

Green's function-based methods in chemistry

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Today's program

• Charged excitations

- One-shot GW (G_0W_0)
- Partially self-consistent eigenvalue GW (evGW)
- Quasiparticle self-consistent GW (qsGW)
- Other self-energies (GF2, SOSEX, T-matrix, etc)

• Neutral excitations

- Random-phase approximation (RPA)
- Configuration interaction with singles (CIS)
- Time-dependent Hartree-Fock (TDHF) or RPA with exchange (RPAx)
- Time-dependent density-functional theory (TDDFT)
- Bethe-Salpeter equation (BSE) formalism

• Correlation energy

- Plasmon (or trace) formula
- Galitski-Migdal formulation
- Adiabatic connection fluctuation-dissipation theorem (ACFDT)

- 1 Context
- 2 Charged excitations
- 3 Neutral excitations
- 4 Correlation energy

Assumptions & Notations

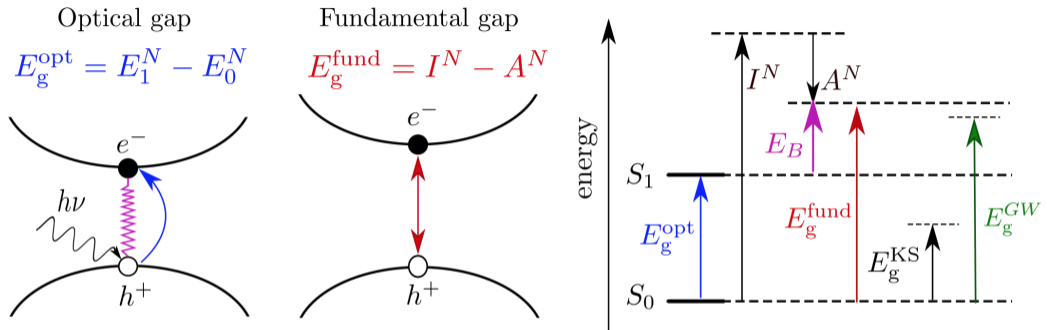
Let's talk about notations

- We consider **closed-shell systems** (2 opposite-spin electrons per orbital)
- We only deal with **singlet excited states** but **triplets** can also be obtained
- Number of **occupied orbitals** O
- Number of **vacant orbitals** V
- **Total number of orbitals** $N = O + V$
- $\phi_p(\mathbf{r})$ is a (real) **spatial orbital**
- i, j, k, l are **occupied orbitals**
- a, b, c, d are **vacant orbitals**
- p, q, r, s are **arbitrary (occupied or vacant) orbitals**
- $\mu, \nu, \lambda, \sigma$ are **basis function indexes**
- m indexes **the OV single excitations** ($i \rightarrow a$)

Useful papers/programs

- **molGW**: Bruneval et al. Comp. Phys. Comm. 208 (2016) 149
- **Turbomole**: van Setten et al. JCTC 9 (2013) 232; Kaplan et al. JCTC 12 (2016) 2528
- **Fiesta**: Blase et al. Chem. Soc. Rev. 47 (2018) 1022
- **FHI-AIMS**: Caruso et al. PRB 86 (2012) 081102
- **Reviews & Books:**
 - Reining, WIREs Comput Mol Sci 2017, e1344. doi: 10.1002/wcms.1344
 - Onida et al. Rev. Mod. Phys. 74 (2002) 601
 - Blase et al. Chem. Soc. Rev. , 47 (2018) 1022
 - Golze et al. Front. Chem. 7 (2019) 377
 - Blase et al. JPCL 11 (2020) 7371
 - Martin, Reining & Ceperley *Interacting Electrons* (Cambridge University Press)
- **GW100**: IPs for a set of 100 molecules. van Setten et al. JCTC 11 (2015) 5665 (<http://gw100.wordpress.com>)

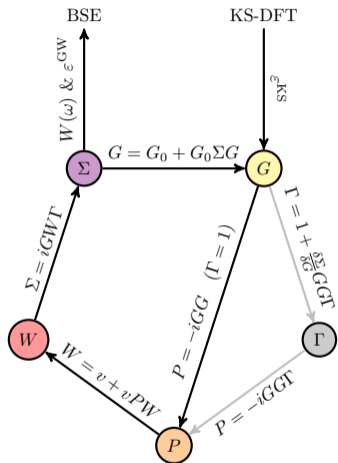
Fundamental and optical gaps (© Bruno Senjean)



$$\underbrace{E_g^{\text{KS}}}_{\text{KS gap}} = \epsilon_{\text{LUMO}}^{\text{KS}} - \epsilon_{\text{HOMO}}^{\text{KS}} \ll \underbrace{E_g^{\text{GW}}}_{\text{GW gap}} = \epsilon_{\text{LUMO}}^{\text{GW}} - \epsilon_{\text{HOMO}}^{\text{GW}} \quad (1)$$

$$\underbrace{E_g^{\text{opt}}}_{\text{optical gap}} = E_1^N - E_0^N = \underbrace{E_g^{\text{fund}}}_{\text{fundamental gap}} + \underbrace{E_B}_{\text{excitonic effect}} \quad (2)$$

Hedin's pentagon



Hedin, Phys Rev 139 (1965) A796

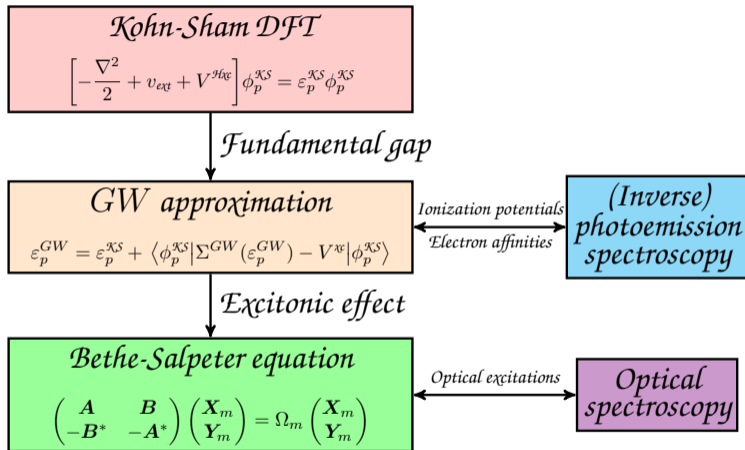
What can you calculate with *GW*?

- Ionization potentials (IPs) given by occupied MO energies
- Electron affinities (EAs) given by virtual MO energies
- Fundamental (HOMO-LUMO) gap (or band gap in solids)
- Correlation and total energies

What can you calculate with BSE?

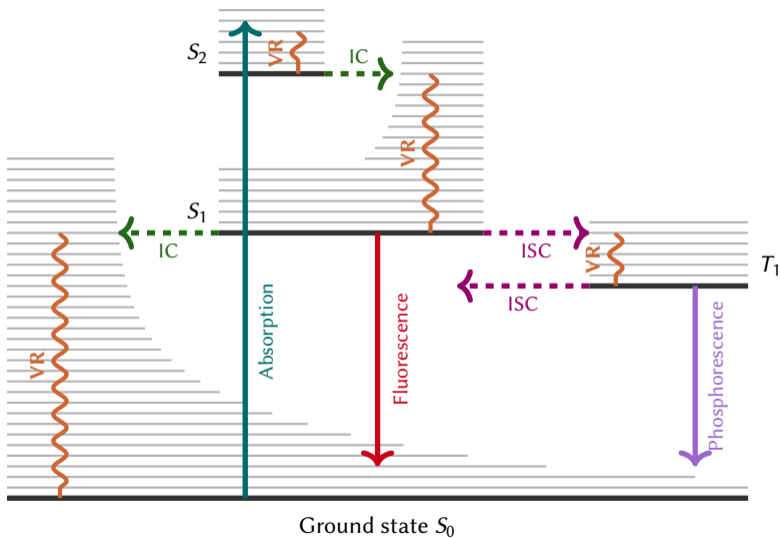
- Singlet and triplet optical excitations (vertical absorption energies)
- Oscillator strengths (absorption intensities)
- Correlation and total energies

The MBPT chain of actions

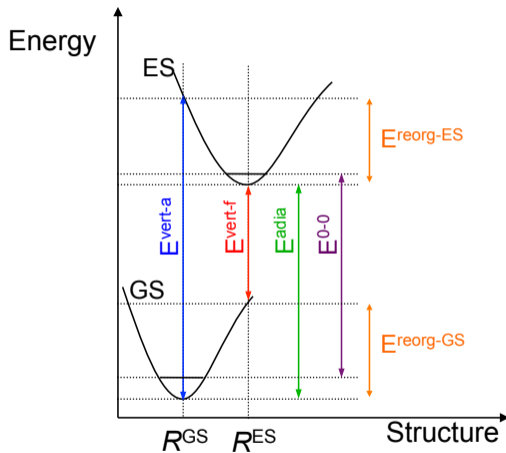


Blase et al. JPCL 11 (2020) 7371

Photochemistry: Jablonski diagram



Photochemistry: absorption, emission, and 0-0



Vertical excitation energies cannot be computed experimentally!!!

- 1 Context
- 2 Charged excitations**
- 3 Neutral excitations
- 4 Correlation energy

Green's function and dynamical screening

One-body Green's function

$$G(\mathbf{r}_1, \mathbf{r}_2; \omega) = \underbrace{\sum_i \frac{\phi_i(\mathbf{r}_1)\phi_i(\mathbf{r}_2)}{\omega - \epsilon_i - i\eta}}_{\text{removal part = IPs}} + \underbrace{\sum_a \frac{\phi_a(\mathbf{r}_1)\phi_a(\mathbf{r}_2)}{\omega - \epsilon_a + i\eta}}_{\text{addition part = EAs}} \quad (3)$$

Polarizability

$$P(\mathbf{r}_1, \mathbf{r}_2; \omega) = -\frac{i}{\pi} \int G(\mathbf{r}_1, \mathbf{r}_2; \omega + \omega') G(\mathbf{r}_1, \mathbf{r}_2; \omega') d\omega' \quad (4)$$

Dielectric function and dynamically-screened Coulomb potential

$$\epsilon(\mathbf{r}_1, \mathbf{r}_2; \omega) = \delta(\mathbf{r}_1 - \mathbf{r}_2) - \int \frac{P(\mathbf{r}_1, \mathbf{r}_3; \omega)}{|\mathbf{r}_2 - \mathbf{r}_3|} d\mathbf{r}_3 \quad (5)$$

$$W(\mathbf{r}_1, \mathbf{r}_2; \omega) = \int \frac{\epsilon^{-1}(\mathbf{r}_1, \mathbf{r}_3; \omega)}{|\mathbf{r}_2 - \mathbf{r}_3|} d\mathbf{r}_3 \quad (6)$$

Dynamical screening in the orbital basis

Spectral representation of W

$$\begin{aligned}
 W_{pq,rs}(\omega) &= \iint \phi_p(\mathbf{r}_1)\phi_q(\mathbf{r}_1) W(\mathbf{r}_1, \mathbf{r}_2; \omega) \phi_r(\mathbf{r}_2)\phi_s(\mathbf{r}_2) d\mathbf{r}_1 d\mathbf{r}_2 \\
 &= \underbrace{(pq|rs)}_{\text{(static) exchange part}} + \underbrace{2 \sum_m (pq|m)(rs|m) \left[\frac{1}{\omega - \Omega_m^{\text{RPA}} + i\eta} - \frac{1}{\omega + \Omega_m^{\text{RPA}} - i\eta} \right]}_{\text{(dynamical) correlation part } W_{pq,rs}^c(\omega)} \quad (7)
 \end{aligned}$$

Electron repulsion integrals (ERIs)

$$(pq|rs) = \iint \frac{\phi_p(\mathbf{r}_1)\phi_q(\mathbf{r}_1)\phi_r(\mathbf{r}_2)\phi_s(\mathbf{r}_2)}{|\mathbf{r}_1 - \mathbf{r}_2|} d\mathbf{r}_1 d\mathbf{r}_2 \quad (8)$$

Screened ERIs (or spectral weights)

$$(pq|m) = \sum_{ia} (pq|ia) (\mathbf{X}_m^{\text{RPA}} + \mathbf{Y}_m^{\text{RPA}})_{ia} \quad (9)$$

Computation of the dynamical screening

Direct (ph-)RPA calculation (pseudo-hermitian linear problem)

$$\begin{pmatrix} \mathbf{A}^{\text{RPA}} & \mathbf{B}^{\text{RPA}} \\ -\mathbf{B}^{\text{RPA}} & -\mathbf{A}^{\text{RPA}} \end{pmatrix} \cdot \begin{pmatrix} \mathbf{X}_m^{\text{RPA}} \\ \mathbf{Y}_m^{\text{RPA}} \end{pmatrix} = \Omega_m^{\text{RPA}} \begin{pmatrix} \mathbf{X}_m^{\text{RPA}} \\ \mathbf{Y}_m^{\text{RPA}} \end{pmatrix} \quad (10)$$

For singlet states: $A_{ia,jb}^{\text{RPA}} = \delta_{ij}\delta_{ab}(\epsilon_a - \epsilon_i) + 2(ia|bj)$ $B_{ia,jb}^{\text{RPA}} = 2(ia|jb)$ (11)

Non-hermitian to hermitian

$$(\mathbf{A} - \mathbf{B})^{1/2} \cdot (\mathbf{A} + \mathbf{B}) \cdot (\mathbf{A} - \mathbf{B})^{1/2} \cdot \mathbf{Z}_m = \Omega_m^2 \mathbf{Z}_m \quad (12)$$

$$(\mathbf{X}_m + \mathbf{Y}_m) = \Omega_m^{-1/2} (\mathbf{A} - \mathbf{B})^{+1/2} \cdot \mathbf{Z}_m \quad (13)$$

$$(\mathbf{X}_m - \mathbf{Y}_m) = \Omega_m^{+1/2} (\mathbf{A} - \mathbf{B})^{-1/2} \cdot \mathbf{Z}_m \quad (14)$$

Tamm-Dancoff approximation (TDA)

$$\mathbf{B} = \mathbf{0} \quad \Rightarrow \quad \mathbf{A} \cdot \mathbf{X}_m = \Omega_m^{\text{TDA}} \mathbf{X}_m \quad (15)$$

The self-energy

GW Self-energy

$$\underbrace{\Sigma^{\text{xc}}(\mathbf{r}_1, \mathbf{r}_2; \omega)}_{\text{GW self-energy}} = \underbrace{\Sigma^{\text{x}}(\mathbf{r}_1, \mathbf{r}_2)}_{\text{exchange}} + \underbrace{\Sigma^{\text{c}}(\mathbf{r}_1, \mathbf{r}_2; \omega)}_{\text{correlation}} = \frac{i}{2\pi} \int G(\mathbf{r}_1, \mathbf{r}_2; \omega + \omega') W(\mathbf{r}_1, \mathbf{r}_2; \omega') e^{i\eta\omega'} d\omega' \quad (16)$$

Exchange part of the (static) self-energy

$$\Sigma_{pq}^{\text{x}} = - \sum_i (pi|i q) \quad (17)$$

Correlation part of the (dynamical) self-energy

$$\Sigma_{pq}^{\text{c}}(\omega) = 2 \sum_{im} \frac{(pi|m)(qi|m)}{\omega - \epsilon_i + \Omega_m^{\text{RPA}} - i\eta} + 2 \sum_{am} \frac{(pa|m)(qa|m)}{\omega - \epsilon_a - \Omega_m^{\text{RPA}} + i\eta} \quad (18)$$

Quasiparticle equation

Dyson equation

$$[G(\mathbf{r}_1, \mathbf{r}_2; \omega)]^{-1} = \underbrace{[G_{\text{KS}}(\mathbf{r}_1, \mathbf{r}_2; \omega)]^{-1}}_{\text{KS Green's function}} + \underbrace{\Sigma^{\text{xc}}(\mathbf{r}_1, \mathbf{r}_2; \omega) - v^{\text{xc}}(\mathbf{r}_1)}_{\text{KS potential}} \delta(\mathbf{r}_1 - \mathbf{r}_2) \quad (19)$$

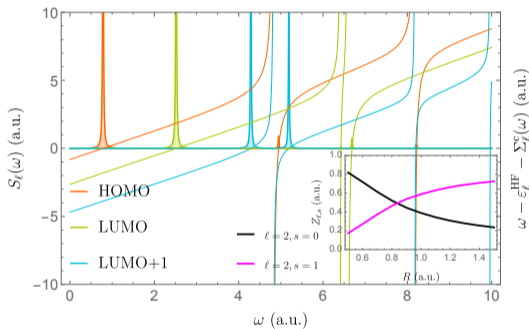
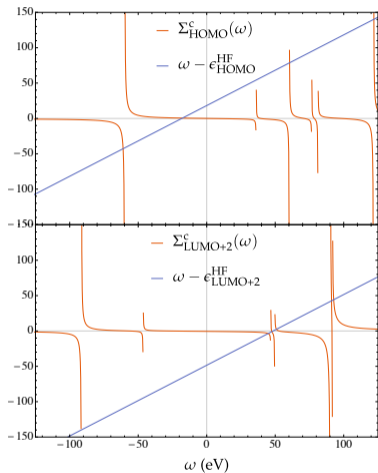
Non-linear quasiparticle (QP) equation

$$\omega = \epsilon_p^{\text{KS}} + \Sigma_{pp}^{\text{xc}}(\omega) - V_p^{\text{xc}} \quad \text{with} \quad V_p^{\text{xc}} = \int \phi_p(\mathbf{r}) v^{\text{xc}}(\mathbf{r}) \phi_p(\mathbf{r}) d\mathbf{r} \quad (20)$$

Linearized QP equation

$$\Sigma_{pp}^{\text{xc}}(\omega) \approx \Sigma_{pp}^{\text{xc}}(\epsilon_p^{\text{KS}}) + (\omega - \epsilon_p^{\text{KS}}) \left. \frac{\partial \Sigma_{pp}^{\text{xc}}(\omega)}{\partial \omega} \right|_{\omega=\epsilon_p^{\text{KS}}} \Rightarrow \epsilon_p^{\text{GW}} = \epsilon_p^{\text{KS}} + Z_p [\Sigma_{pp}^{\text{xc}}(\epsilon_p^{\text{KS}}) - V_p^{\text{xc}}] \quad (21)$$

$$\underbrace{Z_p}_{\text{renormalization factor}} = \left[1 - \left. \frac{\partial \Sigma_{pp}^{\text{xc}}(\omega)}{\partial \omega} \right|_{\omega=\epsilon_p^{\text{KS}}} \right]^{-1} \quad \text{with} \quad 0 \leq Z_p \leq 1 \quad (22)$$

Solutions of the non-linear QP equation: $evGW@HF/6-31G$ for H_2 at $R = 1$ bohr

Loos et al, JCTC 14 (2018) 3071

Vénil et al, JCTC 14 (2018) 5220

Acronyms

- perturbative GW , one-shot GW , or G_0W_0
- $evGW$ or eigenvalue-only (partially) self-consistent GW
- $qsGW$ or quasiparticle (partially) self-consistent GW
- $scGW$ or (fully) self-consistent GW

Perturbative GW with linearized solution

procedure G_0W_0 LIN@KS

Perform KS calculation to get ϵ^{KS} , \mathbf{c}^{KS} , and \mathbf{V}^{xc}

AO to MO transformation for ERIs: $(\mu\nu|\lambda\sigma) \xrightarrow{\mathbf{c}^{\text{KS}}} (pq|rs)$

Construct RPA matrices \mathbf{A}^{RPA} and \mathbf{B}^{RPA} from ϵ^{KS} and $(pq|rs)$

Compute RPA eigenvalues Ω^{RPA} and eigenvectors $\mathbf{X}^{\text{RPA}} + \mathbf{Y}^{\text{RPA}}$

▷ This is a $\mathcal{O}(N^6)$ step!

Form screened ERIs $(pq|m)$

for $p = 1, \dots, N$ do

 Compute diagonal of the self-energy $\Sigma_{pp}^c(\omega)$ at $\omega = \epsilon_p^{\text{KS}}$

 Compute renormalization factors Z_p

 Evaluate $\epsilon_p^{G_0W_0} = \epsilon_p^{\text{KS}} + Z_p \left\{ \text{Re}[\Sigma_{pp}^c(\epsilon_p^{\text{KS}})] - V_p^{\text{xc}} \right\}$

end for

end procedure

For contour deformation technique, see, for example, Duchemin & Blase, JCTC 16 (2020) 1742

Example from QuAcK (Ne/cc-pVDZ)

One-shot G0W0 calculation			Linearized G0W0 subroutine		
Iter #	Frame	e_HF (eV)	Sig_c (eV)	Procedure G0WZ LIN	e_QP (eV)
1		-891.591504	18.364427	0.859504	-875.807142
2		-52.218791	4.035435	0.956042	-48.360659
3		-22.647397	1.832273	0.965238	-20.878718
4		-22.647397	1.832273	0.965238	-20.878718
5		-22.647397	1.832273	0.965238	-20.878718
6		46.107752	-0.820124	0.982086	45.302383
7		46.107752	-0.820124	0.982086	45.302383
8		46.107752	-0.820124	0.982086	45.302383
9		54.167043	-1.061182	0.985754	53.121001
10		141.402085	-2.617768	0.898641	139.049684
11		141.402085	-2.617768	0.898641	139.049684
12		141.402085	-2.617768	0.898641	139.049684
13		141.402085	-2.617768	0.898641	139.049684
14		141.402085	-2.617768	0.898641	139.049684
15		282.545807	-3.872629	0.944019	278.890026

G0W0 HOMO energy:	-20.878718 eV
G0W0 LUMO energy:	45.302383 eV
G0W0 HOMO-LUMO gap :	66.181102 eV

RPA@G0W0 total energy :	-128.714946 au
RPA@G0W0 correlation energy:	-0.226138 au
GM@G0W0 total energy :	-128.887856 au
GM@G0W0 correlation energy:	-0.399048 au

<https://github.com/pfloos/QuAcK>

Perturbative GW with graphical solution

procedure G_0W_0 GRAPH@KS

Perform KS calculation to get ϵ^{KS} , \mathbf{c}^{KS} , and \mathbf{V}^{xc}

AO to MO transformation for ERIs: $(\mu\nu|\lambda\sigma) \xrightarrow{\mathbf{c}^{\text{KS}}} (pq|rs)$

Construct RPA matrices \mathbf{A}^{RPA} and \mathbf{B}^{RPA} from ϵ^{KS} and $(pq|rs)$

Compute RPA eigenvalues Ω^{RPA} and eigenvectors $\mathbf{X}^{\text{RPA}} + \mathbf{Y}^{\text{RPA}}$

▷ This is a $\mathcal{O}(N^6)$ step!

Form screened ERIs $(pq|m)$

for $p = 1, \dots, N$ **do**

 Compute diagonal of the self-energy $\Sigma_{pp}^{\text{c}}(\omega)$

 Solve $\omega = \epsilon_p^{\text{KS}} + \text{Re}[\Sigma_{pp}^{\text{c}}(\omega)] - V_p^{\text{xc}}$ to get $\epsilon_p^{G_0W_0}$ via Newton's method

end for

end procedure

Newton's method

https://en.wikipedia.org/wiki/Newton%27s_method

Partially self-consistent eigenvalue GW

procedure EVGW@KS

Perform KS calculation to get ϵ^{KS} , \mathbf{c}^{KS} , and \mathbf{V}^{xc}

AO to MO transformation for ERIs: $(\mu\nu|\lambda\sigma) \xrightarrow{\mathbf{c}^{\text{KS}}} (pq|rs)$

Set $\epsilon^{G_{-1}W_{-1}} = \epsilon^{\text{KS}}$ and $n = 0$

while $\max |\Delta| > \tau$ **do**

Construct RPA matrices \mathbf{A}^{RPA} and \mathbf{B}^{RPA} from $\epsilon^{G_{n-1}W_{n-1}}$ and $(pq|rs)$

Compute RPA eigenvalues Ω^{RPA} and eigenvectors $\mathbf{X}^{\text{RPA}} + \mathbf{Y}^{\text{RPA}}$

Form screened ERIs $(pq|m)$

for $p = 1, \dots, N$ **do**

Compute diagonal of the self-energy $\Sigma_{pp}^{\text{c}}(\omega)$

Solve $\omega = \epsilon_p^{\text{KS}} + \text{Re}[\Sigma_{pp}^{\text{c}}(\omega)] - V_p^{\text{xc}}$ to get $\epsilon_p^{G_n W_n}$

end for

$\Delta = \epsilon^{G_n W_n} - \epsilon^{G_{n-1} W_{n-1}}$

$n \leftarrow n + 1$

end while

end procedure

▷ This is a $\mathcal{O}(N^6)$ step!

Example from QuAcK (Ne/cc-pVDZ)

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-----
Self-consistent evG8W8 calculation
-----
# | e_HF (eV) | Sigma_c (eV) | Z | e_QP (eV) |
-----
1 | -891.591504 | 18.746313 | 0.853211 | -872.845115 |
2 | -52.218791 | 4.097592 | 0.954012 | -48.121107 |
3 | -22.647397 | 1.872062 | 0.963351 | -20.775232 |
4 | -22.647397 | 1.872062 | 0.963351 | -20.775232 |
5 | -22.647397 | 1.872062 | 0.963351 | -20.775232 |
6 | 46.107752 | -0.834752 | 0.981106 | 45.273065 |
7 | 46.107752 | -0.834752 | 0.981106 | 45.273065 |
8 | 46.107752 | -0.834752 | 0.981106 | 45.273065 |
9 | 54.167043 | -1.078523 | 0.984963 | 53.088542 |
10 | 141.402085 | -3.068193 | 0.763837 | 138.333929 |
11 | 141.402085 | -3.068193 | 0.763837 | 138.333929 |
12 | 141.402085 | -3.068193 | 0.763837 | 138.333929 |
13 | 141.402085 | -3.068193 | 0.763837 | 138.333929 |
14 | 141.402085 | -3.068193 | 0.763837 | 138.333929 |
15 | 282.545807 | -4.009519 | 0.941599 | 278.536345 |
-----
Iteration 8
Convergence = 0.00000
-----
evGW HOMO energy: -20.775232 eV
evGW LUMO energy: 45.273065 eV
evGW HOMO-LUMO gap : 66.048297 eV
-----
RPA@evGW total energy : -128.715585 au
RPA@evGW correlation energy: -0.226777 au
GM@evGW total energy : -128.898601 au
GM@evGW correlation energy: -0.409794 au
-----

```

<https://github.com/pfloos/QuAcK>

Quasiparticle self-consistent GW (qsGW)

procedure qsGW

Perform HF calculation to get ϵ^{HF} and \mathbf{c}^{HF} (optional)

Set $\epsilon^{G_{-1}W_{-1}} = \epsilon^{\text{HF}}$, $\mathbf{c}^{G_{-1}W_{-1}} = \mathbf{c}^{\text{HF}}$ and $n = 0$

while $\max|\Delta| > \tau$ **do**

AO to MO transformation for ERIs: $(\mu\nu|\lambda\sigma) \xrightarrow{\mathbf{c}^{G_{n-1}W_{n-1}}} (pq|rs)$

▷ This is a $\mathcal{O}(N^5)$ step!

Construct RPA matrices \mathbf{A}^{RPA} and \mathbf{B}^{RPA} from $\epsilon^{G_{n-1}W_{n-1}}$ and $(pq|rs)$

Compute RPA eigenvalues Ω^{RPA} and eigenvectors $\mathbf{X}^{\text{RPA}} + \mathbf{Y}^{\text{RPA}}$

▷ This is a $\mathcal{O}(N^6)$ step!

Form screened ERIs $(pq|m)$

Evaluate $\Sigma^{\text{c}}(\epsilon^{G_{n-1}W_{n-1}})$ and form $\tilde{\Sigma}^{\text{c}} \leftarrow [\Sigma^{\text{c}}(\epsilon^{G_{n-1}W_{n-1}})^{\dagger} + \Sigma^{\text{c}}(\epsilon^{G_{n-1}W_{n-1}})]/2$

Form \mathbf{F}^{HF} from $\mathbf{c}^{G_{n-1}W_{n-1}}$ and then $\tilde{\mathbf{F}} = \mathbf{F}^{\text{HF}} + \tilde{\Sigma}^{\text{c}}$

Diagonalize $\tilde{\mathbf{F}}$ to get $\epsilon^{G_nW_n}$ and $\mathbf{c}^{G_nW_n}$

$\Delta = \epsilon^{G_nW_n} - \epsilon^{G_{n-1}W_{n-1}}$

$n \leftarrow n + 1$

end while

end procedure

Example from QuAcK (Ne/cc-pVDZ)

```

Self-consistent qsG16W16 calculation
-----
# | e_HF (eV) | Sig_c (eV) | Z | e_QP (eV) |
-----
1 | -891.591504 | 18.755754 | 0.853363 | -873.652325 |
2 | -52.218791 | 4.058060 | 0.954380 | -48.405559 |
3 | -22.647397 | 1.855512 | 0.963520 | -21.066156 |
4 | -22.647397 | 1.855512 | 0.963520 | -21.066156 |
5 | -22.647397 | 1.855512 | 0.963520 | -21.066156 |
6 | 46.107752 | -0.848683 | 0.980977 | 45.067534 |
7 | 46.107752 | -0.848683 | 0.980977 | 45.067534 |
8 | 46.107752 | -0.848683 | 0.980977 | 45.067534 |
9 | 54.167043 | -1.102700 | 0.984676 | 52.926661 |
10 | 141.402085 | -3.043127 | 0.776916 | 138.069002 |
11 | 141.402085 | -3.043127 | 0.776916 | 138.069002 |
12 | 141.402085 | -3.043127 | 0.776916 | 138.069002 |
13 | 141.402085 | -3.043127 | 0.776916 | 138.069002 |
14 | 141.402085 | -3.043127 | 0.776916 | 138.069002 |
15 | 282.545807 | -3.998794 | 0.941677 | 278.156200 |
-----
Iteration 16
Convergence = 0.00001

qsGW HOMO energy: -21.066156 eV
qsGW LUMO energy: 45.067534 eV
qsGW HOMO-LUMO gap : 66.133690 eV

qsGW total energy: -128.488468 au
qsGW exchange energy: -12.101095 au
GM@qsGW correlation energy: -0.410249 au
RPA@qsGW correlation energy: -0.227077 au

```

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Summary
-----
One-electron energy: -182.4760110151 au
Kinetic energy: 128.2215634186 au
Potential energy: -310.6975744337 au
-----
Two-electron energy: 53.9875434022 au
Hartree energy: 66.0886388591 au
Exchange energy: -12.1010954570 au
Correlation energy: -0.4102491313 au
-----
Electronic energy: -128.4884676130 au
Nuclear repulsion: 0.0000000000 au
qsGW energy: -128.4884676130 au
-----
Dipole moment (Debye)
begin:center X Y Z Tot.
\ 0.000000 0.000000 0.000000 0.000000
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```

<https://github.com/pfloos/QuAcK>

Other self-energies

Second-order Green's function (GF2) [Hirata et al. JCP 147 (2017) 044108]

$$\Sigma_{pq}^{\text{GF2}}(\omega) = \frac{1}{2} \sum_{iab} \frac{\langle iq||ab\rangle \langle ab||ip\rangle}{\omega + \epsilon_i - \epsilon_a - \epsilon_b} + \frac{1}{2} \sum_{ija} \frac{\langle aq||ij\rangle \langle ij||ap\rangle}{\omega + \epsilon_a - \epsilon_i - \epsilon_j} \quad (23)$$

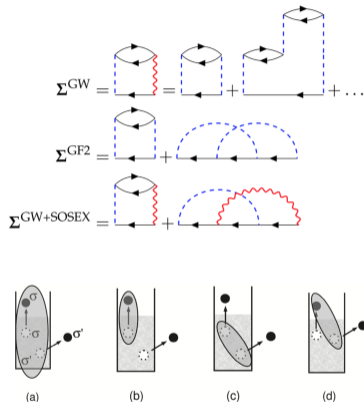
T-matrix [Romaniello et al. PRB 85 (2012) 155131; Zhang et al. JPCL 8 (2017) 3223]

$$\Sigma_{pq}^{\text{GT}}(\omega) = \sum_{im} \frac{\langle pi|\chi_m^{N+2}\rangle \langle qi|\chi_m^{N+2}\rangle}{\omega + \epsilon_i - \Omega_m^{N+2}} + \sum_{am} \frac{\langle pa|\chi_m^{N-2}\rangle \langle qa|\chi_m^{N-2}\rangle}{\omega + \epsilon_i - \Omega_m^{N-2}} \quad (24)$$

$$\langle pi|\chi_m^{N+2}\rangle = \sum_{c<d} \langle pi||cd\rangle \chi_{cd}^{N+2,m} + \sum_{k<l} \langle pi||kl\rangle \gamma_{kl}^{N+2,m} \quad (25)$$

$$\langle pa|\chi_m^{N-2}\rangle = \sum_{c<d} \langle pa||cd\rangle \chi_{cd}^{N-2,m} + \sum_{k<l} \langle pa||kl\rangle \gamma_{kl}^{N-2,m} \quad (26)$$

pp-RPA problem:
$$\begin{pmatrix} \mathbf{A} & \mathbf{B} \\ -\mathbf{B}^\top & -\mathbf{C} \end{pmatrix} \cdot \begin{pmatrix} \mathbf{X}_m^{N\pm 2} \\ \mathbf{Y}_m^{N\pm 2} \end{pmatrix} = \Omega_m^{N\pm 2} \begin{pmatrix} \mathbf{X}_m^{N\pm 2} \\ \mathbf{Y}_m^{N\pm 2} \end{pmatrix} \quad (27)$$



Martin, Reining & Ceperley, Interacting Electrons (Cambridge University Press)

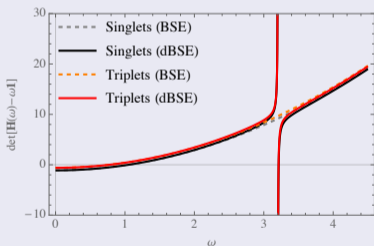
- 1 Context
- 2 Charged excitations
- 3 Neutral excitations**
- 4 Correlation energy

Dynamical vs static kernels

A non-linear BSE problem [Strinati, Riv. Nuovo Cimento 11 (1988) 1]

$$\begin{pmatrix} \mathbf{A}(\omega) & \mathbf{B}(\omega) \\ -\mathbf{B}(-\omega) & -\mathbf{A}(-\omega) \end{pmatrix} \cdot \begin{pmatrix} \mathbf{X} \\ \mathbf{Y} \end{pmatrix} = \omega \begin{pmatrix} \mathbf{X} \\ \mathbf{Y} \end{pmatrix} \quad \text{Hard to solve!} \quad (28)$$

Static BSE vs dynamic BSE for HeH⁺/STO-3G



Dynamical kernels can give you more than static kernels... Sometimes, too much...

Authier & Loos, JCP 153 (2020) 184105 [see also Romaniello et al, JCP 130 (2009) 044108]

Löwdin partitioning technique

Folding or dressing process

$$\underbrace{H \cdot c = \omega c}_{\text{A large linear system with } N \text{ solutions...}} \Rightarrow \begin{pmatrix} \overbrace{H_1}^{N_1 \times N_1} & h^\top \\ h & \underbrace{H_2}_{N_2 \times N_2} \end{pmatrix} \cdot \begin{pmatrix} c_1 \\ c_2 \end{pmatrix} = \omega \begin{pmatrix} c_1 \\ c_2 \end{pmatrix} \quad N = N_1 + N_2 \quad (29)$$

$$\text{Row \#2: } h \cdot c_1 + H_2 \cdot c_2 = \omega c_2 \quad \Rightarrow \quad c_2 = (\omega \mathbf{1} - H_2)^{-1} \cdot h \cdot c_1 \quad (30)$$

$$\text{Row \#1: } H_1 \cdot c_1 + h^\top \cdot c_2 = \omega c_1 \quad \Rightarrow \quad \underbrace{\tilde{H}_1(\omega) \cdot c_1 = \omega c_1}_{\text{A smaller non-linear system with } N \text{ solutions...}} \quad (31)$$

$$\underbrace{\tilde{H}_1(\omega)}_{\text{Effective Hamiltonian}} = H_1 + h^\top \cdot (\omega \mathbf{1} - H_2)^{-1} \cdot h \quad (32)$$

$$\text{Static approx. (e.g. } \omega = 0\text{):} \quad \underbrace{\tilde{H}_1(\omega = 0)}_{\text{A smaller linear system with } N_1 \text{ solutions...}} = H_1 - \underbrace{h^\top \cdot H_2^{-1} \cdot h}_{\text{approximations possible...}} \quad (33)$$

TD-DFT and BSE in practice: Casida-like equations

Linear response problem

$$\begin{pmatrix} \mathbf{A} & \mathbf{B} \\ -\mathbf{B} & -\mathbf{A} \end{pmatrix} \cdot \begin{pmatrix} \mathbf{X}_m \\ \mathbf{Y}_m \end{pmatrix} = \Omega_m \begin{pmatrix} \mathbf{X}_m \\ \mathbf{Y}_m \end{pmatrix}$$

Blue pill: TD-DFT within the **adiabatic** approximation

$$A_{ia,jb} = \left(\epsilon_a^{\text{KS}} - \epsilon_i^{\text{KS}} \right) \delta_{ij} \delta_{ab} + 2(ia|bj) + f_{ia,bj}^{\text{xc}} \quad B_{ia,jb} = 2(ia|jb) + f_{ia,jb}^{\text{xc}} \quad (34)$$

$$f_{ia,bj}^{\text{xc}} = \iint \phi_i(\mathbf{r}) \phi_a(\mathbf{r}) \frac{\delta^2 E^{\text{xc}}}{\delta \rho(\mathbf{r}) \delta \rho(\mathbf{r}')} \phi_b(\mathbf{r}) \phi_j(\mathbf{r}) d\mathbf{r} d\mathbf{r}' \quad (35)$$

Red pill: BSE within the **static** approximation

$$A_{ia,jb} = \left(\epsilon_a^{\text{GW}} - \epsilon_i^{\text{GW}} \right) \delta_{ij} \delta_{ab} + 2(ia|bj) - W_{ij,ba}^{\text{stat}} \quad B_{ia,jb} = 2(ia|jb) - W_{ib,ja}^{\text{stat}} \quad (36)$$

$$W_{ij,ab}^{\text{stat}} \equiv W_{ij,ab}(\omega = 0) = (ij|ab) - W_{ij,ab}^{\text{c}}(\omega = 0) \quad (37)$$

The bridge between TD-DFT and BSE

TD-DFT	Connection	BSE
One-point density $\rho(1)$	$\rho(1) = -iG(11^+)$	Two-point Green's function $G(12)$
Two-point susceptibility $\chi(12) = \frac{\partial \rho(1)}{\partial U(2)}$	$\chi(12) = -iL(12; 1^+2^+)$	Four-point susceptibility $L(12; 34) = \frac{\partial G(13)}{\partial U(42)}$
Two-point kernel $K(12) = v(12) + \frac{\partial v^{xc}(1)}{\partial \rho(2)}$		Four-point kernel $i\Xi(1234) = v(13)\delta(12)\delta(34) - \frac{\partial \Sigma^{xc}(12)}{\partial G(34)}$

For dynamical correction within BSE, see, for example, [Loos & Blase, JCP 153 \(2020\) 114120](#)

BSE in a computer

Vertical excitation energies from BSE

procedure BSE@GW

Compute GW quasiparticle energies ϵ_p^{GW} at the G_0W_0 , $evGW$, or $qsGW$ level

Compute static screening $W_{pq,rs}^{stat}$

Construct BSE matrices A^{BSE} and B^{BSE} from ϵ_p^{GW} , $(pq|rs)$, and $W_{pq,rs}^{stat}$

Compute lowest eigenvalues Ω_m^{BSE} and eigenvectors $X_m^{BSE} + Y_m^{BSE}$ (optional) ▷ This is a $\mathcal{O}(N^4)$ step!

end procedure

Removing the correlation part: TDHF and CIS

Linear response problem

$$\begin{pmatrix} \mathbf{A} & \mathbf{B} \\ -\mathbf{B} & -\mathbf{A} \end{pmatrix} \cdot \begin{pmatrix} \mathbf{X}_m \\ \mathbf{Y}_m \end{pmatrix} = \Omega_m \begin{pmatrix} \mathbf{X}_m \\ \mathbf{Y}_m \end{pmatrix}$$

TDHF = RPA with exchange (RPAx)

$$A_{ia,jb} = \left(\epsilon_a^{\text{HF}} - \epsilon_i^{\text{HF}} \right) \delta_{ij} \delta_{ab} + 2(ia|bj) - (ij|ba) \quad B_{ia,jb} = 2(ia|jb) - (ib|ja) \quad (38)$$

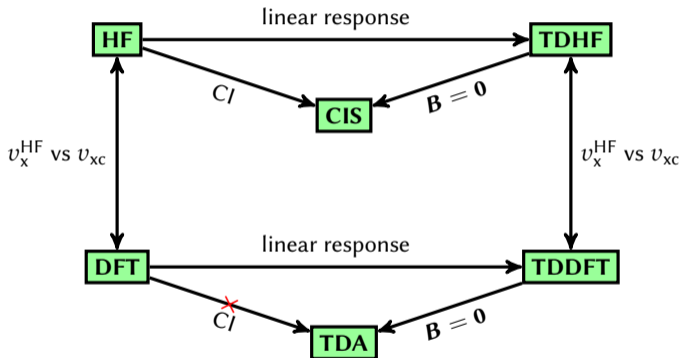
Linear response problem within the Tamm-Dancoff approximation

$$\mathbf{A} \cdot \mathbf{X}_m = \Omega_m \mathbf{X}_m \quad (39)$$

TDHF within TDA = CIS

$$A_{ia,jb} = \left(\epsilon_a^{\text{HF}} - \epsilon_i^{\text{HF}} \right) \delta_{ij} \delta_{ab} + 2(ia|bj) - (ij|ba) \quad (40)$$

Relationship between CIS, TDHF, DFT and TDDFT



Linear response

General linear response problem

procedure LINEAR RESPONSE

Compute \mathbf{A} matrix at a given level of theory (RPA, RPAx, TD-DFT, BSE, etc)

if TDA then

Diagonalize \mathbf{A} to get Ω_m^{TDA} and $\mathbf{X}_m^{\text{TDA}}$

else

Compute \mathbf{B} matrix at a given level of theory

Diagonalize $\mathbf{A} - \mathbf{B}$ to form $(\mathbf{A} - \mathbf{B})^{1/2}$

Form and diagonalize $(\mathbf{A} - \mathbf{B})^{1/2} \cdot (\mathbf{A} + \mathbf{B}) \cdot (\mathbf{A} - \mathbf{B})^{1/2}$ to get Ω_m^2 and \mathbf{Z}_m

Compute $\sqrt{\Omega_m^2}$ and $(\mathbf{X}_m + \mathbf{Y}_m) = \Omega_m^{-1/2} (\mathbf{A} - \mathbf{B})^{1/2} \cdot \mathbf{Z}_m$

end if

end procedure

Form linear response matrices

Linear-response matrices for BSE

procedure FORM **A** FOR SINGLET STATES

Set **A** = **0**

$ia \leftarrow 0$

for $i = 1, \dots, O$ **do**

for $a = 1, \dots, V$ **do**

$ia \leftarrow ia + 1$

$jb \leftarrow 0$

for $j = 1, \dots, O$ **do**

for $b = 1, \dots, V$ **do**

$jb \leftarrow jb + 1$

$$A_{ia,jb} = \delta_{ij}\delta_{ab}(\epsilon_a^{\text{GW}} - \epsilon_i^{\text{GW}}) + 2(ia|bj) - (ij|ba) + W_{ij,ba}^c(\omega = 0)$$

end for

end for

end for

end for

end procedure

Properties

Oscillator strength (length gauge)

$$f_m = \frac{2}{3} \Omega_m [(\mu_m^x)^2 + (\mu_m^y)^2 + (\mu_m^z)^2] \quad (41)$$

Transition dipole

$$\mu_m^x = \sum_{ia} (i|x|a)(X_m + Y_m)_{ia} \quad (p|x|q) = \int \phi_p(\mathbf{r}) x \phi_q(\mathbf{r}) d\mathbf{r} \quad (42)$$

Monitoring possible spin contamination [Monino & Loos, JCTC 17 (2021) 2852]

$$\langle \hat{S}^2 \rangle_m = \langle \hat{S}^2 \rangle_0 + \underbrace{\Delta \langle \hat{S}^2 \rangle_m}_{\text{JCP 134101 (2011) 134}} \quad \langle \hat{S}^2 \rangle_0 = \frac{n_\alpha - n_\beta}{2} \left(\frac{n_\alpha - n_\beta}{2} + 1 \right) + n_\beta + \sum_p (p_\alpha | p_\beta) \quad (43)$$

Example from QuAcK (H₂O/cc-pVDZ)

```

-----
Excitation n. 1: 8.411378 eV f = 0.0255 <S**2> = 0.0000
-----
5 -> 6 = 0.704168
-----
Excitation n. 2: 10.496539 eV f = 0.0000 <S**2> = 0.0000
-----
5 -> 7 = 0.699391
5 -> 8 = -0.095559
-----
Excitation n. 3: 11.080888 eV f = 0.0924 <S**2> = 0.0000
-----
4 -> 6 = -0.703496
-----
Excitation n. 4: 13.165908 eV f = 0.0706 <S**2> = 0.0000
-----
4 -> 7 = 0.701946
-----
Excitation n. 5: 14.913736 eV f = 0.2678 <S**2> = 0.0000
-----
3 -> 6 = 0.704100

```

```

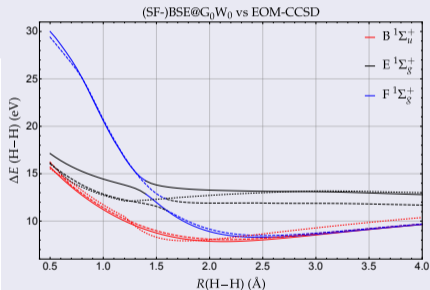
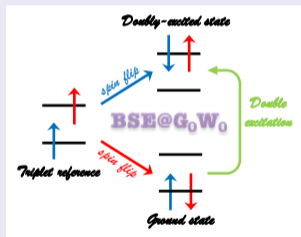
-----
Excitation n. 1: 7.632804 eV f = 0.0000 <S**2> = 2.0000
-----
5 -> 6 = 0.700599
5 -> 9 = -0.089914
-----
Excitation n. 2: 9.897068 eV f = 0.0000 <S**2> = 2.0000
-----
4 -> 6 = -0.695522
4 -> 9 = 0.093664
-----
Excitation n. 3: 10.002114 eV f = 0.0000 <S**2> = 2.0000
-----
5 -> 7 = 0.695328
5 -> 8 = -0.117774
-----
Excitation n. 4: 11.995497 eV f = 0.0000 <S**2> = 2.0000
-----
3 -> 6 = 0.228354
4 -> 7 = 0.651412
4 -> 8 = -0.135998
-----
Excitation n. 5: 13.698483 eV f = 0.0000 <S**2> = 2.0000
-----
3 -> 6 = -0.656938
3 -> 9 = 0.101160
4 -> 7 = 0.234306

```

<https://github.com/pfloos/QuAcK>

Open-shell systems and double excitations

Spin-flip formalism (H2/cc-pVQZ)



```

Excitation n. 1: -4.891498 eV f = 0.0000 <S**2> = 0.0217
-----
1A -> 2B = 0.111265
1A -> 4B = -0.121073
2A -> 1B = 0.961211
2A -> 3B = -0.212041

Excitation n. 2: 0.691826 eV f = 0.0000 <S**2> = 1.9964
-----
1A -> 1B = -0.680242
1A -> 3B = 0.202252 0.70
2A -> 2B = -0.590377
2A -> 4B = 0.373424

Excitation n. 3: 5.625694 eV f = 0.0000 <S**2> = 0.1795
-----
1A -> 1B = -0.617840
1A -> 3B = 0.196687
2A -> 2B = 0.753811

Excitation n. 4: 7.474558 eV f = 0.0000 <S**2> = 0.9821
-----
1A -> 4B = -0.111548
2A -> 1B = -0.231266
2A -> 3B = -0.960135
  
```

Monino & Loos, JCTC 17 (2021) 2852

- 1 Context
- 2 Charged excitations
- 3 Neutral excitations
- 4 Correlation energy

Correlation energy at the *GW* or BSE levelRPA@*GW* correlation energy: plasmon (or trace) formula

$$E_c^{\text{RPA}} = \frac{1}{2} \left[\sum_p \Omega_m^{\text{RPA}} - \text{Tr}(\mathbf{A}^{\text{RPA}}) \right] = \frac{1}{2} \sum_m (\Omega_m^{\text{RPA}} - \Omega_m^{\text{TDA}})$$

Galitskii-Migdal functional

$$E_c^{\text{GM}} = \frac{-i}{2} \sum_{pq} \int \frac{d\omega}{2\pi} \Sigma_{pq}^c(\omega) G_{pq}(\omega) e^{i\omega\eta} = 4 \sum_{ia} \sum_m \frac{(ai|m)^2}{\epsilon_a^{\text{GW}} - \epsilon_i^{\text{GW}} + \Omega_m^{\text{RPA}}}$$

ACFDT@BSE@*GW* correlation energy from the adiabatic connection

$$E_c^{\text{ACFDT}} = \frac{1}{2} \int_0^1 \text{Tr}(\mathbf{K} \mathbf{P}^\lambda) d\lambda \quad (44)$$

Adiabatic connection fluctuation dissipation theorem (ACFDT)

Adiabatic connection

$$E_c^{\text{ACFDT}} = \frac{1}{2} \int_0^1 \text{Tr}(\mathbf{K} \mathbf{P}^\lambda) d\lambda \stackrel{\text{quad}}{\approx} \frac{1}{2} \sum_{k=1}^K w_k \text{Tr}(\mathbf{K} \mathbf{P}^{\lambda_k}) \quad (45)$$

λ is the **strength** of the electron-electron interaction:

- $\lambda = 0$ for the **non-interacting system**
- $\lambda = 1$ for the **physical system**

Interaction kernel

$$\mathbf{K} = \begin{pmatrix} \tilde{\mathbf{A}} & \tilde{\mathbf{B}} \\ \tilde{\mathbf{B}} & \tilde{\mathbf{A}} \end{pmatrix} \quad \tilde{A}_{ia,jb} = 2(ia|bj) \quad \tilde{B}_{ia,jb} = 2(ia|jb) \quad (46)$$

Correlation part of the two-particle density matrix

$$\mathbf{P}^\lambda = \begin{pmatrix} \mathbf{Y}^\lambda \cdot (\mathbf{Y}^\lambda)^\top & \mathbf{Y}^\lambda \cdot (\mathbf{X}^\lambda)^\top \\ \mathbf{X}^\lambda \cdot (\mathbf{Y}^\lambda)^\top & \mathbf{X}^\lambda \cdot (\mathbf{X}^\lambda)^\top \end{pmatrix} - \begin{pmatrix} \mathbf{0} & \mathbf{0} \\ \mathbf{0} & \mathbf{1} \end{pmatrix} \quad (47)$$

Gaussian quadrature

Numerical integration by quadrature

“A K -point *Gaussian quadrature* rule is a quadrature rule constructed to yield an exact result for polynomials up to degree $2K - 1$ by a suitable choice of the *roots* x_k and *weights* w_k for $k = 1, \dots, K$.”

$$\int_a^b f(x) w(x) dx \approx \sum_k^K \underbrace{w_k}_{\text{weights}} f(\underbrace{x_k}_{\text{roots}}) \quad (48)$$

Quadrature rules

Interval $[a, b]$	Weight function $w(x)$	Orthogonal polynomials	Name
$[-1, 1]$	1	Legendre $P_n(x)$	Gauss-Legendre
$(-1, 1)$	$(1-x)^\alpha(1+x)^\beta$, $\alpha, \beta > -1$	Jacobi $P_n^{\alpha, \beta}(x)$	Gauss-Jacobi
$(-1, 1)$	$1/\sqrt{1-x^2}$	Chebyshev (1st kind) $T_n(x)$	Gauss-Chebyshev
$[-1, 1]$	$\sqrt{1-x^2}$	Chebyshev (2nd kind) $U_n(x)$	Gauss-Chebyshev
$[0, \infty)$	$\exp(-x)$	Laguerre $L_n(x)$	Gauss-Laguerre
$[0, \infty)$	$x^\alpha \exp(-x)$, $\alpha > -1$	Generalized Laguerre $L_n^\alpha(x)$	Gauss-Laguerre
$(-\infty, \infty)$	$\exp(-x^2)$	Hermite $H_n(x)$	Gauss-Hermite

https://en.wikipedia.org/wiki/Gaussian_quadrature

ACFDT at the RPA/RPax level

RPA matrix elements

$$A_{ia,jb}^{\lambda,RPA} = \delta_{ij}\delta_{ab}(\epsilon_a^{\text{HF}} - \epsilon_i^{\text{HF}}) + 2\lambda(ia|bj) \quad B_{ia,jb}^{\lambda,RPA} = 2\lambda(ia|jb) \quad (49)$$

$$E_c^{\text{RPA}} = \frac{1}{2} \int_0^1 \text{Tr}(\mathbf{K}\mathbf{P}^\lambda) d\lambda = \frac{1}{2} \left[\sum_m \Omega_m^{\text{RPA}} - \text{Tr}(\mathbf{A}^{\text{RPA}}) \right] \quad (50)$$

RPax matrix elements

$$A_{ia,jb}^{\lambda,RPax} = \delta_{ij}\delta_{ab}(\epsilon_a^{\text{HF}} - \epsilon_i^{\text{HF}}) + \lambda[2(ia|bj) - (ij|ab)] \quad B_{ia,jb}^{\lambda,RPax} = \lambda[2(ia|jb) - (ib|aj)] \quad (51)$$

$$E_c^{\text{RPax}} = \frac{1}{2} \int_0^1 \text{Tr}(\mathbf{K}\mathbf{P}^\lambda) d\lambda \neq \frac{1}{2} \left[\sum_m \Omega_m^{\text{RPax}} - \text{Tr}(\mathbf{A}^{\text{RPax}}) \right] \quad (52)$$

If exchange added to kernel, i.e., $\mathbf{K} = \mathbf{K}^x$, then [Angyan et al. JCTC 7 (2011) 3116]

$$E_c^{\text{RPax}} = \frac{1}{4} \int_0^1 \text{Tr}(\mathbf{K}^x\mathbf{P}^\lambda) d\lambda = \frac{1}{4} \left[\sum_m \Omega_m^{\text{RPax}} - \text{Tr}(\mathbf{A}^{\text{RPax}}) \right] \quad (53)$$

ACFDT at the BSE level

BSE matrix elements

$$A_{ia,jb}^{\lambda,\text{BSE}} = \delta_{ij}\delta_{ab}(\epsilon_a^{\text{GW}} - \epsilon_i^{\text{GW}}) + \lambda \left[2(ia|bj) - W_{ij,ab}^{\lambda}(\omega = 0) \right] \quad B_{ia,jb}^{\lambda,\text{BSE}} = \lambda \left[2(ia|jb) - W_{ib,ja}^{\lambda}(\omega = 0) \right] \quad (54)$$

$$E_c^{\text{BSE}} = \frac{1}{2} \int_0^1 \text{Tr}(\mathbf{K}\mathbf{P}^{\lambda}) d\lambda \neq \frac{1}{2} \left[\sum_m \Omega_m^{\text{BSE}} - \text{Tr}(\mathbf{A}^{\text{BSE}}) \right] \quad (55)$$

 λ -dependent screening

$$W_{pq,rs}^{\lambda}(\omega) = (pq|rs) + 2 \sum_m (pq|m)^{\lambda} (rs|m)^{\lambda} \left[\frac{1}{\omega - \Omega_m^{\lambda,\text{RPA}} + i\eta} - \frac{1}{\omega + \Omega_m^{\lambda,\text{RPA}} - i\eta} \right] \quad (56)$$

$$(pq|m)^{\lambda} = \sum_{ia} (pq|ia) (\mathbf{X}_m^{\lambda,\text{RPA}} + \mathbf{Y}_m^{\lambda,\text{RPA}})_{ia} \quad (57)$$

ACFDT in a computer

ACFDT correlation energy from BSE

procedure ACFDT FOR BSE

Compute GW quasiparticle energies ϵ^{GW} and interaction kernel \mathbf{K}

Get Gauss-Legendre weights and roots $\{w_k, \lambda_k\}_{1 \leq k \leq K}$

$E_c \leftarrow 0$

for $k = 1, \dots, K$ **do**

Compute static screening elements $W_{pq,rs}^{\lambda_k}(\omega = 0)$

Perform BSE calculation at $\lambda = \lambda_k$ to get \mathbf{X}^{λ_k} and \mathbf{Y}^{λ_k} \triangleright This is a $\mathcal{O}(N^6)$ step done many times!

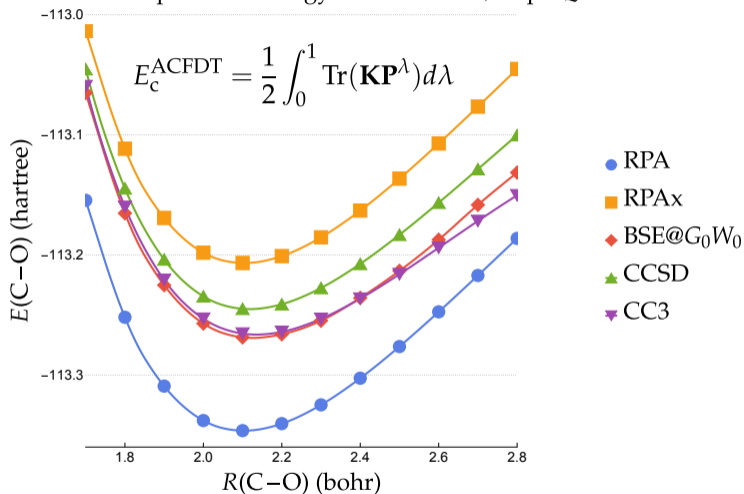
Form two-particle density matrix \mathbf{P}^{λ_k}

$E_c \leftarrow E_c + w_k \text{Tr}(\mathbf{K}\mathbf{P}^{\lambda_k})$

end for

end procedure

Ground-state potential energy surface of CO/cc-pVQZ



Loos et al. JPLCL 11 (2020) 3536

Useful papers/programs

- **molGW**: Bruneval et al. *Comp. Phys. Comm.* 208 (2016) 149
- **Turbomole**: van Setten et al. *JCTC* 9 (2013) 232; Kaplan et al. *JCTC* 12 (2016) 2528
- **Fiesta**: Blase et al. *Chem. Soc. Rev.* 47 (2018) 1022
- **FHI-AIMS**: Caruso et al. *PRB* 86 (2012) 081102
- **Reviews & Books:**
 - Reining, *WIREs Comput Mol Sci* 2017, e1344. doi: 10.1002/wcms.1344
 - Onida et al. *Rev. Mod. Phys.* 74 (2002) 601
 - Blase et al. *Chem. Soc. Rev.* , 47 (2018) 1022
 - Golze et al. *Front. Chem.* 7 (2019) 377
 - Blase et al. *JPCL* 11 (2020) 7371
 - Martin, Reining & Ceperley *Interacting Electrons* (Cambridge University Press)
- **GW100**: IPs for a set of 100 molecules. van Setten et al. *JCTC* 11 (2015) 5665 (<http://gw100.wordpress.com>)