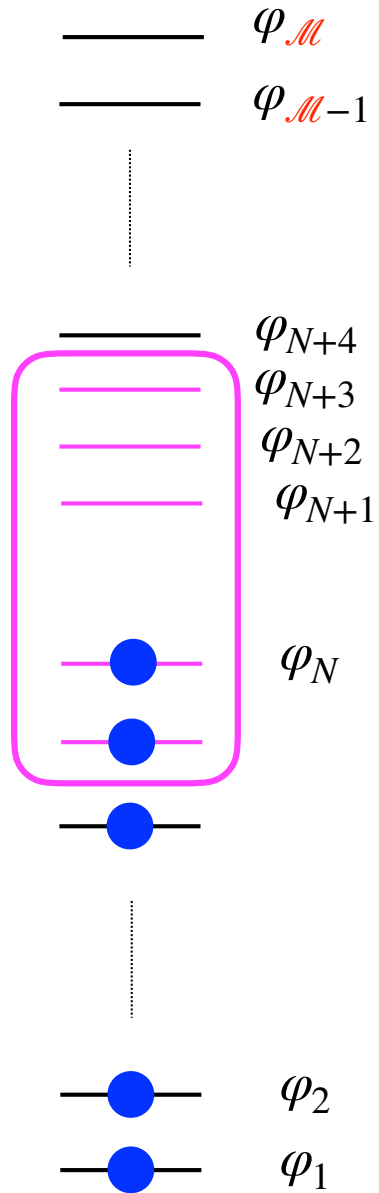
Complete active space CI (CAS-CI) method



Complete active space CI (CAS-CI) method



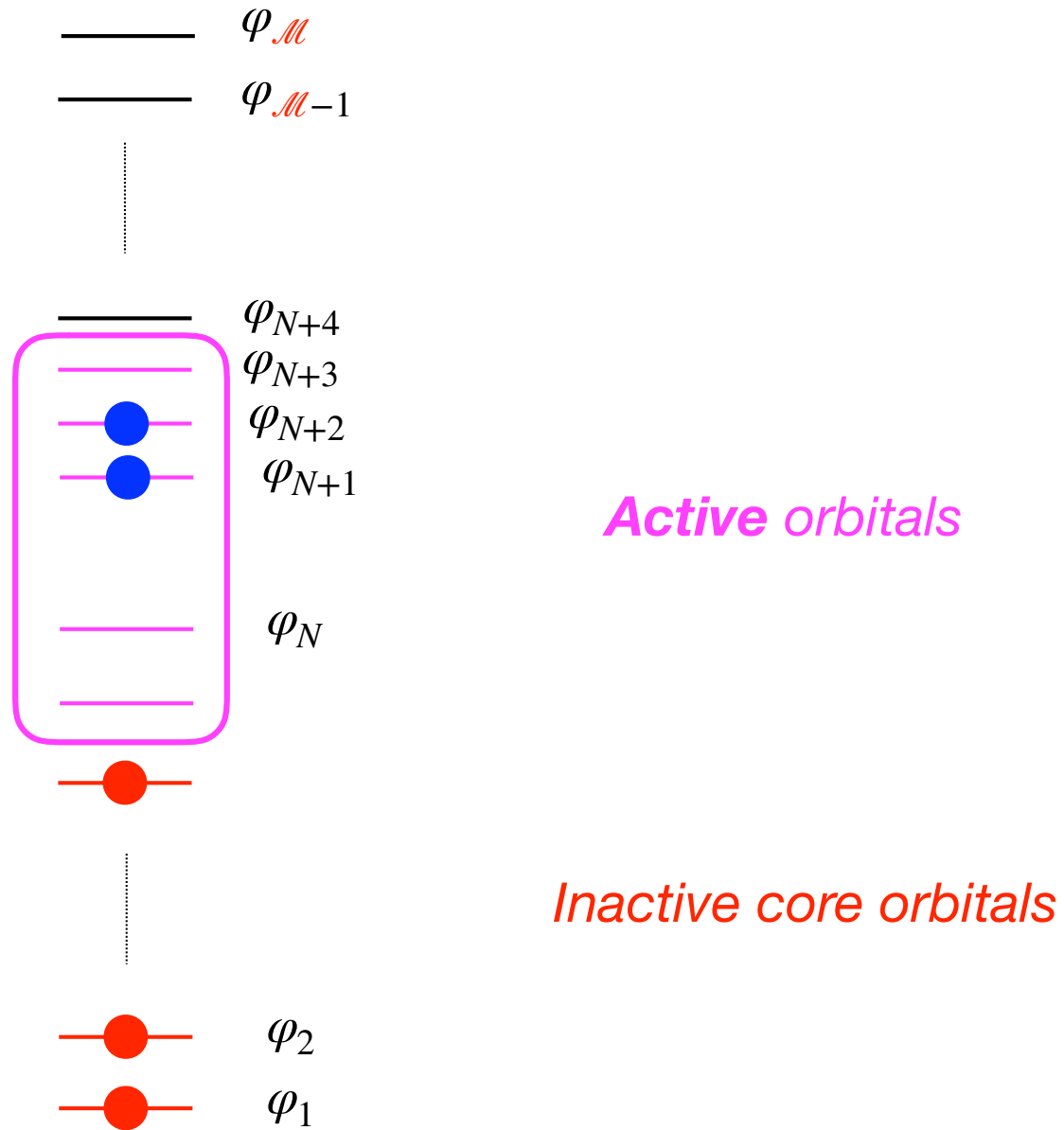
Complete active space CI (CAS-CI) method



or

Etc...

Complete active space CI (CAS-CI) method



Projection operators in the N -electron space

$$|\Psi\rangle = \sum_{\mathcal{J}}^{\text{determinants}} C_{\mathcal{J}} |\text{det}_{\mathcal{J}}\rangle$$

Projection operators in the N -electron space

$$|\Psi\rangle = \sum_{\mathcal{J}}^{\text{determinants}} C_{\mathcal{J}} |\text{det}_{\mathcal{J}}\rangle$$



$$\langle \text{det}_{\mathcal{L}} | \Psi \rangle = C_{\mathcal{L}}$$

Projection operators in the N -electron space

$$|\Psi\rangle = \sum_{\mathcal{J}}^{\text{determinants}} C_{\mathcal{J}} |\text{det}_{\mathcal{J}}\rangle$$



$$\langle \text{det}_{\mathcal{L}} | \Psi \rangle = C_{\mathcal{L}}$$



$$\sum_{\mathcal{L}}^{\text{determinants}} |\text{det}_{\mathcal{L}}\rangle \langle \text{det}_{\mathcal{L}} | \Psi \rangle = |\Psi\rangle$$

Projection operators in the N -electron space

determinants

$$\sum_{\mathcal{L}} |\det_{\mathcal{L}}\rangle \langle \det_{\mathcal{L}}| \equiv \hat{1} \quad \leftarrow \text{Identity operator } (\hat{1} |\Psi\rangle = |\Psi\rangle)$$

determinants

$$\sum_{\mathcal{L}} |\det_{\mathcal{L}}\rangle \langle \det_{\mathcal{L}}| \Psi\rangle = |\Psi\rangle$$

Projection operators in the N -electron space

determinants

$$\sum_{\mathcal{L}} |\det_{\mathcal{L}}\rangle \langle \det_{\mathcal{L}}| \equiv \hat{1}$$

“resolution of the identity”

Projection operators in the N -electron space

determinants

$$\sum_{\mathcal{L}} |\det_{\mathcal{L}}\rangle \langle \det_{\mathcal{L}}| \equiv \hat{1}$$

“resolution of the identity”



determinants

$$\sum_{\mathcal{L} \in \text{CAS}}$$

$$|\det_{\mathcal{L}}\rangle \langle \det_{\mathcal{L}}| +$$

determinants

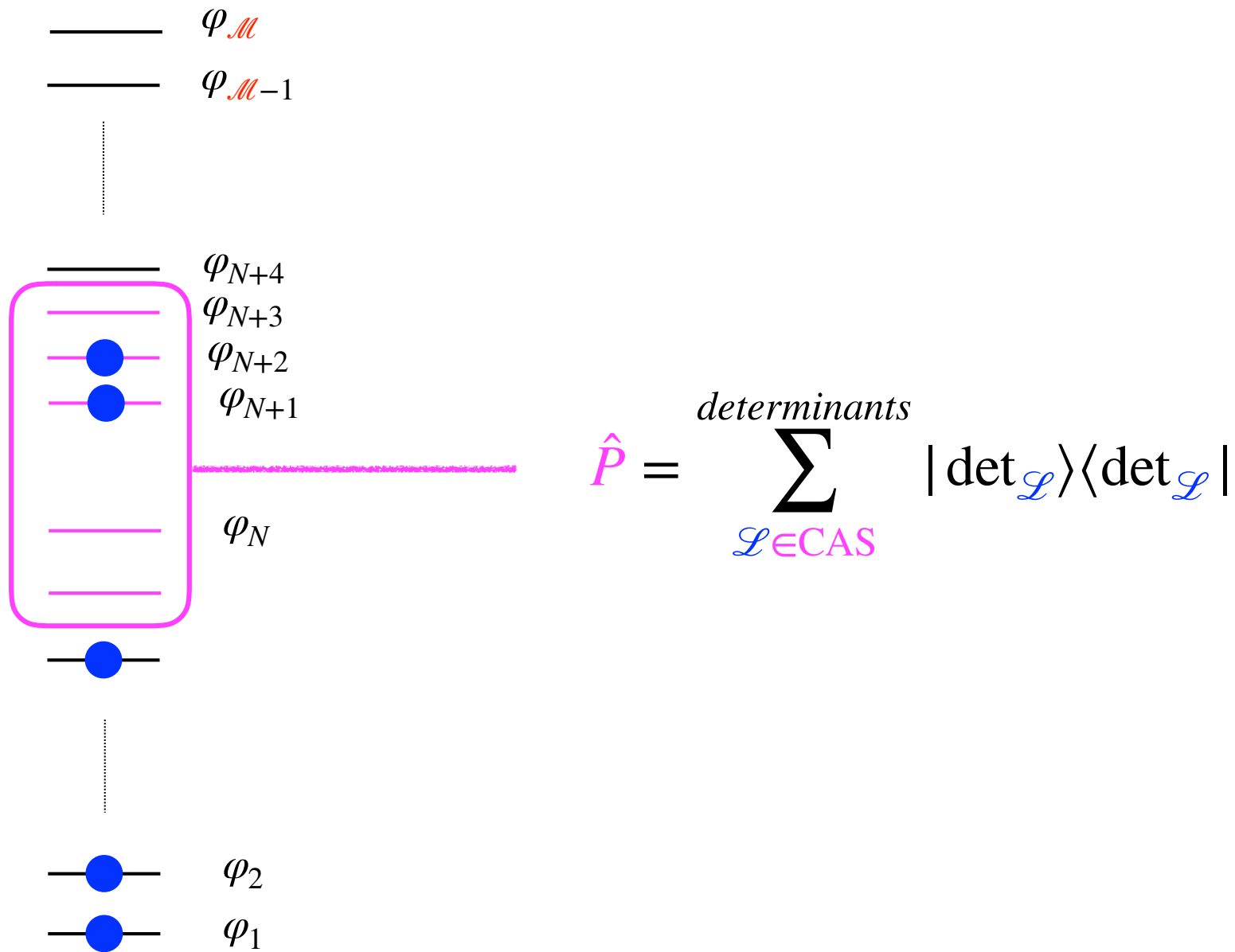
$$\sum_{\mathcal{L} \notin \text{CAS}}$$

$$|\det_{\mathcal{L}}\rangle \langle \det_{\mathcal{L}}| \equiv \hat{1}$$

Projection operators in the N -electron space

$$\underbrace{\sum_{\mathcal{L} \in \text{CAS}}^{\text{determinants}} |\text{det}_{\mathcal{L}}\rangle \langle \text{det}_{\mathcal{L}}| + \sum_{\mathcal{L} \notin \text{CAS}}^{\text{determinants}} |\text{det}_{\mathcal{L}}\rangle \langle \text{det}_{\mathcal{L}}|}_{\hat{P}} \equiv \hat{1}$$

Complete active space (CAS) and effective Hamiltonian



Projection operators in the N -electron space

$$\underbrace{\sum_{\mathcal{L} \in \text{CAS}}^{\text{determinants}} |\text{det}_{\mathcal{L}}\rangle \langle \text{det}_{\mathcal{L}}|}_{\hat{P}} + \underbrace{\sum_{\mathcal{L} \notin \text{CAS}}^{\text{determinants}} |\text{det}_{\mathcal{L}}\rangle \langle \text{det}_{\mathcal{L}}|}_{\hat{Q}} \equiv \hat{1}$$

Effective Hamiltonian

$$\hat{H}|\Psi\rangle = E|\Psi\rangle$$

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$$|\Psi\rangle = (\hat{P} + \hat{Q})|\Psi\rangle = \hat{P}|\Psi\rangle + \hat{Q}|\Psi\rangle$$

Effective Hamiltonian

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$$|\Psi\rangle = (\hat{P} + \hat{Q})|\Psi\rangle = \hat{P}|\Psi\rangle + \hat{Q}|\Psi\rangle$$



To-be-embedded part of the N -electron wave function

Effective Hamiltonian

$$\hat{H}|\Psi\rangle = E|\Psi\rangle$$

$$|\Psi\rangle = (\hat{P} + \hat{Q})|\Psi\rangle = \hat{P}|\Psi\rangle + \hat{Q}|\Psi\rangle$$

Can be determined
from $\hat{P}|\Psi\rangle$

Proof:


$$\hat{H}|\Psi\rangle = E|\Psi\rangle$$

$$\hat{Q}\hat{H}|\Psi\rangle = E\hat{Q}|\Psi\rangle$$

$$\hat{Q}\hat{H}(\hat{P} + \hat{Q})|\Psi\rangle = E\hat{Q}|\Psi\rangle$$

$$\hat{Q}\hat{H}\hat{P}|\Psi\rangle + \hat{Q}\hat{H}\hat{Q}|\Psi\rangle = E\hat{Q}|\Psi\rangle$$

$$\hat{Q}\hat{H}\hat{P}|\Psi\rangle = (E\hat{Q} - \hat{Q}\hat{H}\hat{Q})\hat{Q}|\Psi\rangle$$

$$\hat{Q}^2 = \hat{Q}$$


$$\hat{Q}|\Psi\rangle = [E\hat{Q} - \hat{Q}\hat{H}\hat{Q}]^{-1} \hat{Q}\hat{H}\hat{P}|\Psi\rangle$$

Effective Hamiltonian

$$\hat{H}|\Psi\rangle = E|\Psi\rangle$$

$$\hat{P}\hat{H}|\Psi\rangle = E\hat{P}|\Psi\rangle$$

$$\hat{P}\hat{H}(\hat{P} + \hat{Q})|\Psi\rangle = E\hat{P}|\Psi\rangle$$

$$\hat{P}\hat{H}\hat{P}|\Psi\rangle + \hat{P}\hat{H}\hat{Q}|\Psi\rangle = E\hat{P}|\Psi\rangle$$

$$\hat{P}\hat{H}\hat{P}|\Psi\rangle + \hat{P}\hat{H}\left[E\hat{Q} - \hat{Q}\hat{H}\hat{Q}\right]^{-1}\hat{Q}\hat{H}\hat{P}|\Psi\rangle = E\hat{P}|\Psi\rangle$$

$$\hat{H}_{\text{eff}}(E)\hat{P}|\Psi\rangle = E\hat{P}|\Psi\rangle$$

Effective Hamiltonian

$$\hat{H}_{\text{eff}}(E) \hat{P} |\Psi\rangle = E \hat{P} |\Psi\rangle$$

$$\hat{H}_{\text{eff}}(E) = \hat{P} \hat{H} \hat{P} + \hat{P} \hat{H} \left[E \hat{Q} - \hat{Q} \hat{H} \hat{Q} \right]^{-1} \hat{Q} \hat{H} \hat{P}$$

Effective Hamiltonian

$$\hat{H}_{\text{eff}}(E) \hat{P} |\Psi\rangle = E \hat{P} |\Psi\rangle$$

Embedded wave function

$$\hat{H}_{\text{eff}}(E) = \hat{P}\hat{H}\hat{P} + \hat{P}\hat{H} \left[E\hat{Q} - \hat{Q}\hat{H}\hat{Q} \right]^{-1} \hat{Q}\hat{H}\hat{P}$$

Effective Hamiltonian

$$\hat{H}_{\text{eff}}(E) \hat{P} |\Psi\rangle = E \hat{P} |\Psi\rangle$$

“Embedding Hamiltonian”

$$\hat{H}_{\text{eff}}(E) = \hat{P} \hat{H} \hat{P} + \hat{P} \hat{H} \left[E \hat{Q} - \hat{Q} \hat{H} \hat{Q} \right]^{-1} \hat{Q} \hat{H} \hat{P}$$

Effective Hamiltonian

$$\hat{H}_{\text{eff}}(E) \hat{P} |\Psi\rangle = E \hat{P} |\Psi\rangle$$

Energy (i.e frequency- or time-) dependent

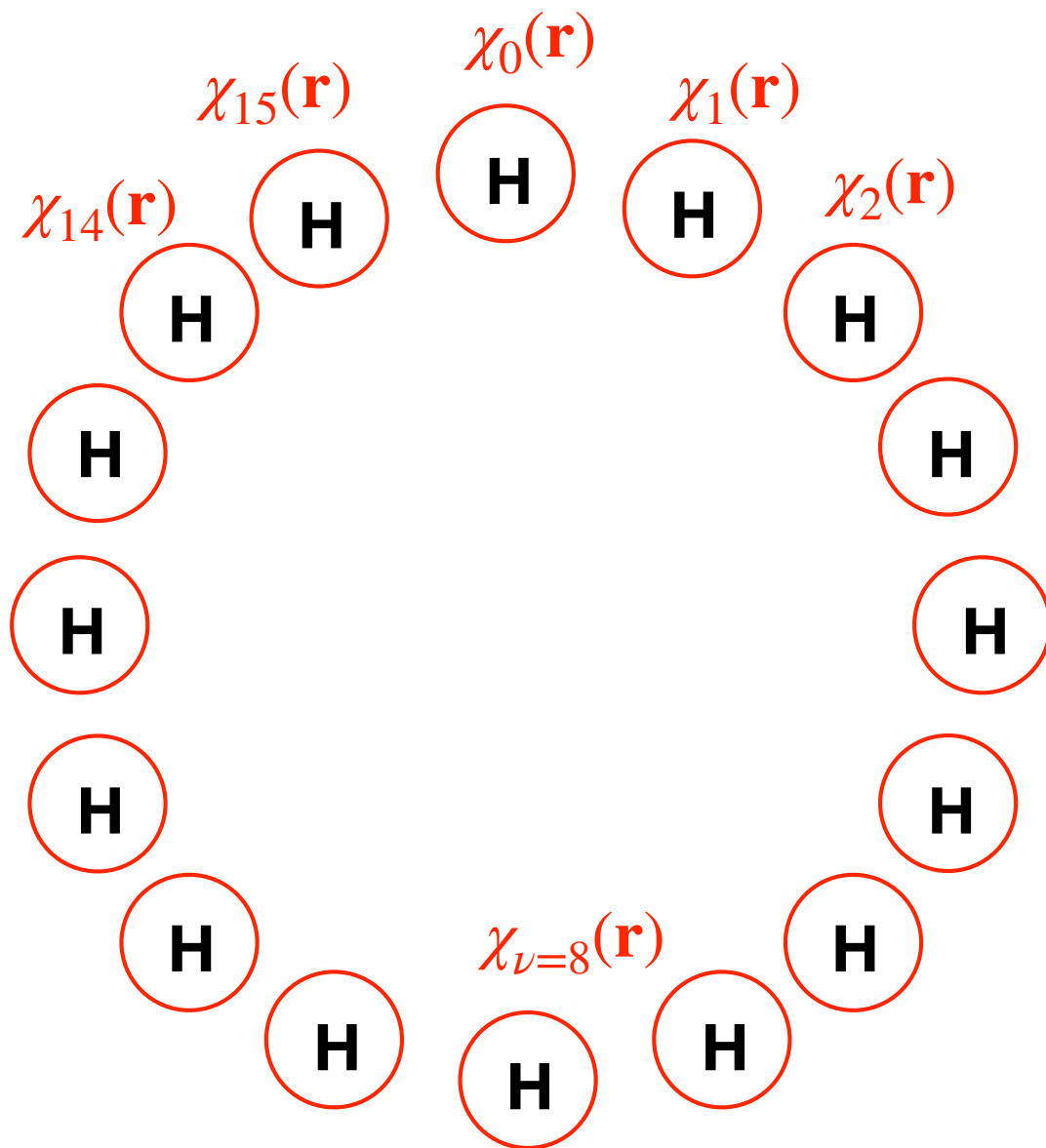
$$\hat{H}_{\text{eff}}(E) = \hat{P} \hat{H} \hat{P} + \hat{P} \hat{H} \left[E \hat{Q} - \hat{Q} \hat{H} \hat{Q} \right]^{-1} \hat{Q} \hat{H} \hat{P}$$

Strategy 3: Quantum embedding of localized orbitals

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*In this third approach we will proceed with a **fragmentation** of the **electronic Hamiltonian** (written in a **localised** orbital basis) and the design of embedding orbitals*

Prototypical ring of $L = 16$ hydrogen atoms

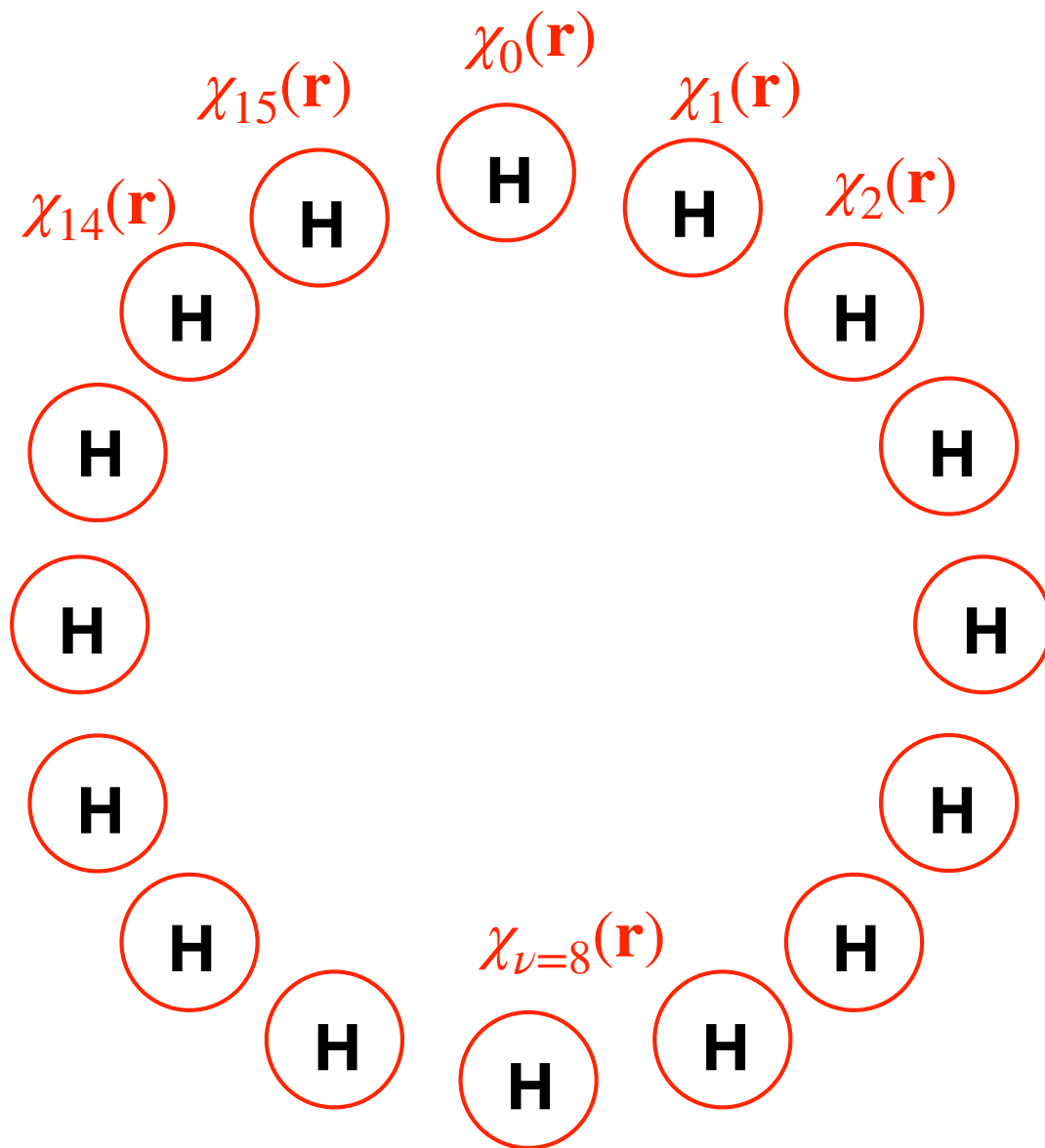


Molecular spin-orbitals

$$\varphi_P(\mathbf{x}) = \sum_{\nu} C_{\nu P} \chi_{\nu}(\mathbf{x})$$

Delocalised over the ring

Prototypical ring of $L = 16$ hydrogen atoms



Delocalized molecular spin-orbitals

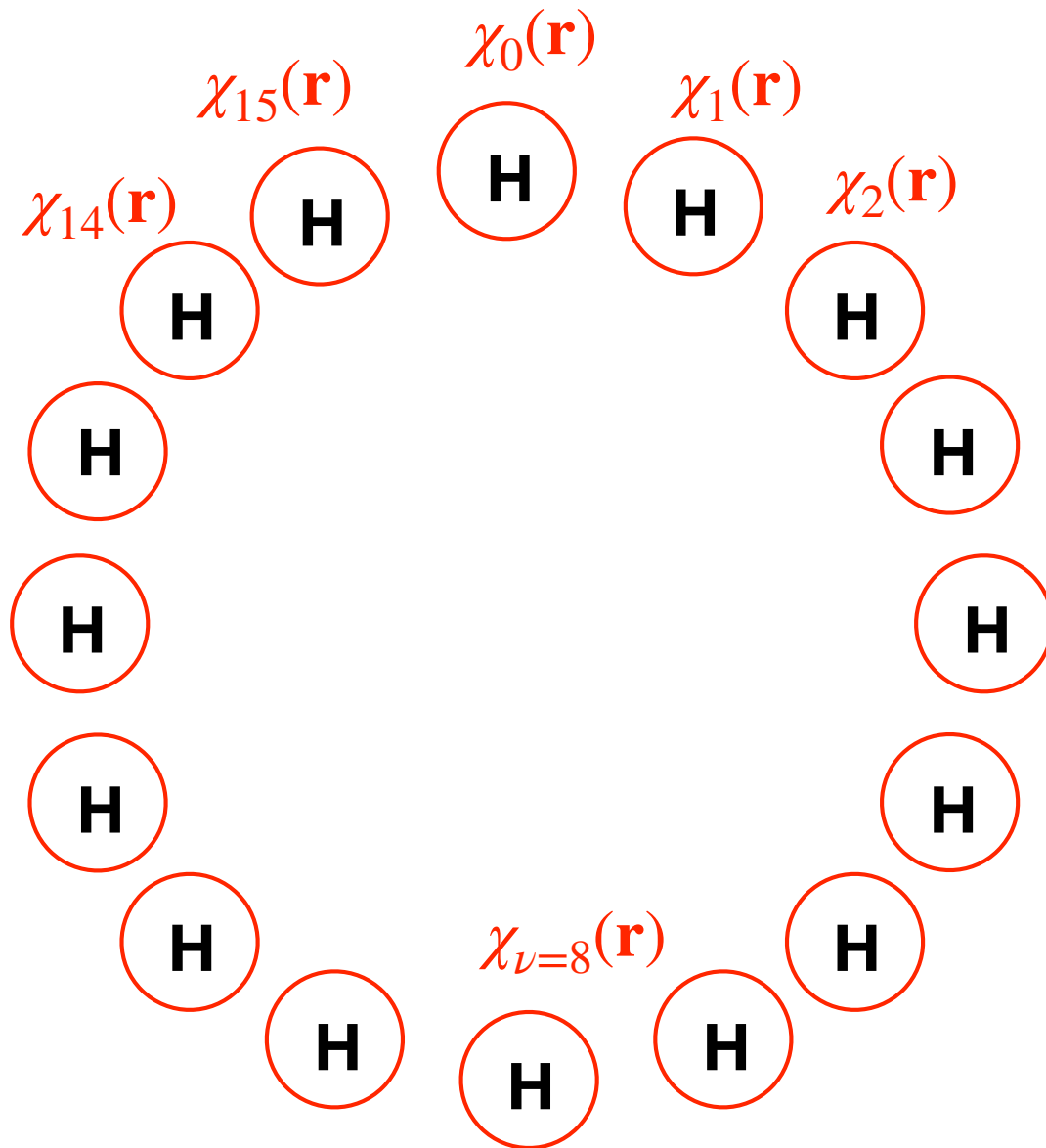
$$\varphi_P(\mathbf{x}) = \sum_{\nu} C_{\nu P} \chi_{\nu}(\mathbf{x})$$



Dirac notation

$$|\varphi_P\rangle = \sum_{\nu} C_{\nu P} |\chi_{\nu}\rangle$$

Prototypical ring of $L = 16$ hydrogen atoms



Delocalized molecular spin-orbitals

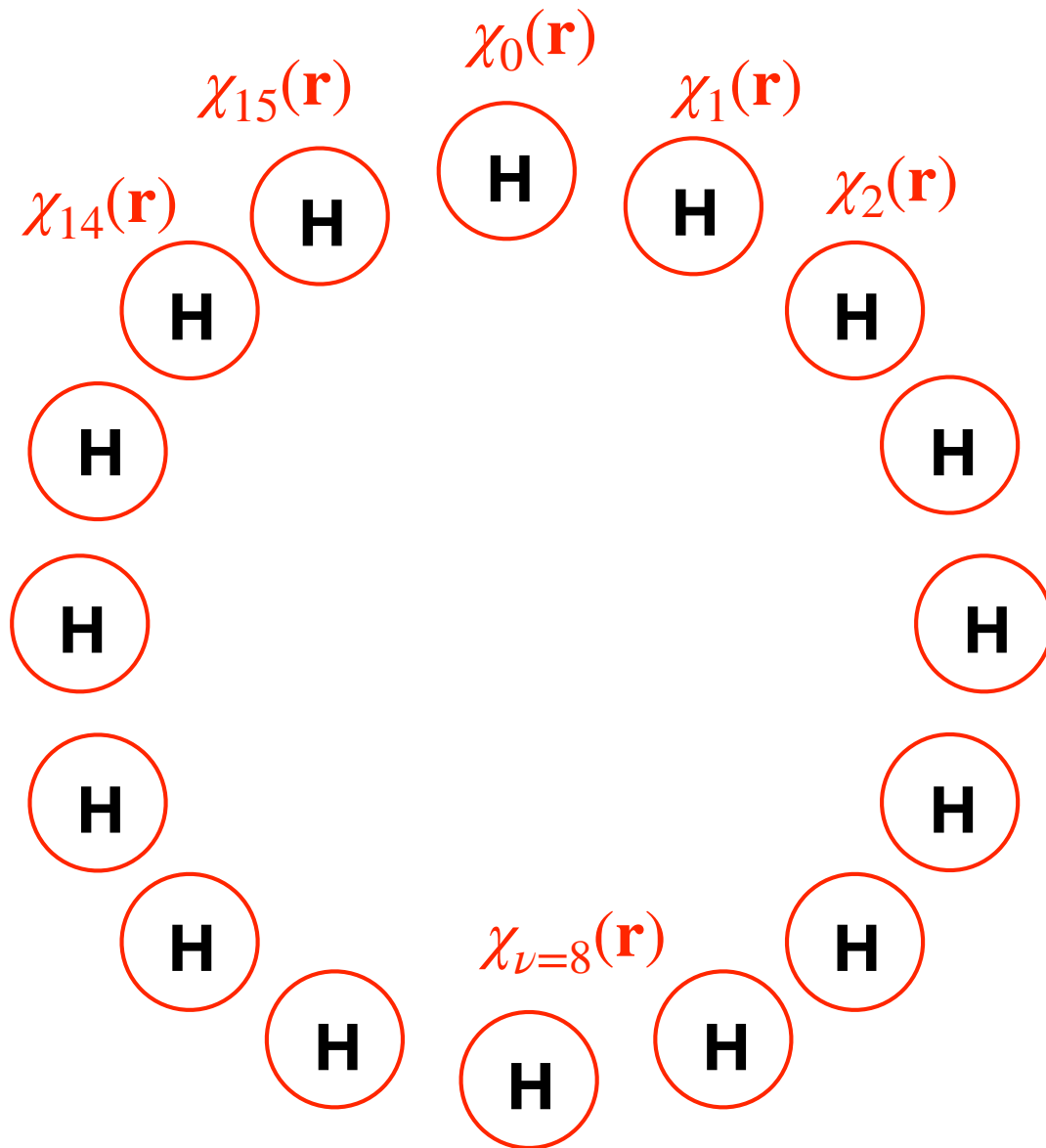
$$\varphi_P(\mathbf{x}) = \sum_{\nu} C_{\nu P} \chi_{\nu}(\mathbf{x})$$



Second-quantized notation

$$\hat{a}_P^{\dagger} |\text{vac}\rangle = \sum_{\nu} C_{\nu P} \hat{c}_{\nu}^{\dagger} |\text{vac}\rangle$$

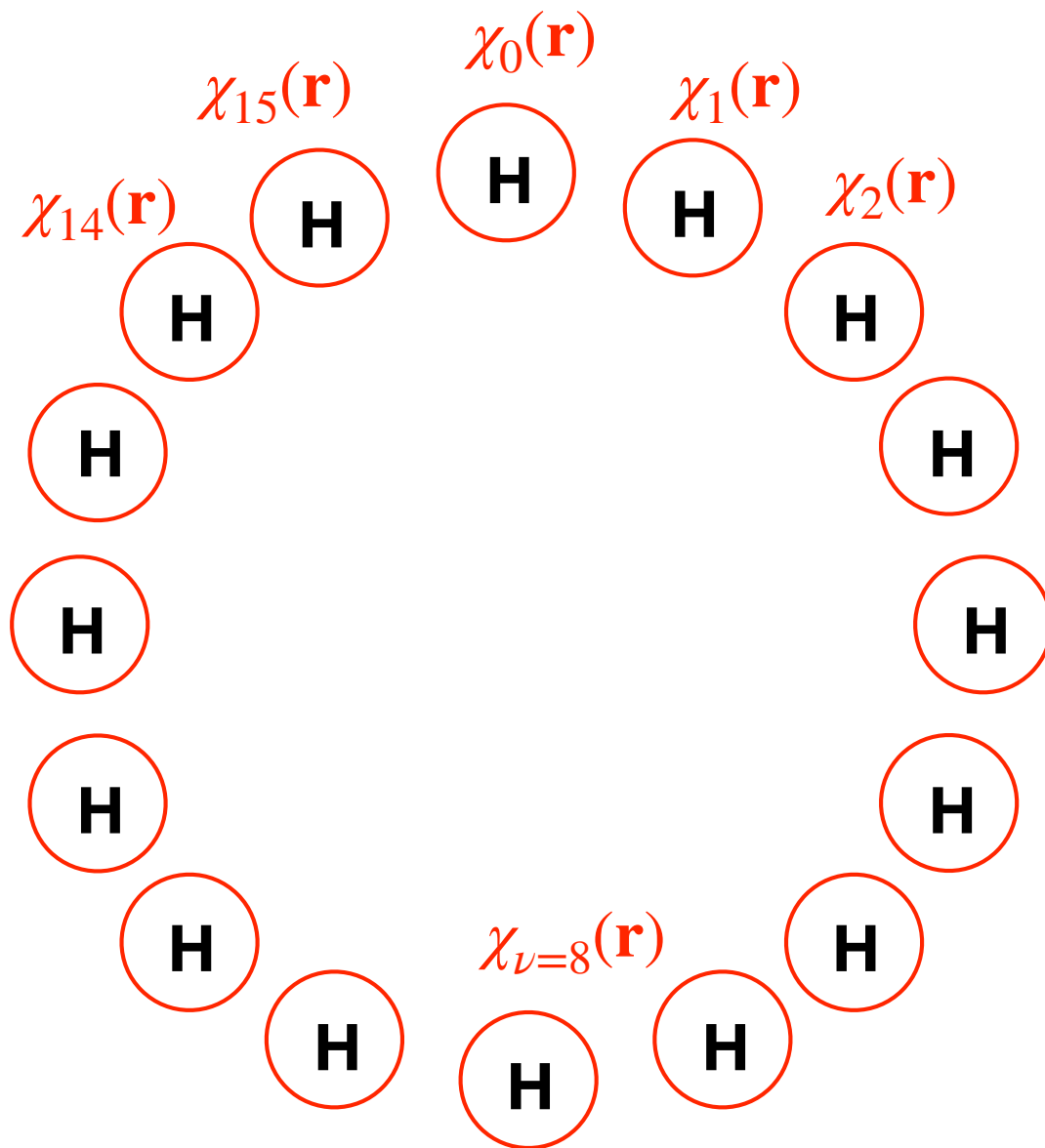
Prototypical ring of $L = 16$ hydrogen atoms



Creates an electron that occupies the **localised** spin-orbital χ_{ν}

$$\hat{a}_P^\dagger |\text{vac}\rangle = \sum_{\nu} C_{\nu P} \hat{c}_{\nu}^\dagger |\text{vac}\rangle$$

Prototypical ring of $L = 16$ hydrogen atoms



Delocalized molecular spin-orbitals

$$\varphi_P(\mathbf{x}) = \sum_{\nu} C_{\nu P} \chi_{\nu}(\mathbf{x})$$



Second-quantized notation

$$\hat{a}_P^{\dagger} = \sum_{\nu} C_{\nu P} \hat{c}_{\nu}^{\dagger}$$

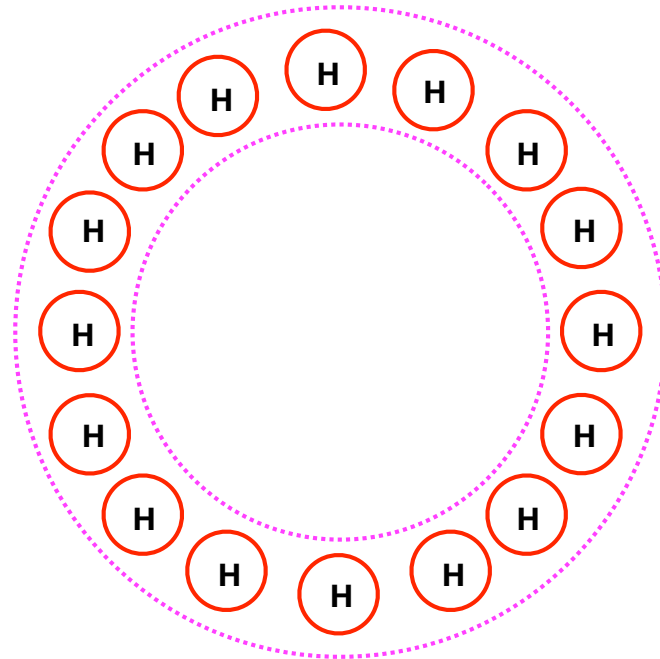
Change of representation in second quantization

$$\hat{a}_P^\dagger = \sum_{\nu} C_{\nu P} \hat{c}_{\nu}^\dagger$$

Change of representation in second quantization

Chemist's **delocalized**
representation

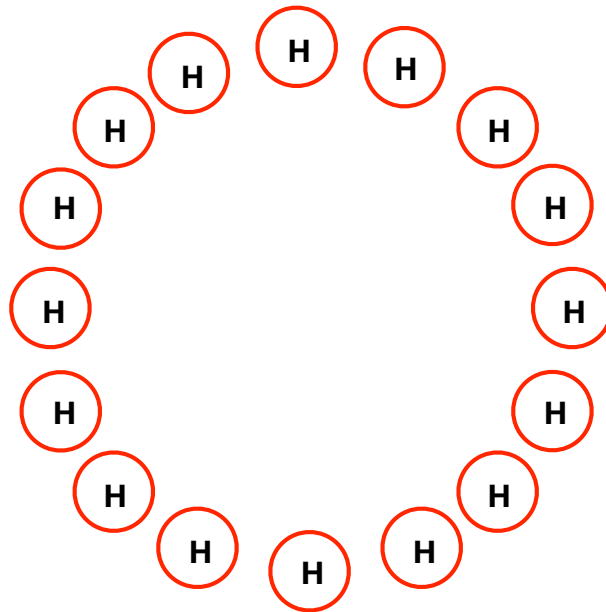
$$\hat{a}_P^\dagger = \sum_{\nu} C_{\nu P} \hat{c}_{\nu}^\dagger$$



$$\varphi_P(\mathbf{x}) = \sum_{\nu} C_{\nu P} \chi_{\nu}(\mathbf{x})$$

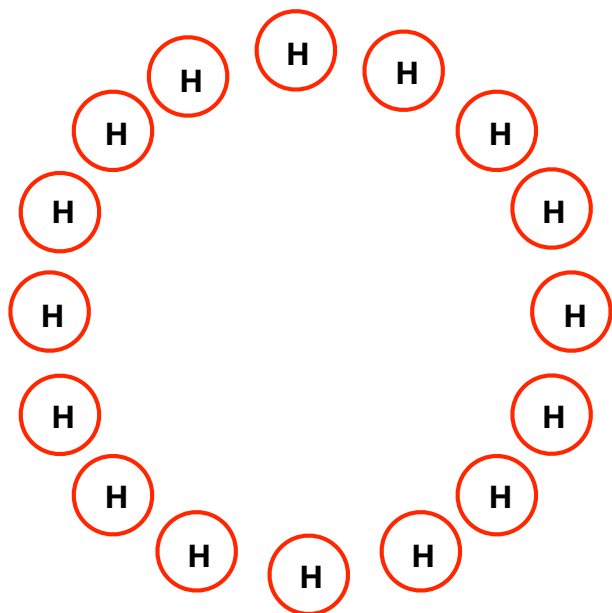
Change of representation in second quantization

$$\hat{a}_P^\dagger = \sum_{\nu} C_{\nu P} \hat{c}_{\nu}^\dagger \longrightarrow \text{Physicist's } \textit{localized} \text{ representation}$$



Change of representation in second quantization

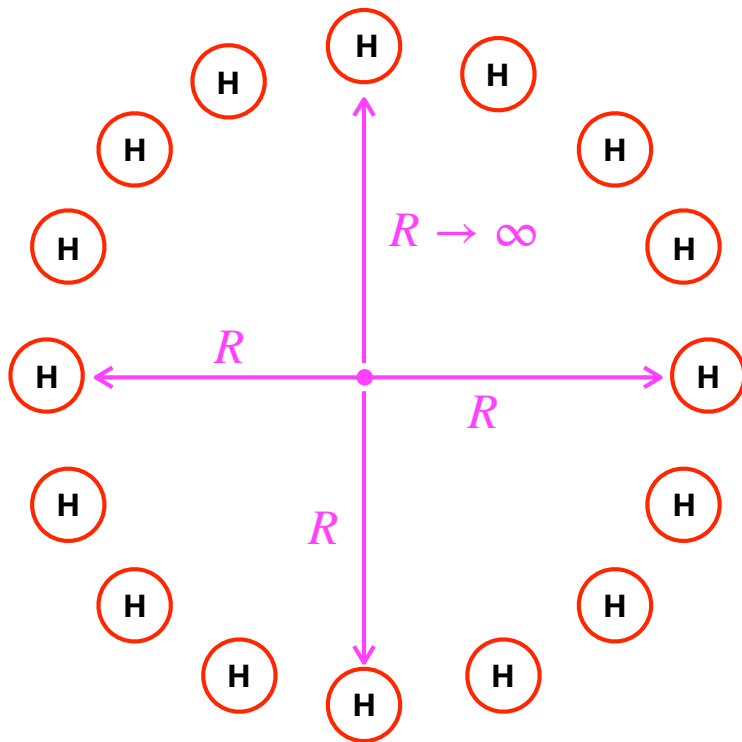
$$\hat{a}_P^\dagger = \sum_{\nu} C_{\nu P} \hat{c}_{\nu}^\dagger \longrightarrow \text{Physicist's } \mathbf{localized} \text{ representation}$$



Allows for a **fragmentation** of the molecule
in the orbital space!

Change of representation in second quantization

$$\hat{a}_P^\dagger = \sum_{\nu} C_{\nu P} \hat{c}_{\nu}^\dagger \longrightarrow \text{Physicist's } \mathbf{localized} \text{ representation}$$



Allows for a **fragmentation** of the molecule
in the orbital space!

This is a relevant strategy for describing
strong (local) electron correlation effects.

Hamiltonian in the localised representation

Delocalized representation

$$\hat{H} \equiv \sum_{PQ} \langle \varphi_P | \hat{h} | \varphi_Q \rangle \hat{a}_P^\dagger \hat{a}_Q + \frac{1}{2} \sum_{PQRS} \langle \varphi_P \varphi_Q | \hat{g} | \varphi_R \varphi_S \rangle \hat{a}_P^\dagger \hat{a}_Q^\dagger \hat{a}_S \hat{a}_R$$

Hamiltonian in the localised representation

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$$\hat{a}_P^\dagger = \sum_{\mu} C_{\mu P} \hat{c}_{\mu}^\dagger$$

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$$\hat{a}_P^\dagger = \sum_{\mu} C_{\mu P} \hat{c}_{\mu}^\dagger$$

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Hamiltonian in the localised representation

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$$\hat{a}_P^\dagger = \sum_{\mu} C_{\mu P} \hat{c}_{\mu}^\dagger$$

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$$\sum_{\mu\nu} \left(\sum_{PQ} C_{\mu P} \langle \varphi_P | \hat{h} | \varphi_Q \rangle C_{\nu Q} \right) \hat{c}_{\mu}^\dagger \hat{c}_{\nu}$$

$$\bar{h}_{\mu\nu}$$

Hamiltonian in the localised representation

Delocalized representation

$$\hat{H} \equiv \sum_{PQ} \langle \varphi_P | \hat{h} | \varphi_Q \rangle \hat{a}_P^\dagger \hat{a}_Q + \frac{1}{2} \sum_{PQRS} \langle \varphi_P \varphi_Q | \hat{g} | \varphi_R \varphi_S \rangle \hat{a}_P^\dagger \hat{a}_Q^\dagger \hat{a}_S \hat{a}_R$$

$$\hat{a}_P^\dagger = \sum_{\mu} C_{\mu P} \hat{c}_{\mu}^\dagger$$

$$\hat{a}_Q = \sum_{\nu} C_{\nu Q} \hat{c}_{\nu}$$

$$\sum_{\mu\nu} \left(\sum_{PQ} C_{\mu P} \langle \varphi_P | \hat{h} | \varphi_Q \rangle C_{\nu Q} \right) \hat{c}_{\mu}^\dagger \hat{c}_{\nu}$$

$\bar{h}_{\mu\nu}$

One-electron integrals
in the localised orbital basis

Hamiltonian in the localised representation

Delocalized representation

$$\hat{H} \equiv \sum_{PQ} \langle \varphi_P | \hat{h} | \varphi_Q \rangle \hat{a}_P^\dagger \hat{a}_Q + \frac{1}{2} \sum_{PQRS} \langle \varphi_P \varphi_Q | \hat{g} | \varphi_R \varphi_S \rangle \hat{a}_P^\dagger \hat{a}_Q^\dagger \hat{a}_S \hat{a}_R$$

III

$$\sum_{\mu\nu} \bar{h}_{\mu\nu} \hat{c}_\mu^\dagger \hat{c}_\nu$$

Localized representation

Hamiltonian in the localised representation

Delocalized representation

$$\hat{H} \equiv \sum_{PQ} \langle \varphi_P | \hat{h} | \varphi_Q \rangle \hat{a}_P^\dagger \hat{a}_Q + \frac{1}{2} \sum_{PQRS} \langle \varphi_P \varphi_Q | \hat{g} | \varphi_R \varphi_S \rangle \hat{a}_P^\dagger \hat{a}_Q^\dagger \hat{a}_S \hat{a}_R$$

III

$$\sum_{\mu\nu} \bar{h}_{\mu\nu} \hat{c}_\mu^\dagger \hat{c}_\nu$$

Localized representation

III

$$\frac{1}{2} \sum_{\mu\nu\lambda\tau} \bar{g}_{\mu\nu\lambda\tau} \hat{c}_\mu^\dagger \hat{c}_\nu^\dagger \hat{c}_\lambda \hat{c}_\tau$$

Localized representation

Hamiltonian in the localised representation

$$\hat{H} \equiv \sum_{PQ} \langle \varphi_P | \hat{h} | \varphi_Q \rangle \hat{a}_P^\dagger \hat{a}_Q + \frac{1}{2} \sum_{PQRS} \langle \varphi_P \varphi_Q | \hat{g} | \varphi_R \varphi_S \rangle \hat{a}_P^\dagger \hat{a}_Q^\dagger \hat{a}_S \hat{a}_R$$

|||

$$\frac{1}{2} \sum_{\mu\nu\lambda\tau} \bar{g}_{\mu\nu\tau\lambda} \hat{c}_\mu^\dagger \hat{c}_\nu^\dagger \hat{c}_\lambda \hat{c}_\tau$$

*Two-electron integrals
in the localised orbital basis*

$$\sum_{PQRS} C_{\mu P} C_{\nu Q} \langle \varphi_P \varphi_Q | \hat{g} | \varphi_R \varphi_S \rangle C_{\tau R} C_{\lambda S}$$

In summary...

$$\hat{H} \equiv \sum_{PQ} \langle \varphi_P | \hat{h} | \varphi_Q \rangle \hat{a}_P^\dagger \hat{a}_Q + \frac{1}{2} \sum_{PQRS} \langle \varphi_P \varphi_Q | \hat{g} | \varphi_R \varphi_S \rangle \hat{a}_P^\dagger \hat{a}_Q^\dagger \hat{a}_S \hat{a}_R$$

Delocalized representation

||| **Localized** representation

$$\sum_{\mu\nu} \bar{h}_{\mu\nu} \hat{c}_\mu^\dagger \hat{c}_\nu + \frac{1}{2} \sum_{\mu\nu\lambda\tau} \bar{g}_{\mu\nu\lambda\tau} \hat{c}_\mu^\dagger \hat{c}_\nu^\dagger \hat{c}_\lambda \hat{c}_\tau$$

In summary...

$$\hat{H} \equiv \sum_{\mu, \nu=1}^{\mathcal{M}} \bar{h}_{\mu\nu} \hat{c}_{\mu}^{\dagger} \hat{c}_{\nu} + \frac{1}{2} \sum_{\mu, \nu, \lambda, \tau=1}^{\mathcal{M}} \bar{g}_{\mu\nu\tau\lambda} \hat{c}_{\mu}^{\dagger} \hat{c}_{\nu}^{\dagger} \hat{c}_{\lambda} \hat{c}_{\tau}$$

In summary...

$$\hat{H} \equiv \sum_{P,Q=1}^{\mathcal{M}} \bar{h}_{PQ} \hat{c}_P^\dagger \hat{c}_Q + \frac{1}{2} \sum_{P,Q,R,S=1}^{\mathcal{M}} \bar{g}_{PQRS} \hat{c}_P^\dagger \hat{c}_Q^\dagger \hat{c}_S \hat{c}_R$$

just labels...

In summary...

$$\hat{H} \equiv \sum_{PQ} \langle \varphi_P | \hat{h} | \varphi_Q \rangle \hat{a}_P^\dagger \hat{a}_Q + \frac{1}{2} \sum_{PQRS} \langle \varphi_P \varphi_Q | \hat{g} | \varphi_R \varphi_S \rangle \hat{a}_P^\dagger \hat{a}_Q^\dagger \hat{a}_S \hat{a}_R$$

Delocalized representation

|||

$$\sum_{PQ} \bar{h}_{PQ} \hat{c}_P^\dagger \hat{c}_Q + \frac{1}{2} \sum_{PQRS} \bar{g}_{PQRS} \hat{c}_P^\dagger \hat{c}_Q^\dagger \hat{c}_S \hat{c}_R$$

Localized representation

Reduced density matrices

$$E_0 = \langle \hat{H} \rangle_{\Psi_0}$$

Reduced density matrices

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Reduced density matrices

$$\begin{aligned} E_0 &= \langle \hat{H} \rangle_{\Psi_0} \\ &= \sum_{PQ} \bar{h}_{PQ} \langle \hat{c}_P^\dagger \hat{c}_Q \rangle_{\Psi_0} + \frac{1}{2} \sum_{PQRS} \bar{g}_{PQRS} \langle \hat{c}_P^\dagger \hat{c}_Q^\dagger \hat{c}_S \hat{c}_R \rangle_{\Psi_0} \end{aligned}$$

Reduced density matrices

$$E_0 = \langle \hat{H} \rangle_{\Psi_0}$$
$$= \sum_{PQ} \bar{h}_{PQ} \langle \hat{c}_P^\dagger \hat{c}_Q \rangle_{\Psi_0} + \frac{1}{2} \sum_{PQRS} \bar{g}_{PQRS} \langle \hat{c}_P^\dagger \hat{c}_Q^\dagger \hat{c}_S \hat{c}_R \rangle_{\Psi_0}$$

One-electron reduced
density matrix (1RDM)

$$\gamma_{PQ} = \langle \hat{c}_P^\dagger \hat{c}_Q \rangle_{\Psi_0}$$

Reduced density matrices

$$E_0 = \langle \hat{H} \rangle_{\Psi_0}$$

$$= \sum_{PQ} \bar{h}_{PQ} \langle \hat{c}_P^\dagger \hat{c}_Q \rangle_{\Psi_0} + \frac{1}{2} \sum_{PQRS} \bar{g}_{PQRS} \langle \hat{c}_P^\dagger \hat{c}_Q^\dagger \hat{c}_S \hat{c}_R \rangle_{\Psi_0}$$

One-electron reduced density matrix (1RDM)

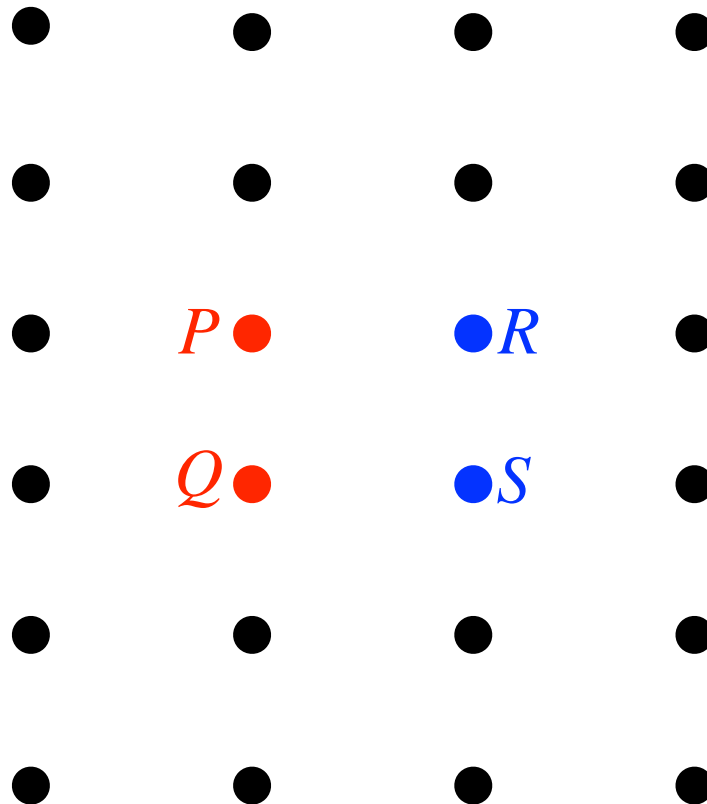
$$\gamma_{PQ} = \langle \hat{c}_P^\dagger \hat{c}_Q \rangle_{\Psi_0}$$

Two-electron reduced density matrix (2RDM)

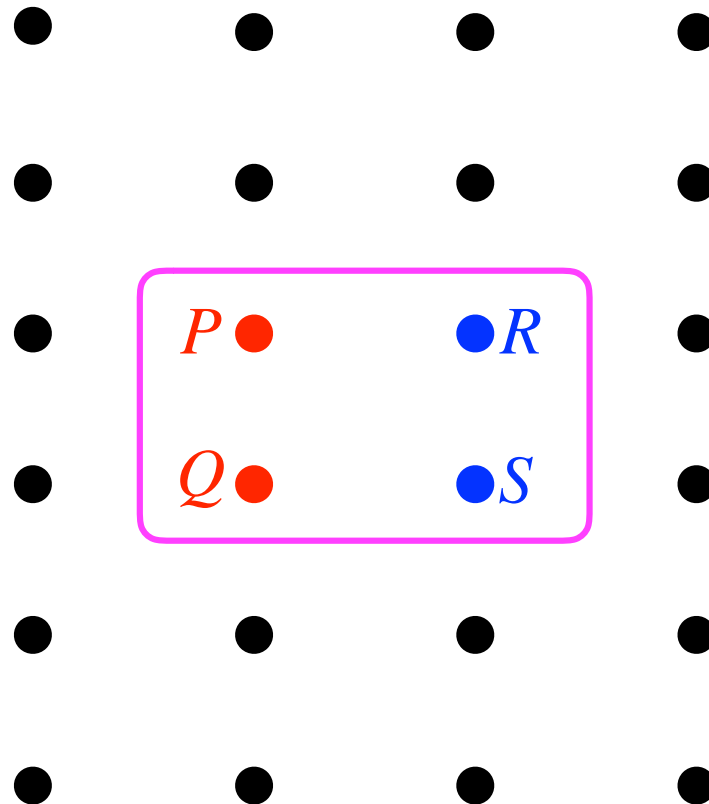
$$\Gamma_{PQRS} = \langle \hat{c}_P^\dagger \hat{c}_Q^\dagger \hat{c}_S \hat{c}_R \rangle_{\Psi_0}$$

Reduced density matrices

Let's consider a 2D lattice of localised spin-orbitals




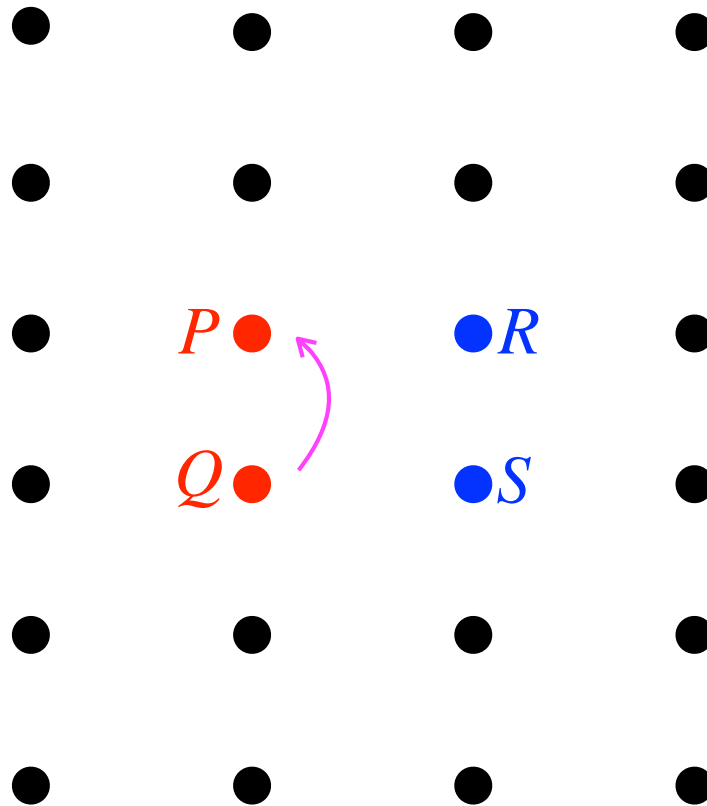
Reduced density matrices




“fragment”

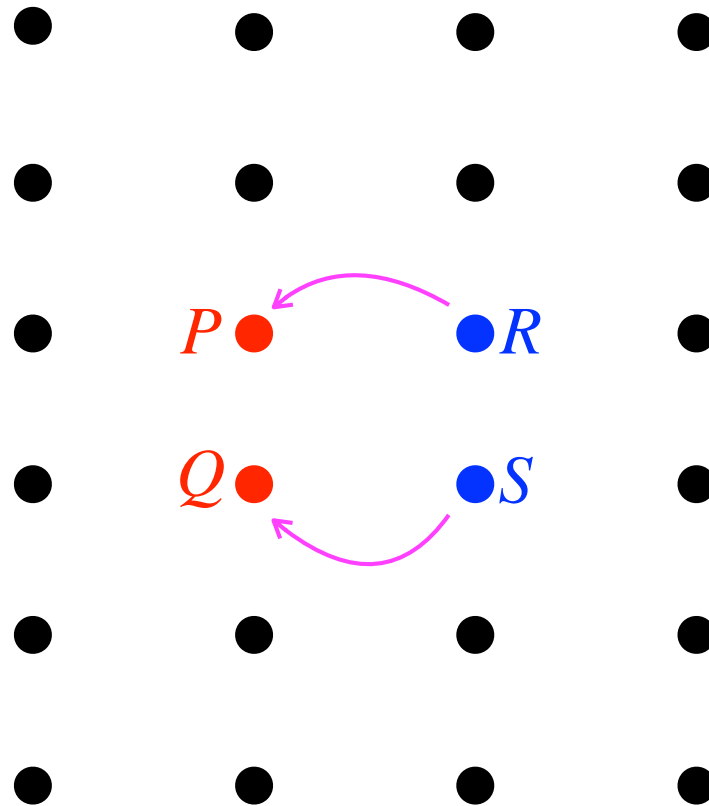
Reduced density matrices

$$\gamma_{PQ} = \langle \Psi_0 | \hat{c}_P^\dagger \hat{c}_Q | \Psi_0 \rangle$$




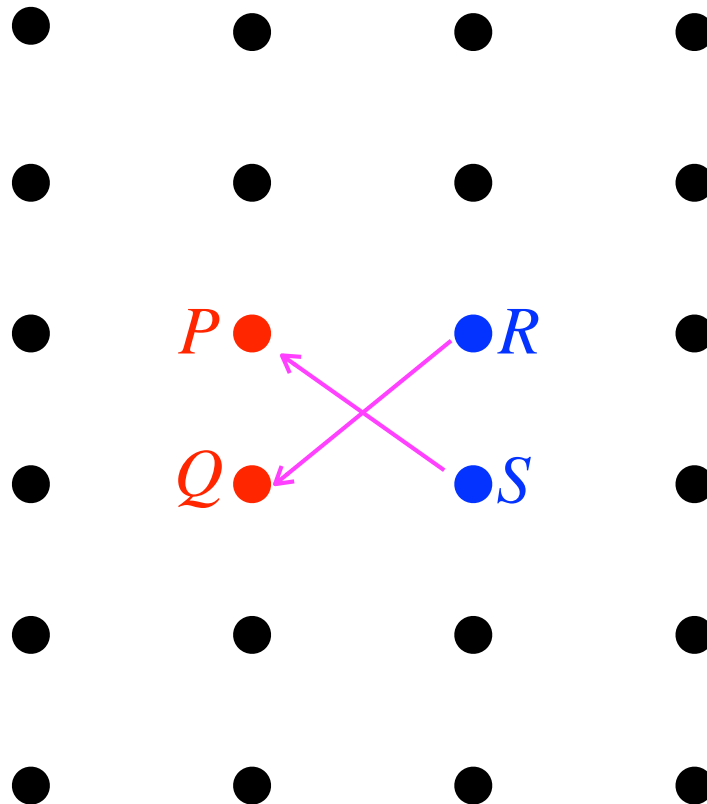
Reduced density matrices

$$\Gamma_{PQSR} = \langle \Psi_0 | \hat{c}_P^\dagger \hat{c}_Q^\dagger \hat{c}_S \hat{c}_R | \Psi_0 \rangle$$




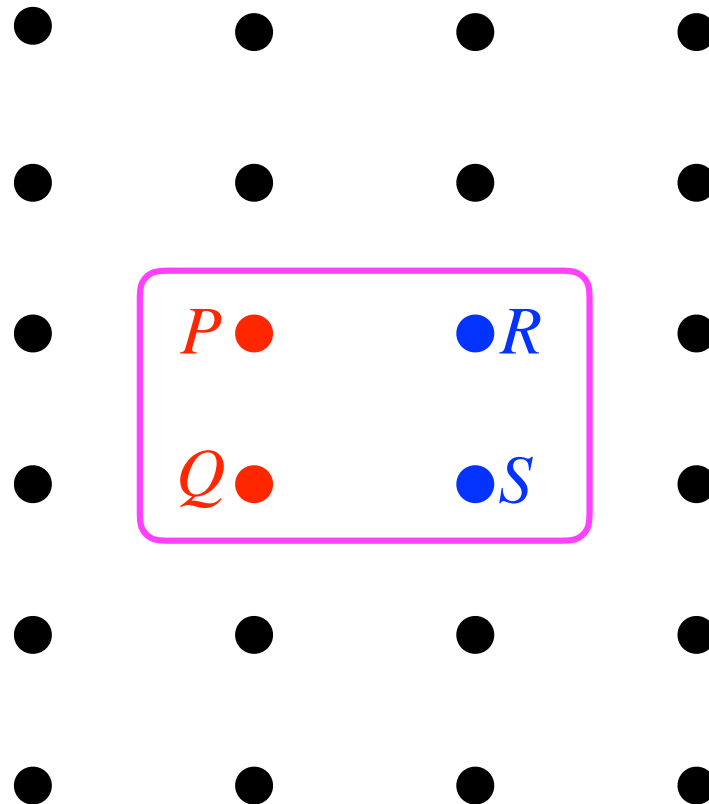
Reduced density matrices

$$\Gamma_{PQSR} = \langle \Psi_0 | \hat{c}_P^\dagger \hat{c}_Q^\dagger \hat{c}_S \hat{c}_R | \Psi_0 \rangle$$



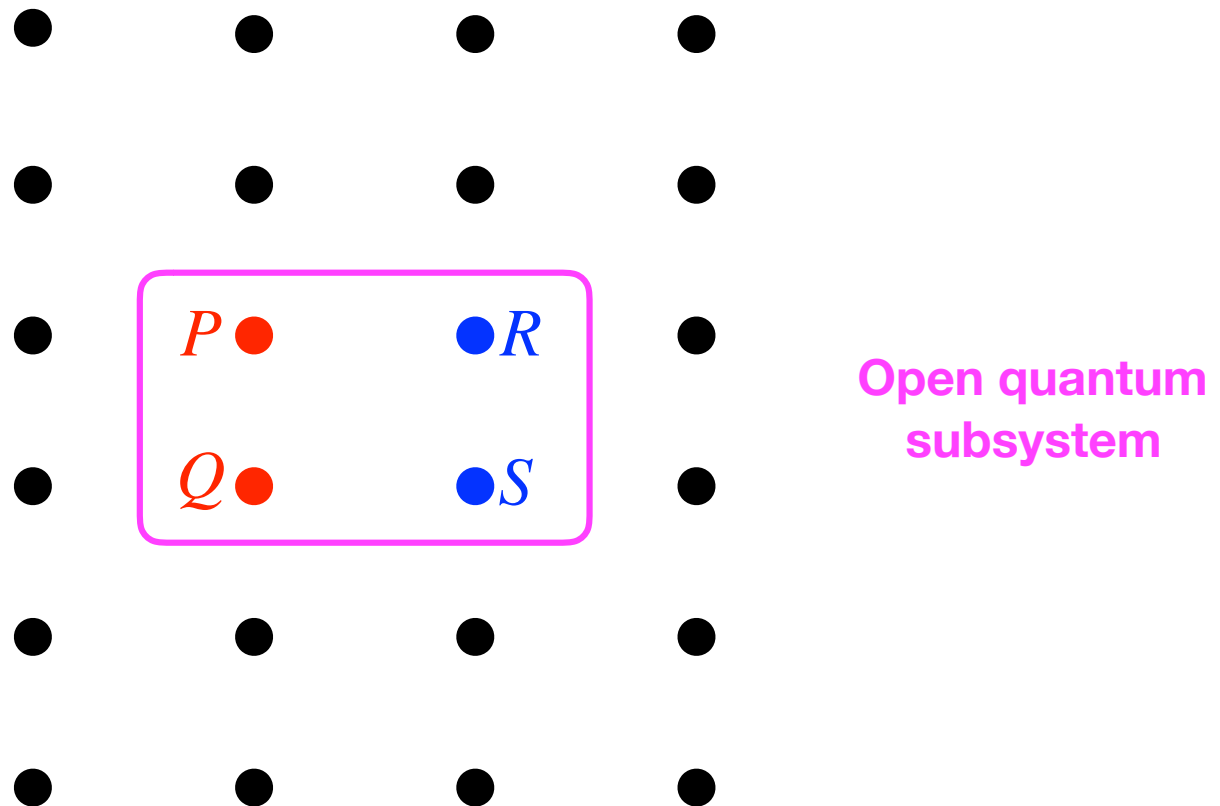
Reduced density matrices

The *PQRS* orbital fragment is **NOT disconnected** from the other orbitals



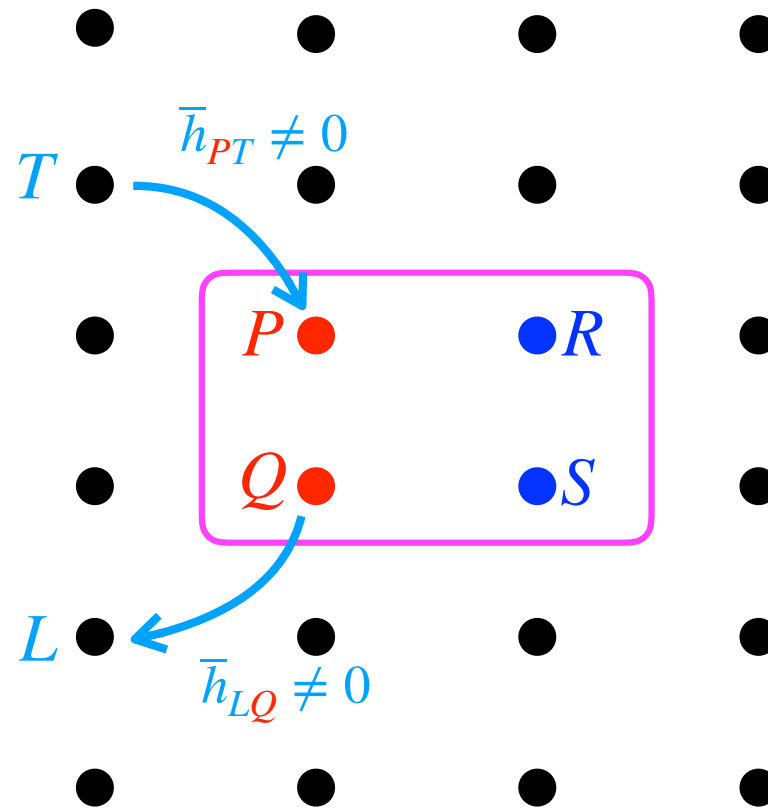
Reduced density matrices

The *PQRS* orbital fragment is **NOT disconnected** from the other orbitals



Reduced density matrices

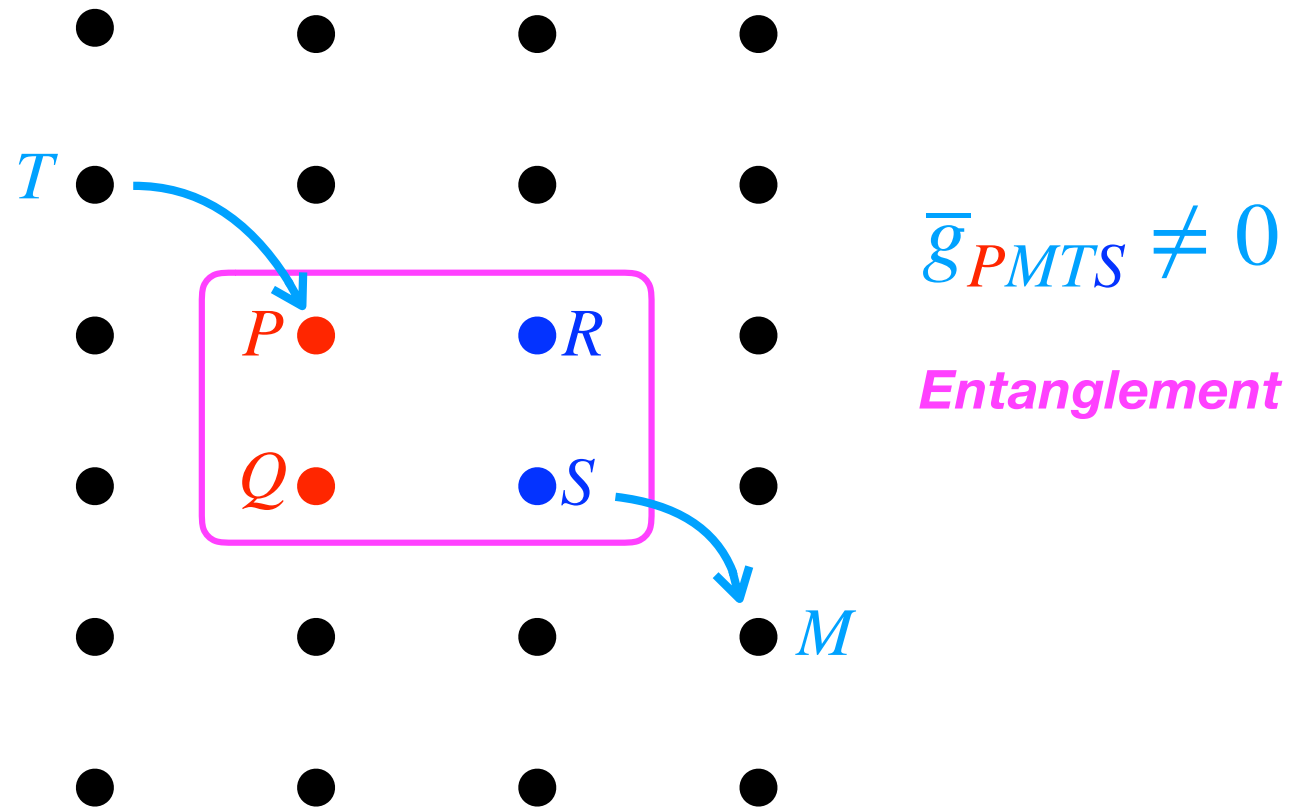
$$\hat{H} \equiv \sum_{PQ} \bar{h}_{PQ} \hat{c}_P^\dagger \hat{c}_Q + \frac{1}{2} \sum_{PQRS} \bar{g}_{PQRS} \hat{c}_P^\dagger \hat{c}_Q^\dagger \hat{c}_S \hat{c}_R$$



Entanglement

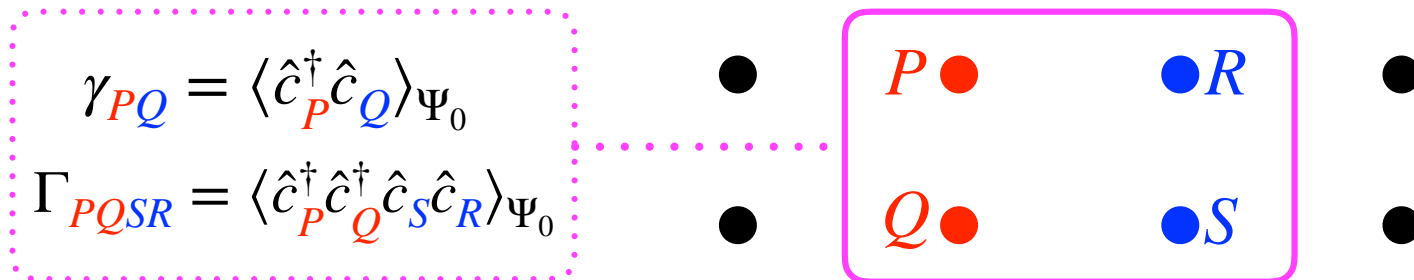
Reduced density matrices

$$\hat{H} \equiv \sum_{PQ} \bar{h}_{PQ} \hat{c}_P^\dagger \hat{c}_Q + \frac{1}{2} \sum_{PQRS} \bar{g}_{PQRS} \hat{c}_P^\dagger \hat{c}_Q^\dagger \hat{c}_S \hat{c}_R$$

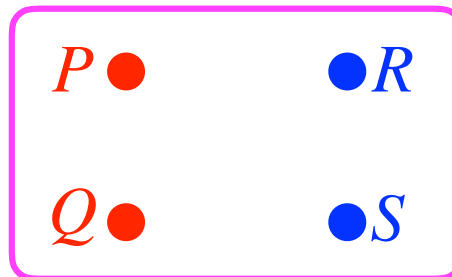


Reduced density matrices

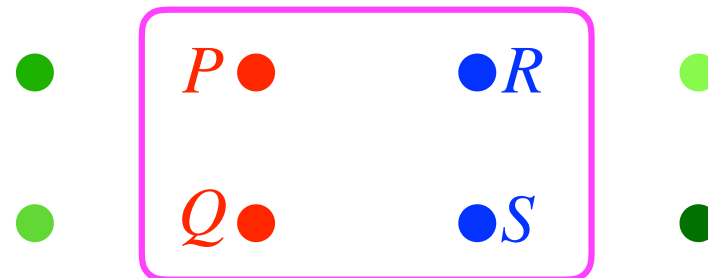
The evaluation of the RDMs requires, in principle, the **full wave function** Ψ_0



Embedding of localised spin-orbitals



Embedding of localised spin-orbitals



Embedded
fragment

Embedding of localised spin-orbitals

Bath orbitals
(reservoir)



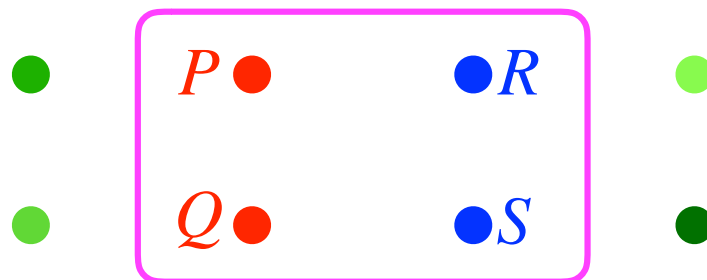
Bath orbitals
(reservoir)



Embedding of localised spin-orbitals

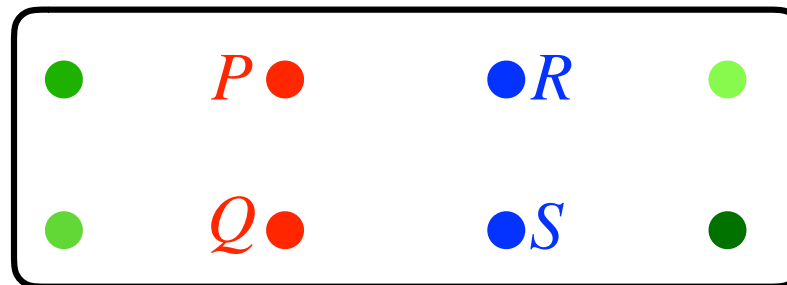
Bath orbitals
(reservoir)

Bath orbitals
(reservoir)



Embedded
fragment

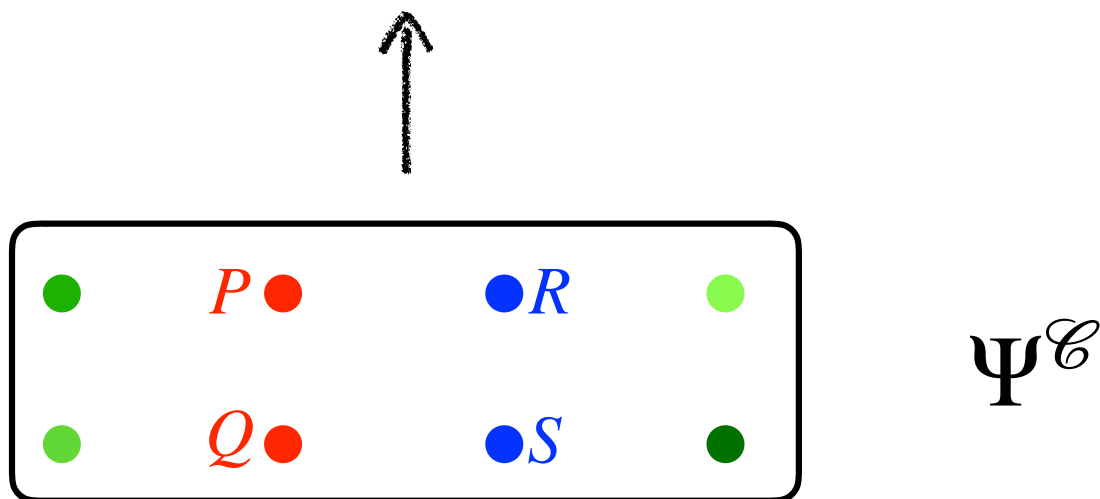
Embedding of localised spin-orbitals



Cluster

Embedding of localised spin-orbitals

Closed (and much smaller than the true system)

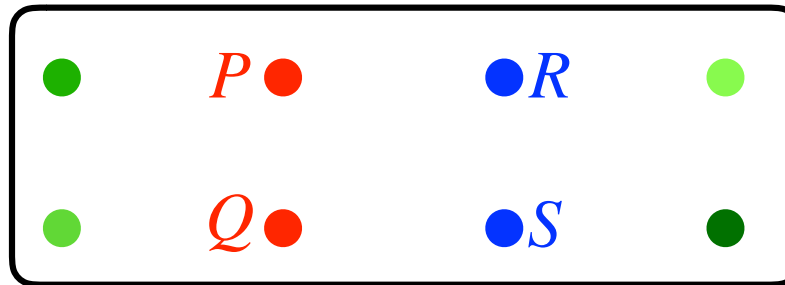


Cluster

Embedding of localised spin-orbitals

$$\gamma_{PQ} = \langle \Psi_0 | \hat{c}_P^\dagger \hat{c}_Q | \Psi_0 \rangle \approx \langle \Psi^{\mathcal{C}} | \hat{c}_P^\dagger \hat{c}_Q | \Psi^{\mathcal{C}} \rangle$$

$$\Gamma_{PQSR} = \langle \Psi_0 | \hat{c}_P^\dagger \hat{c}_Q^\dagger \hat{c}_S \hat{c}_R | \Psi_0 \rangle \approx \langle \Psi^{\mathcal{C}} | \hat{c}_P^\dagger \hat{c}_Q^\dagger \hat{c}_S \hat{c}_R | \Psi^{\mathcal{C}} \rangle$$



$\Psi^{\mathcal{C}}$

Cluster

