EUR Master "Chemistry of Complex Systems"

Exam in theory of extended systems (*part 2***)**

March 2019

one-hour exam

Neither documents nor calculators are allowed. The grading scale might be changed.

1. Questions on the lecture material (6 points)

- a) **[3 pts]** Bloch's theorem is usually explained in the particular case of one-electron systems. How would you apply Bloch's theorem to a system made of many *non-interacting* electrons ? Would the same procedure be exact if we were considering (real) *interacting* electrons instead ? If not, would it be possible to return exactly to a one-electron picture even though we describe interacting electrons ? How ?
- b) **[3 pts]** Do we need to know the non-interacting kinetic energy density functional for performing a practical Kohn–Sham density-functional theory (KS-DFT) calculation ? What are the different strategies that can be adopted for developing density-functional approximations in KS-DFT ?

2. Problem: Ionization potential in density-functional theory (16 points)

Let us consider both interacting ($\lambda = 1$) and non-interacting ($\lambda = 0$) N-electron systems where N is either equal to *N* (the system is then referred to as *neutral*) or $N-1$ (it is then referred to as *ionized*). Their normalized ground-state wavefunctions $\Psi_{\mathcal{N}}^{\lambda,\xi}$ fulfill the following Schrödinger-like equation,

$$
\left(\hat{T} + \lambda \hat{W}_{ee} + \sum_{i=1}^N v^{\lambda,\xi}(\mathbf{r}_i) \times \right) \Psi_{\mathcal{N}}^{\lambda,\xi} = E_{\mathcal{N}}^{\lambda,\xi} \Psi_{\mathcal{N}}^{\lambda,\xi}, \text{ with } \mathcal{N} = N \text{ or } N - 1.
$$
 (1)

The additional parameter *ξ* is a *weight* that will be assigned to the ionized state in the following. The local potential $v^{\lambda,\xi}(\mathbf{r})$ is adjusted such that the so-called ensemble density $n^{\lambda,\xi}(\mathbf{r})$, which is a weighted sum of neutral and ionized densities, *does not vary with ξ and λ*:

$$
n^{\lambda,\xi}(\mathbf{r}) = \left[1 - \frac{(N-1)\xi}{N}\right] \times n_{\Psi_N^{\lambda,\xi}}(\mathbf{r}) + \xi \times n_{\Psi_{N-1}^{\lambda,\xi}}(\mathbf{r}) = n(\mathbf{r}) \quad \forall \lambda, \ \forall \xi. \tag{2}
$$

- a) **[1 pt]** Verify that, by construction, the ensemble density integrates to the number *N* of electrons in the neutral system: $\int d\mathbf{r} \; n^{\lambda,\xi}(\mathbf{r}) = N$.
- b) **[2 pts]** Let us consider the $\xi = 0$ limit of the theory. How would you name the local potentials $v^{\lambda=1,\xi=0}(\mathbf{r})$ and $v^{\lambda=0,\xi=0}(\mathbf{r})$? Are they uniquely defined by the density constraint in Eq. (2)?
- c) **[0.5 pt]** We introduce the following density functionals,

$$
F^{\lambda,\xi}[n] = \left[1 - \frac{(N-1)\xi}{N}\right] \times \left\langle \Psi_N^{\lambda,\xi} \middle| \hat{T} + \lambda \hat{W}_{ee} \middle| \Psi_N^{\lambda,\xi} \right\rangle + \xi \times \left\langle \Psi_{N-1}^{\lambda,\xi} \middle| \hat{T} + \lambda \hat{W}_{ee} \middle| \Psi_{N-1}^{\lambda,\xi} \right\rangle,
$$

\n
$$
E_{\text{Hxc}}^{\xi}[n] = F^{\lambda=1,\xi}[n] - F^{\lambda=0,\xi}[n].
$$
\n(3)

Are they universal ?

d) **[2 pts]** Show that conventional (*N*-electron) KS-DFT is recovered from Eqs. (2) and (3) when $\xi = 0$. How would you name the functionals $F^{\lambda=1,\xi=0}[n]$, $F^{\lambda=0,\xi=0}[n]$, and $E_{\text{Hxc}}^{\xi=0}[n]$? Which energy contributions do they describe ?

e) **[2 pts]** Deduce from Eqs. (1) and (2) that
$$
F^{\lambda,\xi}[n] = \left[1 - \frac{(N-1)\xi}{N}\right] \times E_N^{\lambda,\xi} + \xi \times E_{N-1}^{\lambda,\xi} - \int d\mathbf{r} \, v^{\lambda,\xi}(\mathbf{r}) n(\mathbf{r}).
$$

f) [1.5 pts] Prove that $\frac{\partial E_N^{\lambda,\xi}}{\partial \xi} = \int d\mathbf{r} \frac{\partial v^{\lambda,\xi}(\mathbf{r})}{\partial \xi}$ $\frac{\partial \mathbf{f}}{\partial \mathbf{f}}$ *n*_V_{λ}^{*s*}_{*s*}^{\mathbf{f} (**r**) where $\mathcal{N} = N$ or $N - 1$. **Hint:** prove the} $\text{Hellmann–Feynman theorem, } \frac{\partial E_{\mathcal{N}}^{\lambda,\xi}}{\partial \xi} =$ * Ψ *λ,ξ* $\mathcal N$ *∂H*ˆ *λ,ξ ∂ξ* $\begin{array}{c} \begin{array}{c} \begin{array}{c} \end{array}\\ \begin{array}{c} \end{array} \end{array} \end{array}$ Ψ *λ,ξ* $\mathcal N$ $\left\{\n\begin{aligned}\n\text{where } \hat{H}^{\lambda,\xi} = \hat{T} + \lambda \hat{W}_{\text{ee}} + \sum_{i=1}^{N} v^{\lambda,\xi}(\mathbf{r}_i) \times\n\end{aligned}\n\right\}$ and conclude.

g) **[1.5 pts]** Deduce from Eq. (2) and questions 2. e) and 2. f) that $\frac{\partial F^{\lambda,\xi}[n]}{\partial \xi} = E^{\lambda,\xi}_{N-1} - \frac{(N-1)}{N}$ $\frac{(-1)}{N} \times E_N^{\lambda,\xi}$.

- h) **[1 pt]** Let $I_N = E_{N-1}^{\lambda=1,\xi=0} E_N^{\lambda=1,\xi=0}$ denote the exact ionization potential of a given *N*-electron system. We denote ε_N the HOMO energy (where HOMO stands for *highest occupied molecular orbital*) that would be obtained by applying standard KS-DFT to the latter system. Explain why $\varepsilon_N =$ $\left(E_{N-1}^{\lambda=0,\xi=0}-E_{N}^{\lambda=0,\xi=0}\right]$ $\chi^{(k)}(X) = 0$, $\chi^{(k)}(X) = 0$.
- i) [1.5 pts] Let $C_N = \frac{1}{N}$ *N* $\left(E_N^{\lambda=1,\xi=0} - E_N^{\lambda=0,\xi=0}\right)$ $\binom{N}{N}$. Is C_N expected to be equal to zero in the exact theory ? Justify your answer by considering the exact (interacting) energy decomposition used in KS-DFT.
- j) **[1.5 pts]** Deduce from Eq. (3) and questions 2. g), 2. h), and 2. i) that $I_N = -\varepsilon_N C_N +$ $\partial E_{\rm Hxc}^{\xi}[n]$ *∂ξ ξ*=0 .
- k) **[1.5 pts]** The following fundamental question naturally arises in KS-DFT: *Can we interpret the energy of the HOMO as minus the ionization potential of the true (interacting) system under study?* Explain why this question is of primary importance, for example, in the field of nanotechnology, and bring an answer in the light of question 2. j).