



Introduction to density matrix embedding theory

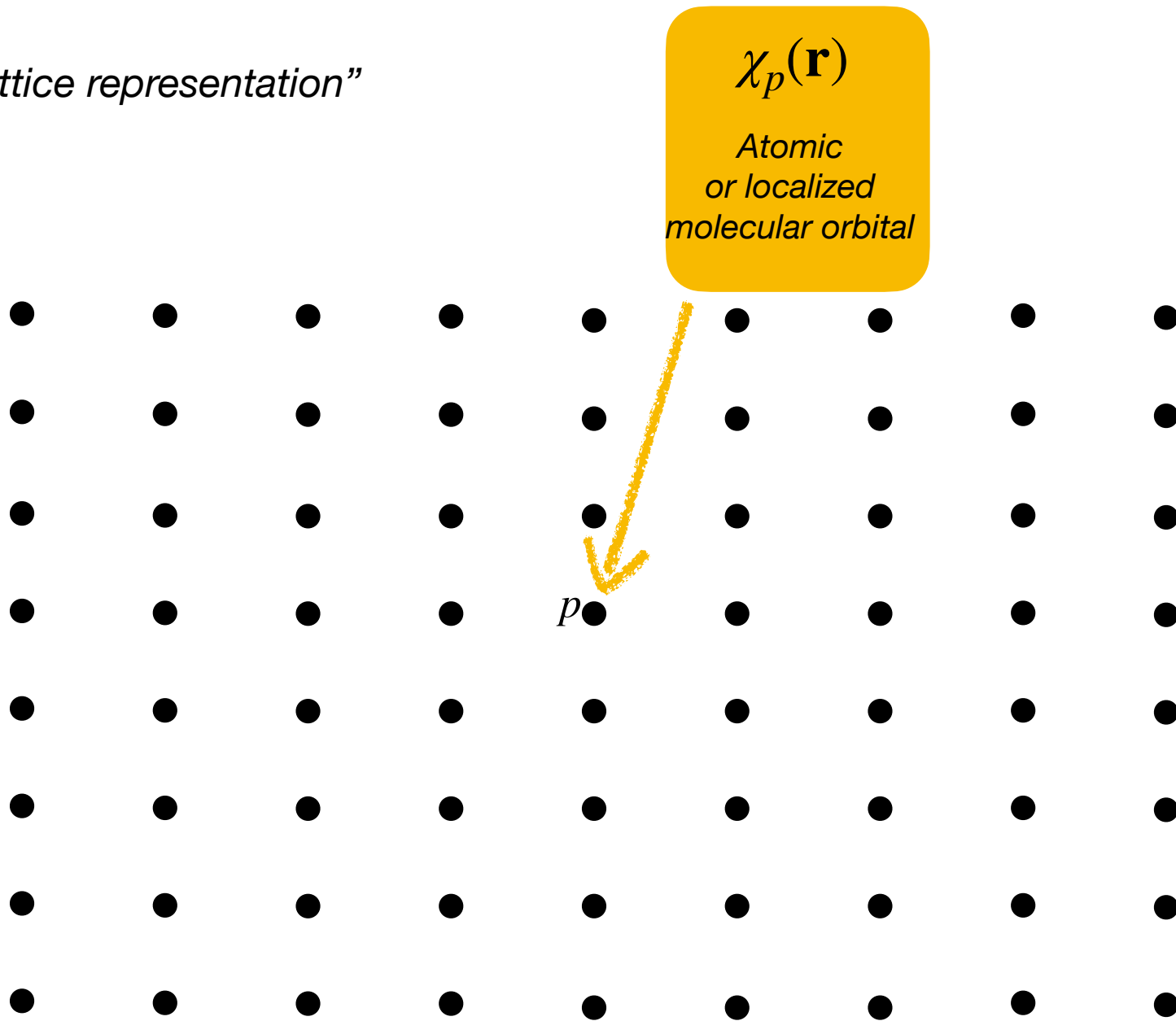
Emmanuel Fromager

*Laboratoire de Chimie Quantique, Institut de Chimie de Strasbourg,
Université de Strasbourg, Strasbourg, France*

On the local treatment of (strong) electron correlation

Local evaluation of the energy (in a *localised* spin-orbital basis)

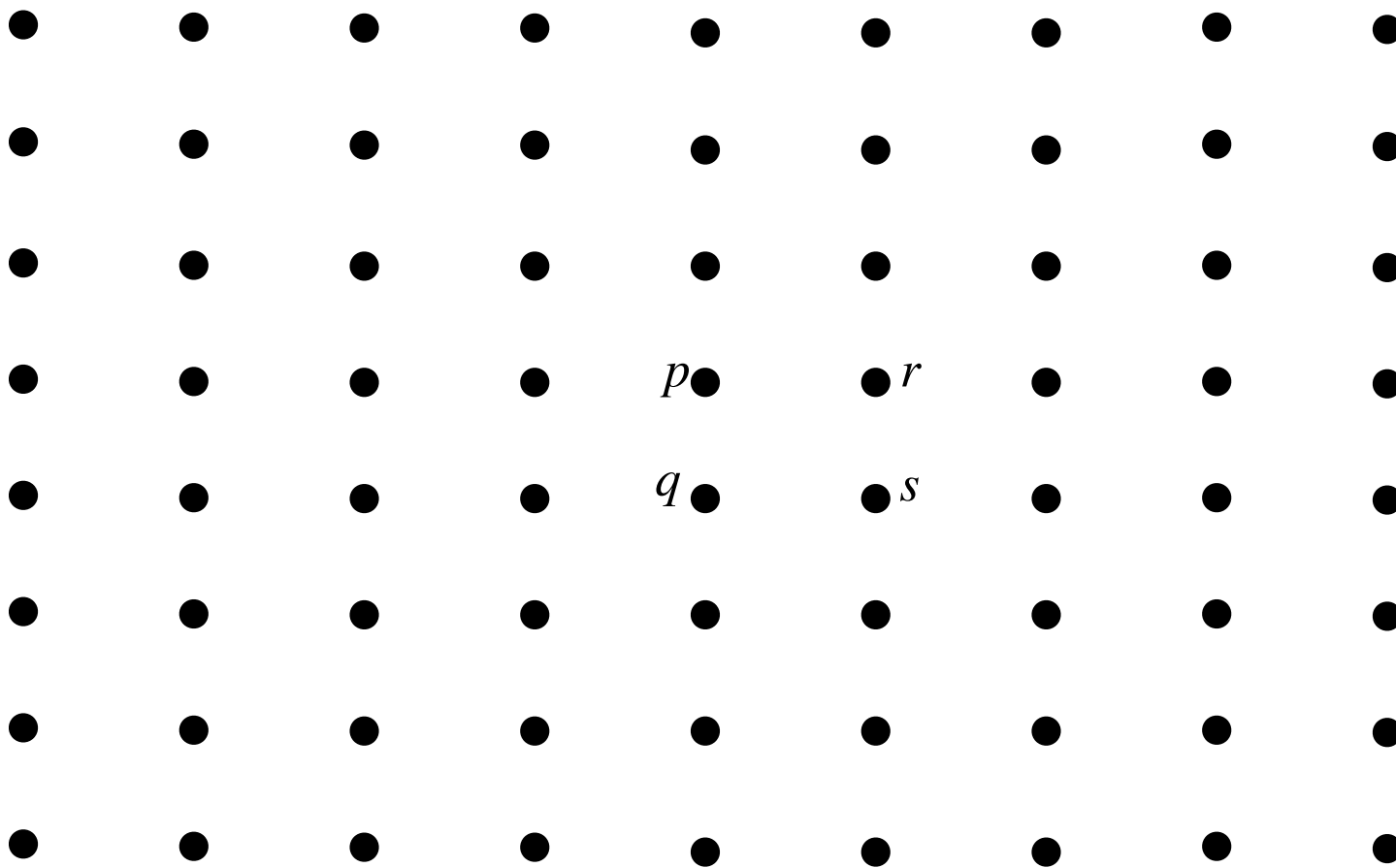
So-called “lattice representation”



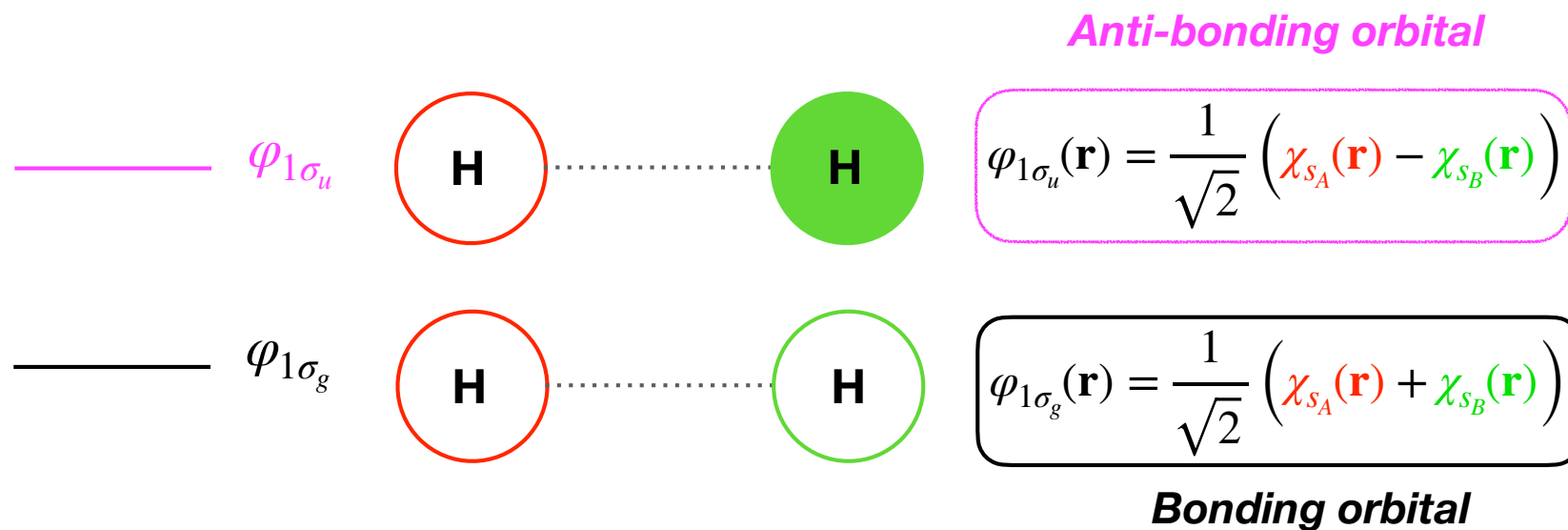
Local evaluation of the energy (in a **localised** spin-orbital basis)

So-called “lattice representation”

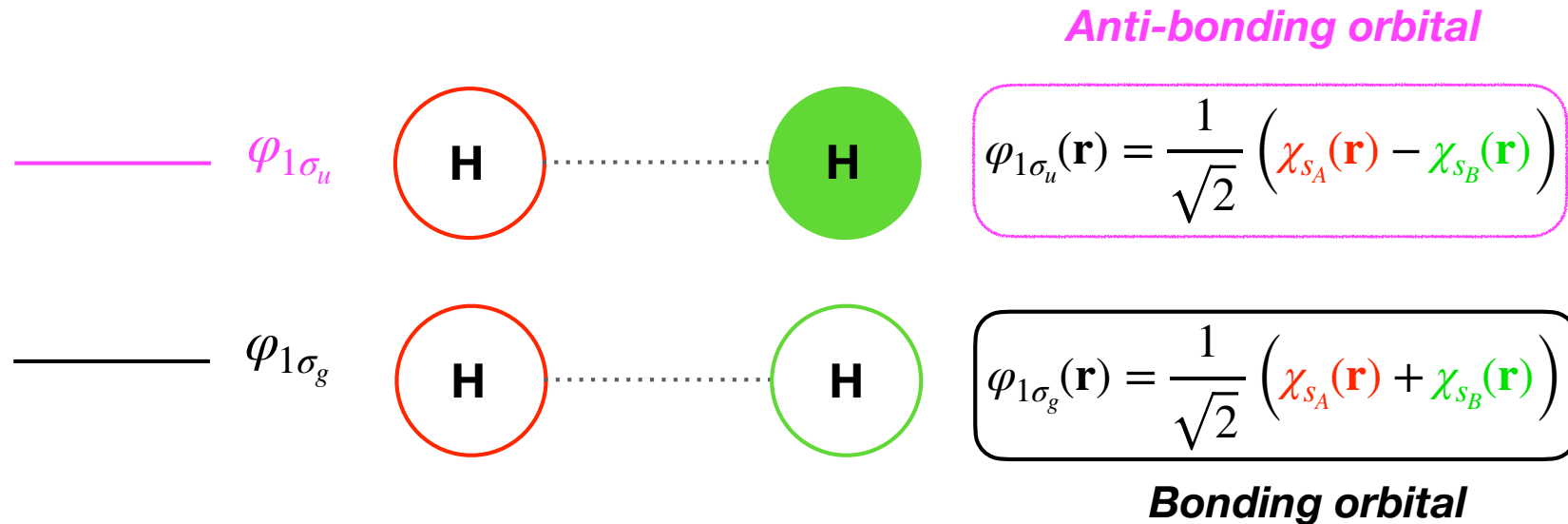
$$\langle \hat{H} \rangle = \sum_{pq} h_{pq} \underbrace{\langle \hat{c}_p^\dagger \hat{c}_q \rangle}_{\text{One-electron density matrix (1RDM)}} + \frac{1}{2} \sum_{pqrs} \langle pq | rs \rangle \underbrace{\langle \hat{c}_p^\dagger \hat{c}_q^\dagger \hat{c}_s \hat{c}_r \rangle}_{\text{Two-electron density matrix (2RDM)}}$$



A brief reminder: Multi-configurational description of the stretched hydrogen molecule



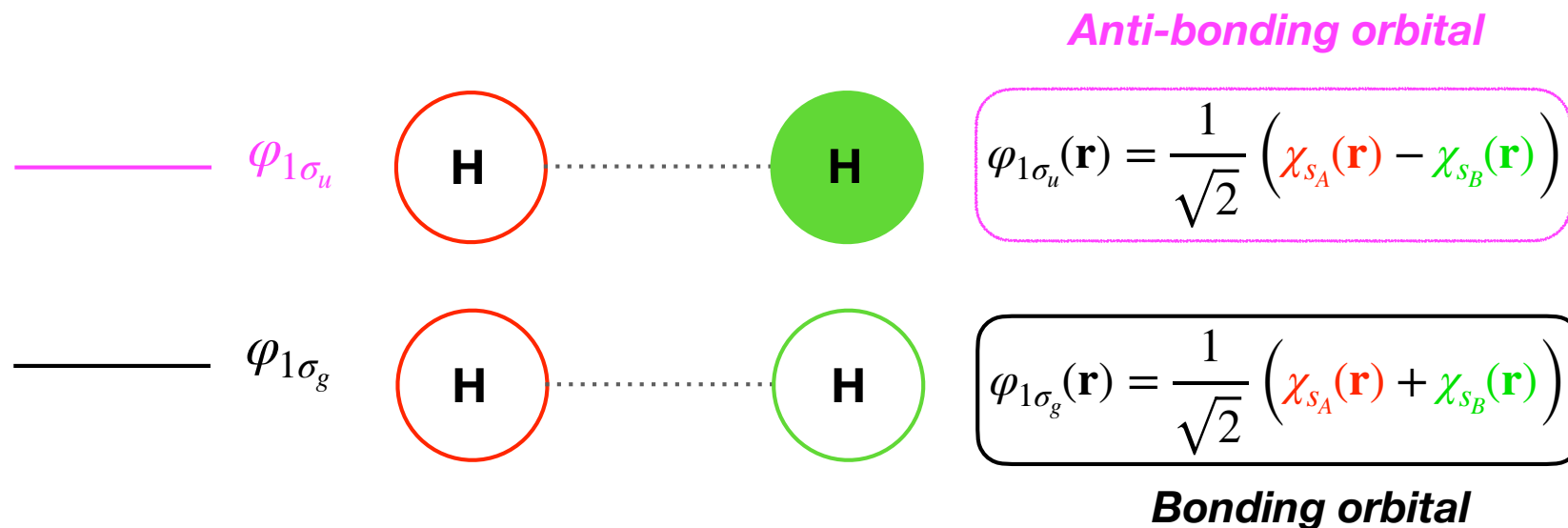
Multi-configurational description of the **stretched hydrogen molecule**



$$\Psi \equiv \frac{1}{\sqrt{2}} \left(\varphi_{1\sigma_g}(\mathbf{r}_1)\varphi_{1\sigma_g}(\mathbf{r}_2) - \varphi_{1\sigma_u}(\mathbf{r}_1)\varphi_{1\sigma_u}(\mathbf{r}_2) \right)$$

← Delocalised picture
(Chemistry)

Multi-configurational description of the stretched hydrogen molecule



$$\Psi \equiv \frac{1}{\sqrt{2}} \left(\varphi_{1\sigma_g}(\mathbf{r}_1)\varphi_{1\sigma_g}(\mathbf{r}_2) - \varphi_{1\sigma_u}(\mathbf{r}_1)\varphi_{1\sigma_u}(\mathbf{r}_2) \right)$$

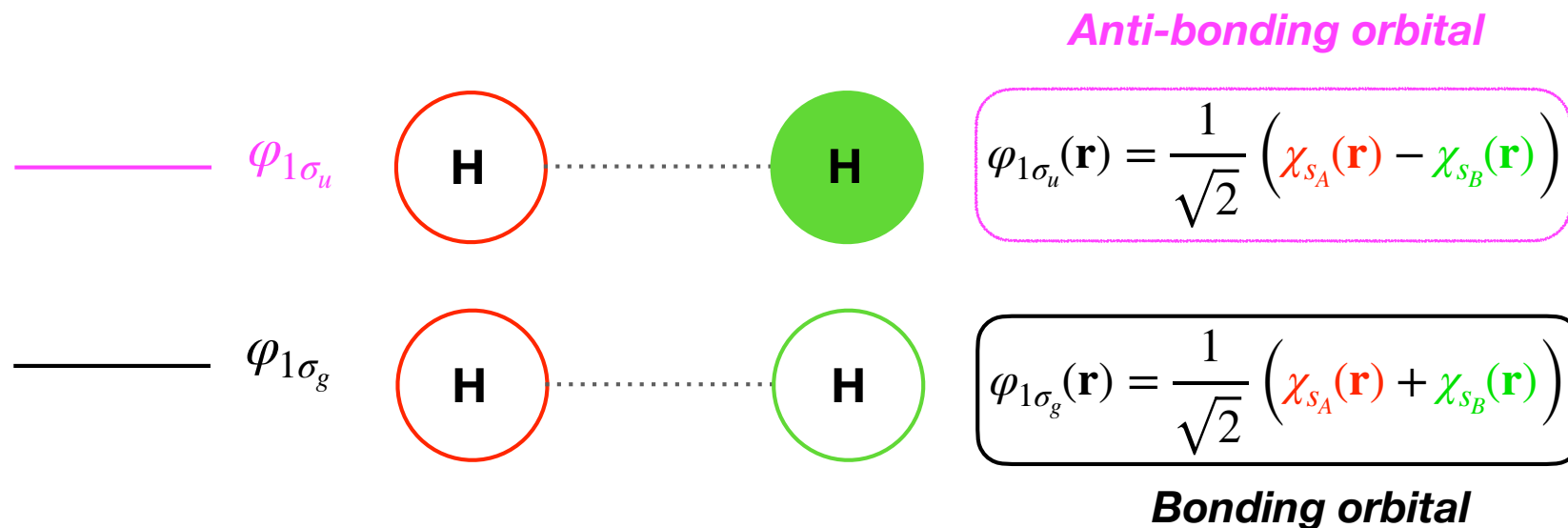
← **Delocalised picture
(Chemistry)**

$$= \frac{1}{\sqrt{2}} \left(\chi_{s_A}(\mathbf{r}_1)\chi_{s_B}(\mathbf{r}_2) + \chi_{s_A}(\mathbf{r}_2)\chi_{s_B}(\mathbf{r}_1) \right)$$

← **Localised picture
(Physics)**

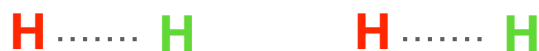


Multi-configurational description of the stretched hydrogen molecule



$$\Psi \equiv \frac{1}{\sqrt{2}} \left(\varphi_{1\sigma_g}(\mathbf{r}_1)\varphi_{1\sigma_g}(\mathbf{r}_2) - \varphi_{1\sigma_u}(\mathbf{r}_1)\varphi_{1\sigma_u}(\mathbf{r}_2) \right) \quad \leftarrow \text{Delocalised picture (Chemistry)}$$

$$= \frac{1}{\sqrt{2}} \left(\chi_{s_A}(\mathbf{r}_1)\chi_{s_B}(\mathbf{r}_2) + \chi_{s_A}(\mathbf{r}_2)\chi_{s_B}(\mathbf{r}_1) \right) \quad \leftarrow \text{Localised picture (Physics)}$$



Consequence of the electronic repulsion on each atom!

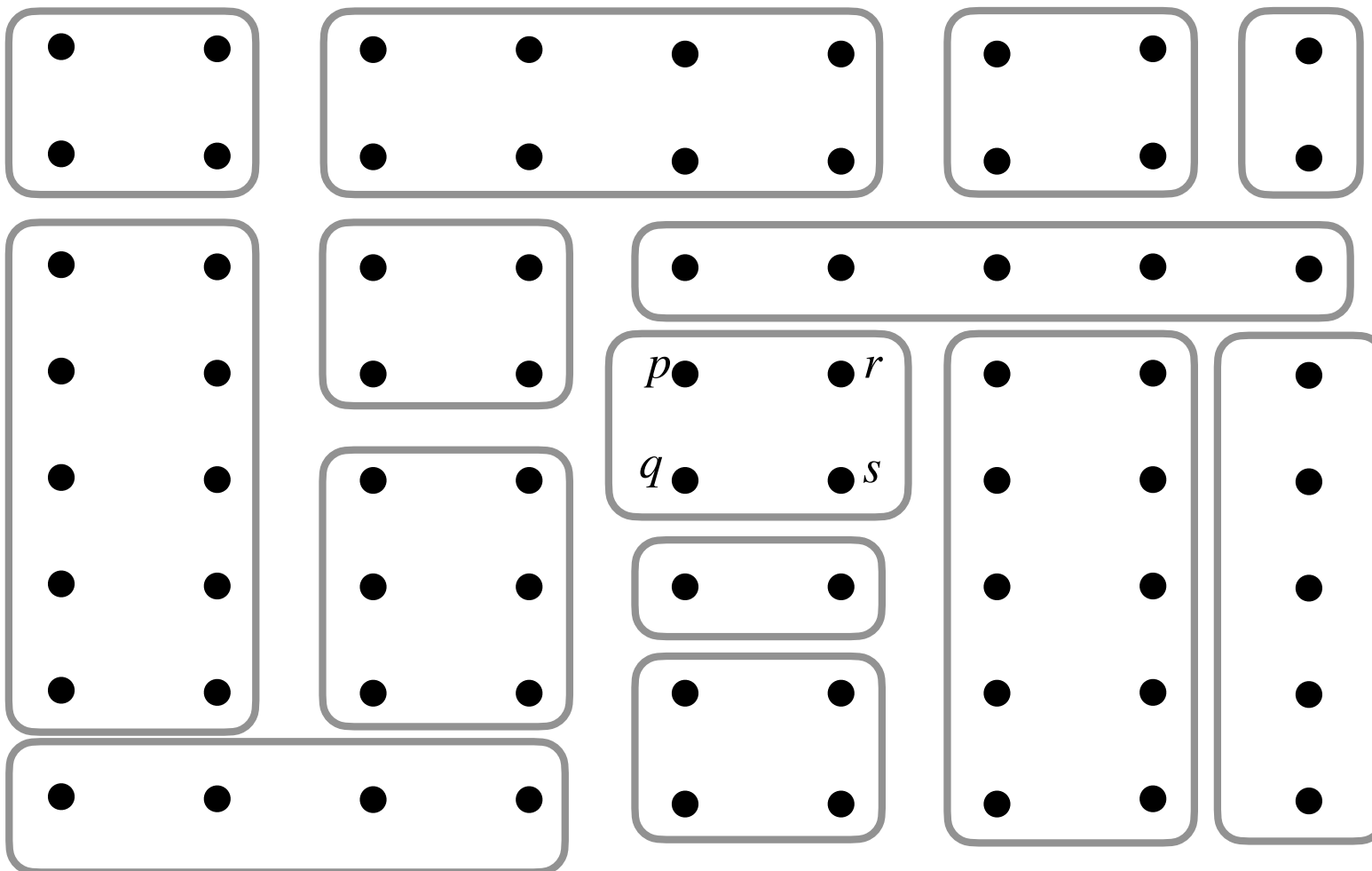
Local evaluation of the energy (in a **localised** spin-orbital basis)

Fragmentation

for treating **strong local electron correlations**

$$\langle \hat{H} \rangle = \sum_{pq} h_{pq} \langle \hat{c}_p^\dagger \hat{c}_q \rangle + \frac{1}{2} \sum_{pqrs} \langle pq | rs \rangle \langle \hat{c}_p^\dagger \hat{c}_q^\dagger \hat{c}_s \hat{c}_r \rangle$$

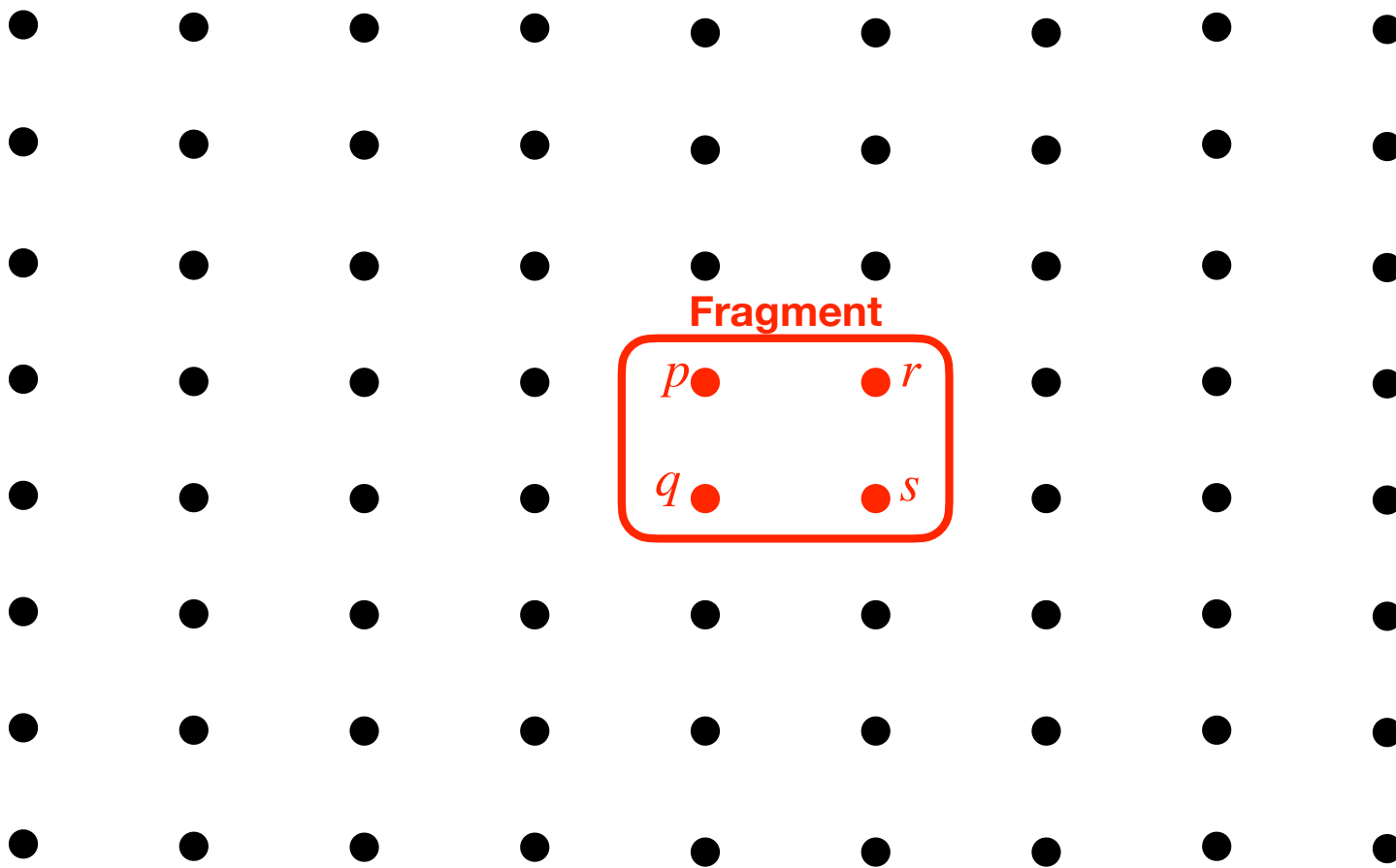
One-electron density matrix (1RDM)
Two-electron density matrix (2RDM)



Local evaluation of the energy (in a **localised** spin-orbital basis)

So-called “lattice representation”

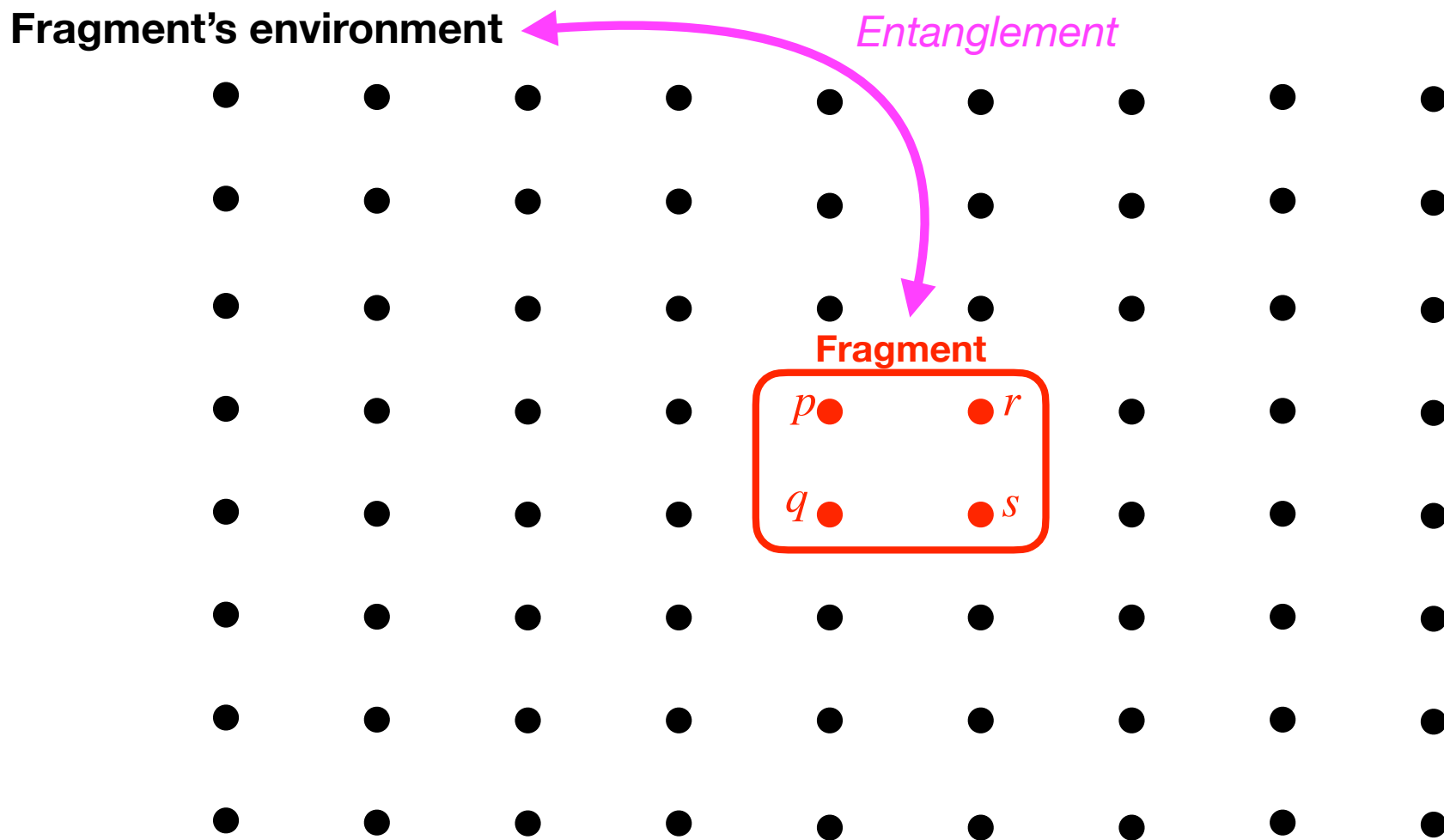
$$\langle \hat{H} \rangle = \sum_{pq} h_{pq} \underbrace{\langle \hat{c}_p^\dagger \hat{c}_q \rangle}_{\text{One-electron density matrix (1RDM)}} + \frac{1}{2} \sum_{pqrs} \langle pq | rs \rangle \underbrace{\langle \hat{c}_p^\dagger \hat{c}_q^\dagger \hat{c}_s \hat{c}_r \rangle}_{\text{Two-electron density matrix (2RDM)}}$$



Local evaluation of the energy (in a *localised spin-orbital basis*)

So-called “lattice representation”

$$\langle \hat{H} \rangle = \sum_{pq} h_{pq} \langle \hat{c}_p^\dagger \hat{c}_q \rangle + \frac{1}{2} \sum_{pqrs} \langle pq | rs \rangle \langle \hat{c}_p^\dagger \hat{c}_q^\dagger \hat{c}_s \hat{c}_r \rangle$$

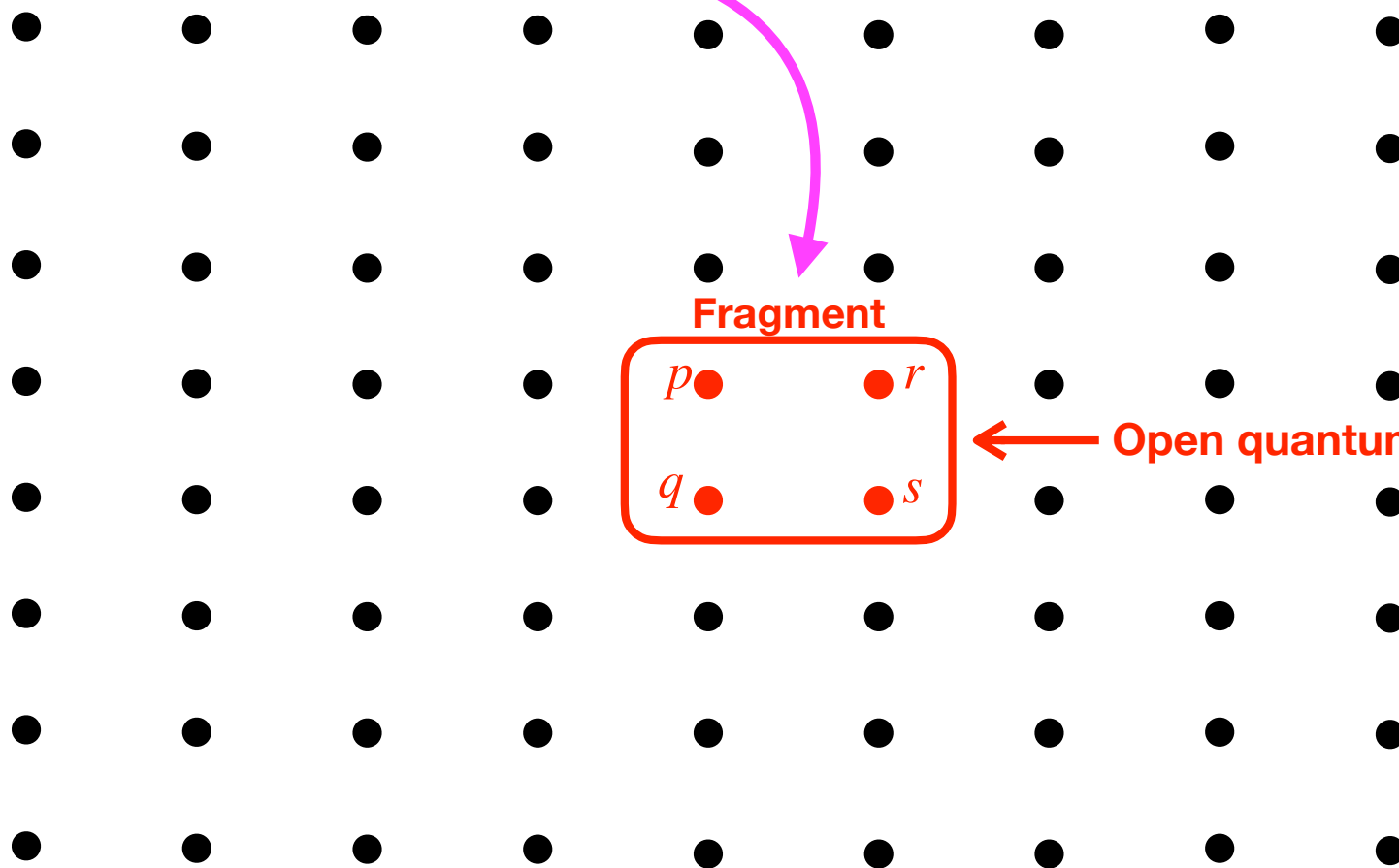


Local evaluation of the energy (in a **localised spin-orbital basis**)

So-called “lattice representation”

$$\langle \hat{H} \rangle = \sum_{pq} h_{pq} \langle \hat{c}_p^\dagger \hat{c}_q \rangle + \frac{1}{2} \sum_{pqrs} \langle pq | rs \rangle \langle \hat{c}_p^\dagger \hat{c}_q^\dagger \hat{c}_s \hat{c}_r \rangle$$

Fragment's environment



Fragment

Open quantum system!

Quantum entanglement of a fragment with its environment

$$\hat{H} \equiv \sum_{PQ} h_{PQ} \hat{c}_P^\dagger \hat{c}_Q + \frac{1}{2} \sum_{PQRS} g_{PQRS} \hat{c}_P^\dagger \hat{c}_Q^\dagger \hat{c}_S \hat{c}_R$$

In principle, we need to **solve the Schrödinger equation** in order to evaluate the (ground-state) energy:

$$\hat{H} |\Psi_0\rangle = E_0 |\Psi_0\rangle$$

Quantum entanglement of a fragment with its environment

$$\hat{H} \equiv \sum_{PQ} h_{PQ} \hat{c}_P^\dagger \hat{c}_Q + \frac{1}{2} \sum_{PQRS} g_{PQRS} \hat{c}_P^\dagger \hat{c}_Q^\dagger \hat{c}_S \hat{c}_R$$

In principle, we need to **solve the Schrödinger equation** in order to evaluate the (ground-state) energy:

$$\hat{H} |\Psi_0\rangle = E_0 |\Psi_0\rangle$$

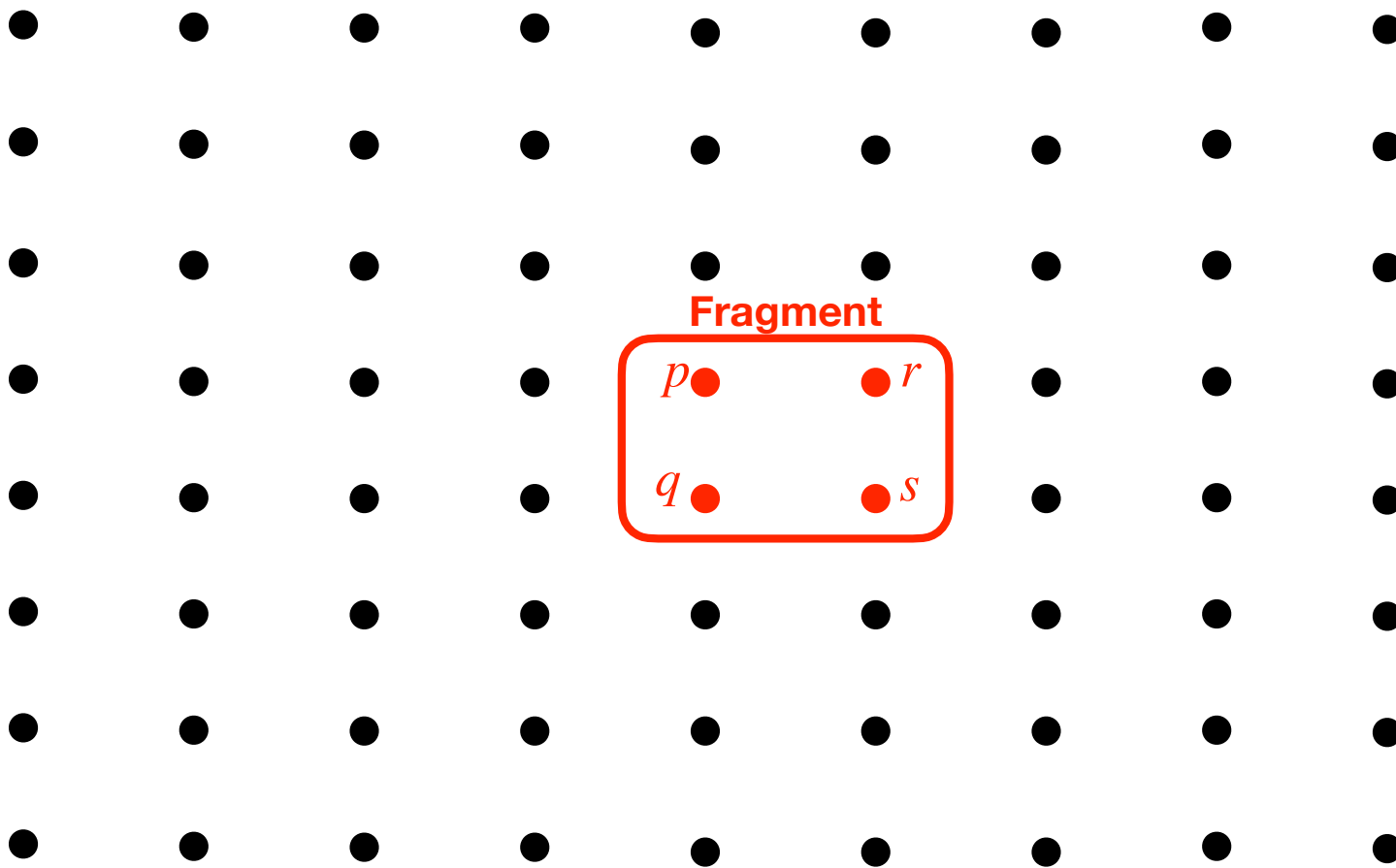
A $|\Psi_0\rangle$ consisting of electrons simply distributed among **disconnected fragments** **cannot** be described by $\hat{H} |\Psi_0\rangle$!

Philosophy of density matrix embedding theory (DMET)

Local evaluation of the energy (in a **localised** spin-orbital basis)

So-called “lattice representation”

$$\langle \hat{H} \rangle = \sum_{pq} h_{pq} \underbrace{\langle \hat{c}_p^\dagger \hat{c}_q \rangle}_{\substack{\text{One-electron} \\ \text{density matrix} \\ \text{(1RDM)}}} + \frac{1}{2} \sum_{pqrs} \langle pq | rs \rangle \underbrace{\langle \hat{c}_p^\dagger \hat{c}_q^\dagger \hat{c}_s \hat{c}_r \rangle}_{\substack{\text{Two-electron} \\ \text{density matrix} \\ \text{(2RDM)}}$$



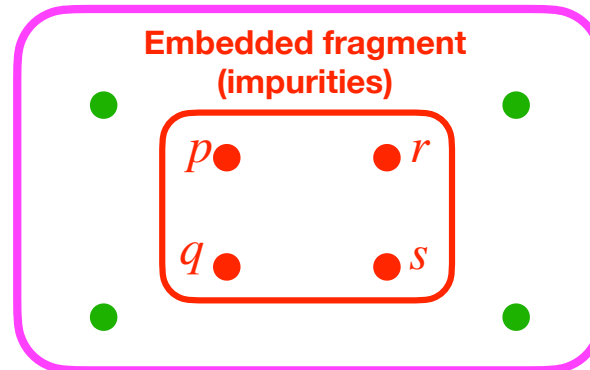
What are we aiming at?

Reduction in size of the problem to be solved:

$$\langle \hat{c}_p^\dagger \hat{c}_q \rangle_{full\ system} \approx \langle \hat{c}_p^\dagger \hat{c}_q \rangle_{\Psi^{\mathcal{C}}}$$
$$\langle \hat{c}_p^\dagger \hat{c}_q^\dagger \hat{c}_s \hat{c}_r \rangle_{full\ system} \approx \langle \hat{c}_p^\dagger \hat{c}_q^\dagger \hat{c}_s \hat{c}_r \rangle_{\Psi^{\mathcal{C}}}$$

Embedding cluster \mathcal{C}

Quantum bath
≡ electronic reservoir



$\Psi^{\mathcal{C}}$

Few-electron
correlated wave function

What are we aiming at?

Reduction in size of the problem to be solved:

$$\langle \hat{c}_p^\dagger \hat{c}_q \rangle_{full\ system} \approx \langle \hat{c}_p^\dagger \hat{c}_q \rangle_{\Psi^{\mathcal{C}}}$$

$$\langle \hat{c}_p^\dagger \hat{c}_q^\dagger \hat{c}_s \hat{c}_r \rangle_{full\ system} \approx \langle \hat{c}_p^\dagger \hat{c}_q^\dagger \hat{c}_s \hat{c}_r \rangle_{\Psi^{\mathcal{C}}}$$

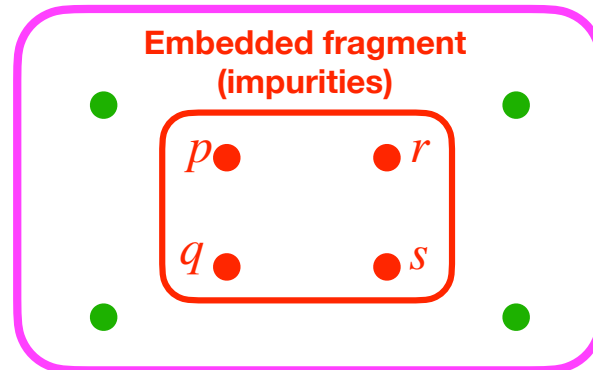
4 orbitals here:
Why and how?



Quantum bath

≡ electronic reservoir

Embedding cluster \mathcal{C}



$\Psi^{\mathcal{C}}$

Few-electron
correlated wave function

G. Knizia and G. K.-L. Chan, *Phys. Rev. Lett.* **109**, 186404 (2012).

S. Wouters, C. A. Jiménez-Hoyos, Q. Sun, and G. K.-L. Chan, *J. Chem. Theory Comput.* **12**, 2706 (2016).

S. Sekaran, M. Tsuchiizu, M. Saubanère, and E. Fromager, *Phys. Rev. B* **104**, 035121 (2021).

S. Sekaran, O. Bindech, and E. Fromager, *J. Chem. Phys.* **159**, 034107 (2023).

What are we aiming at?

Reduction in size of the problem to be solved:

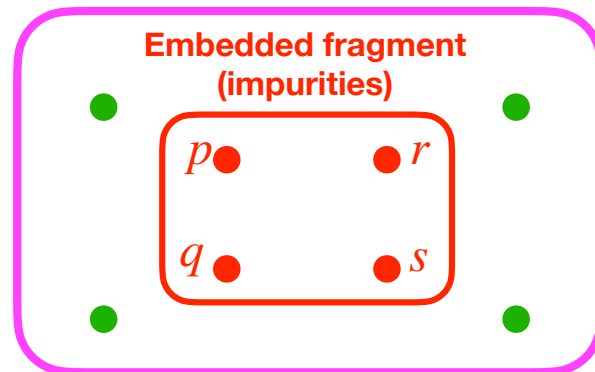
$$\langle \hat{c}_p^\dagger \hat{c}_q \rangle_{full\ system} \approx \langle \hat{c}_p^\dagger \hat{c}_q \rangle_{\Psi^{\mathcal{C}}}$$
$$\langle \hat{c}_p^\dagger \hat{c}_q^\dagger \hat{c}_s \hat{c}_r \rangle_{full\ system} \approx \langle \hat{c}_p^\dagger \hat{c}_q^\dagger \hat{c}_s \hat{c}_r \rangle_{\Psi^{\mathcal{C}}}$$

**4 orbitals here:
Why and how?**



Quantum bath
≡ electronic reservoir

Embedding cluster \mathcal{C}



$\Psi^{\mathcal{C}}$

How many?

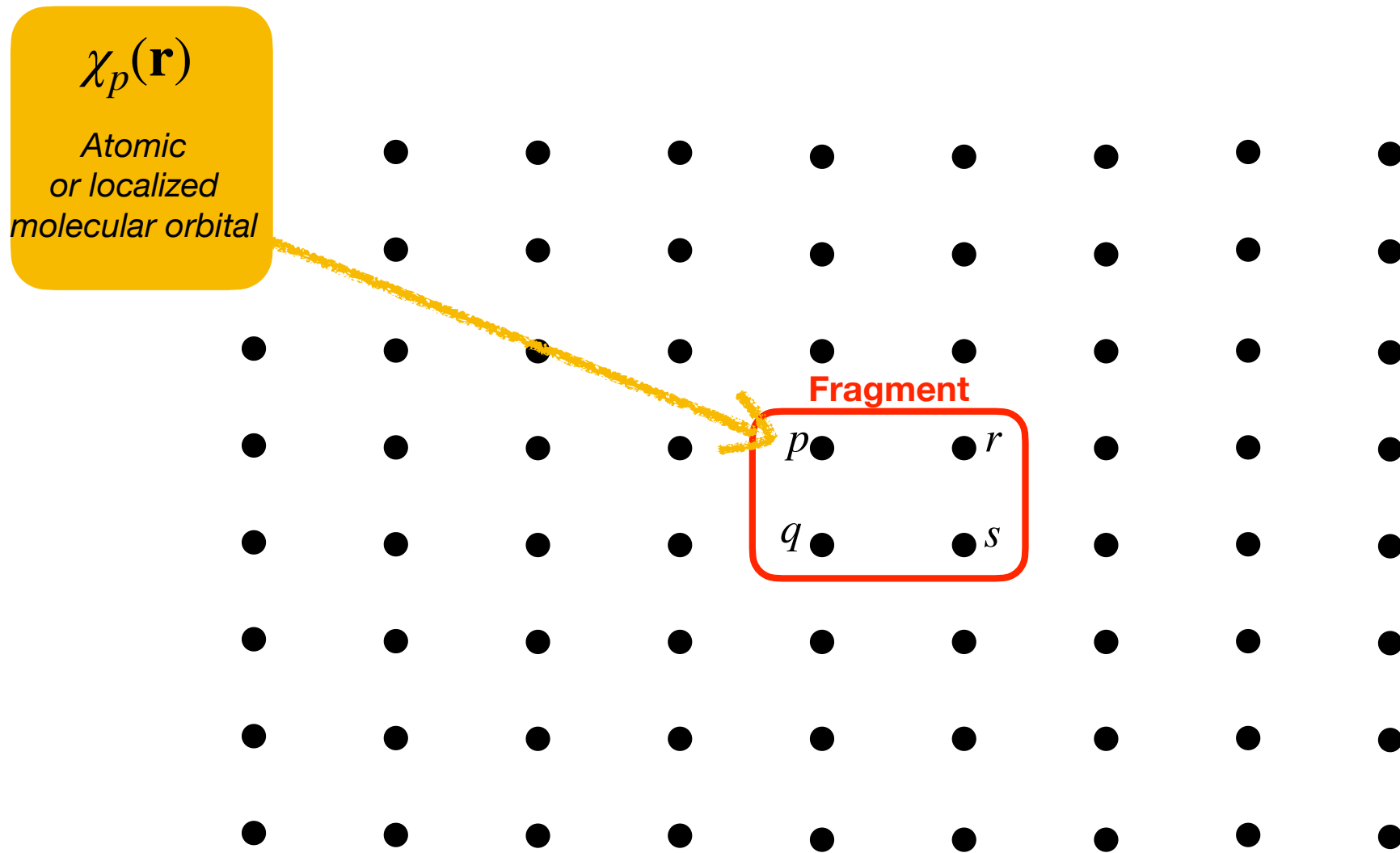


Few-electron
correlated wave function

Mathematical construction of the quantum bath in DMET

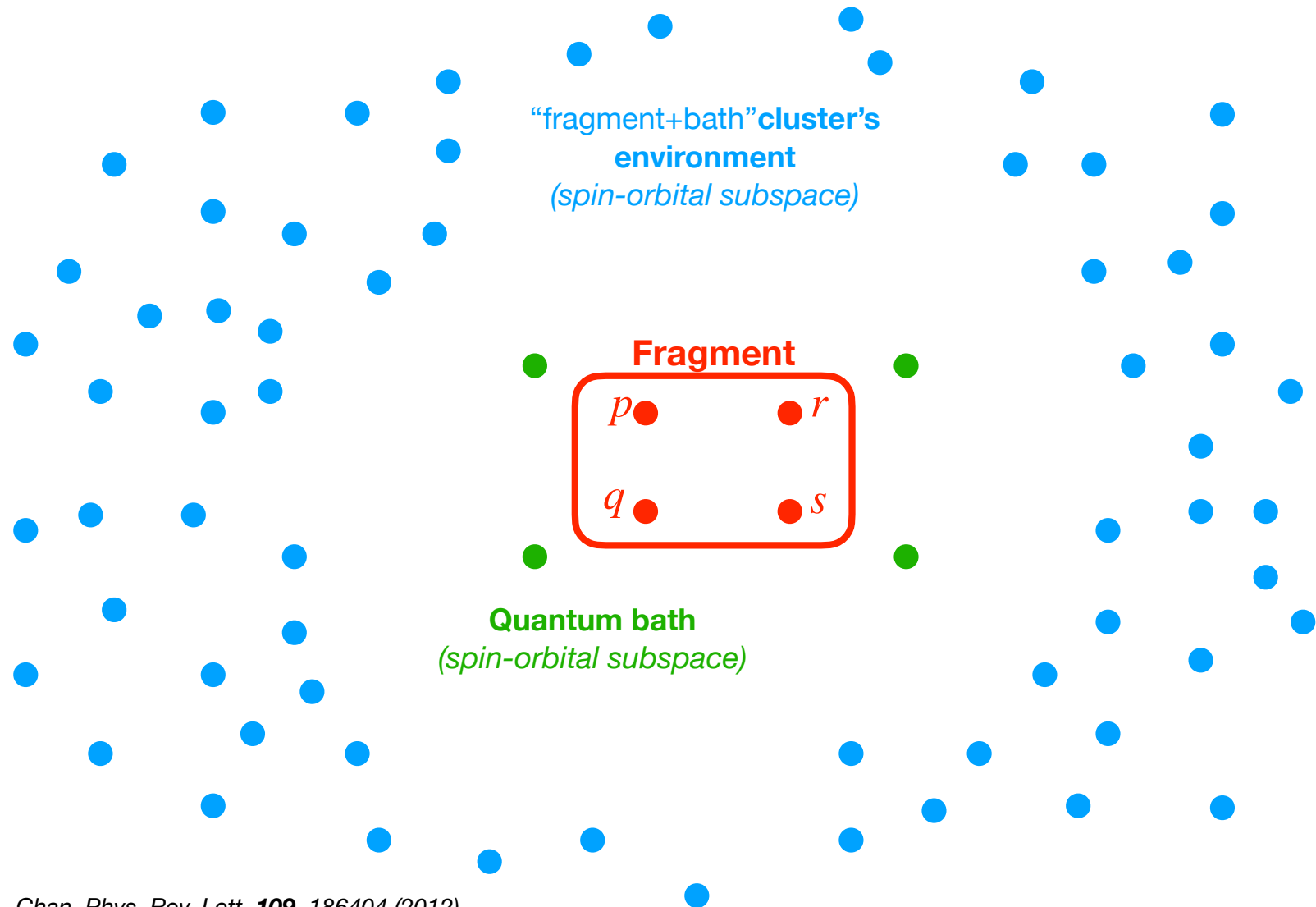
Clusterization through a unitary one-electron transformation

So-called “lattice representation”



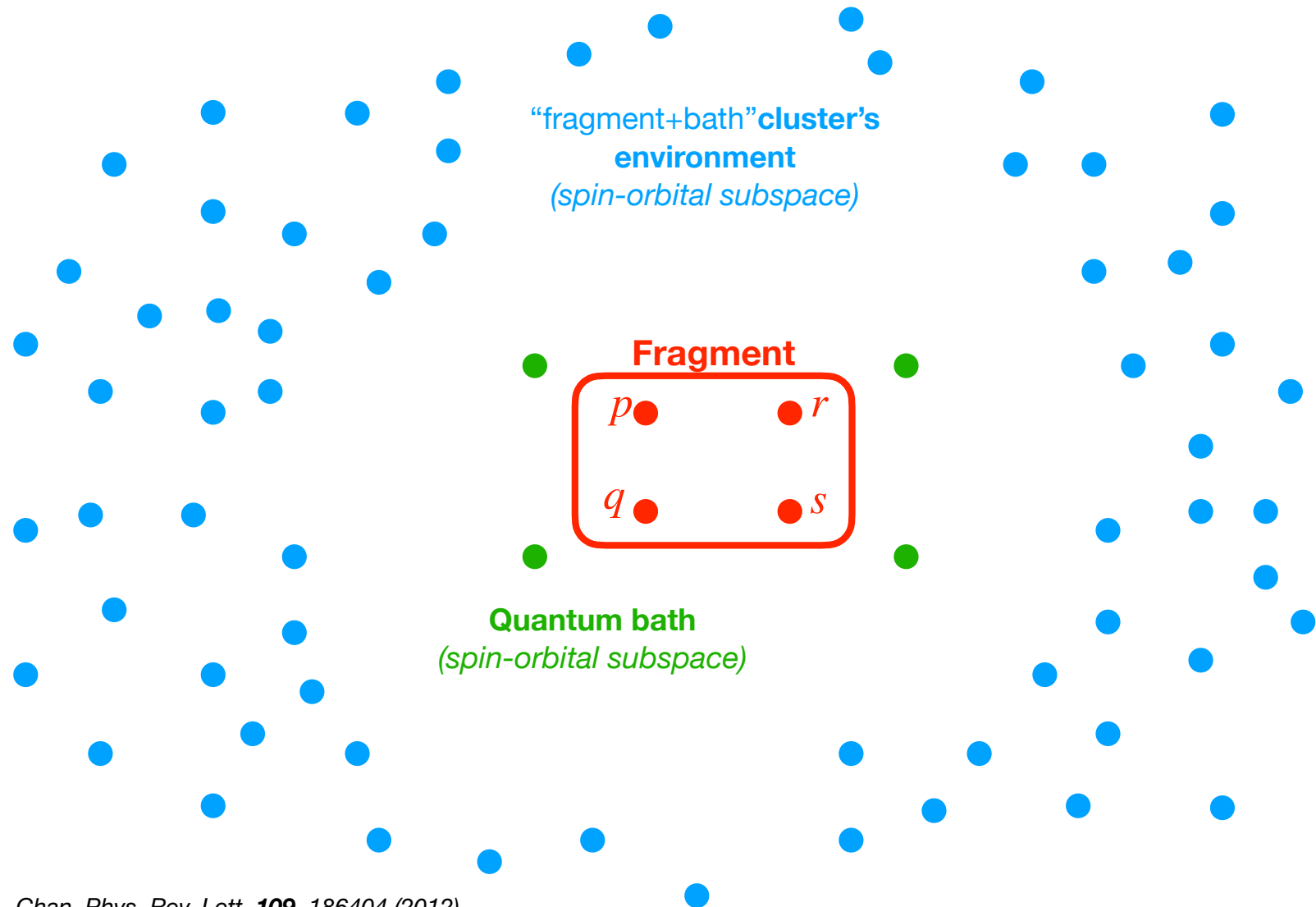
Clusterization through a unitary one-electron transformation

$$|\chi_q\rangle \rightarrow \sum_p^{\text{lattice}} \mathcal{U}_{pq}^* |\chi_p\rangle \quad \Leftrightarrow \quad \langle\chi_q| \rightarrow \sum_p^{\text{lattice}} \mathcal{U}_{pq} \langle\chi_p|$$



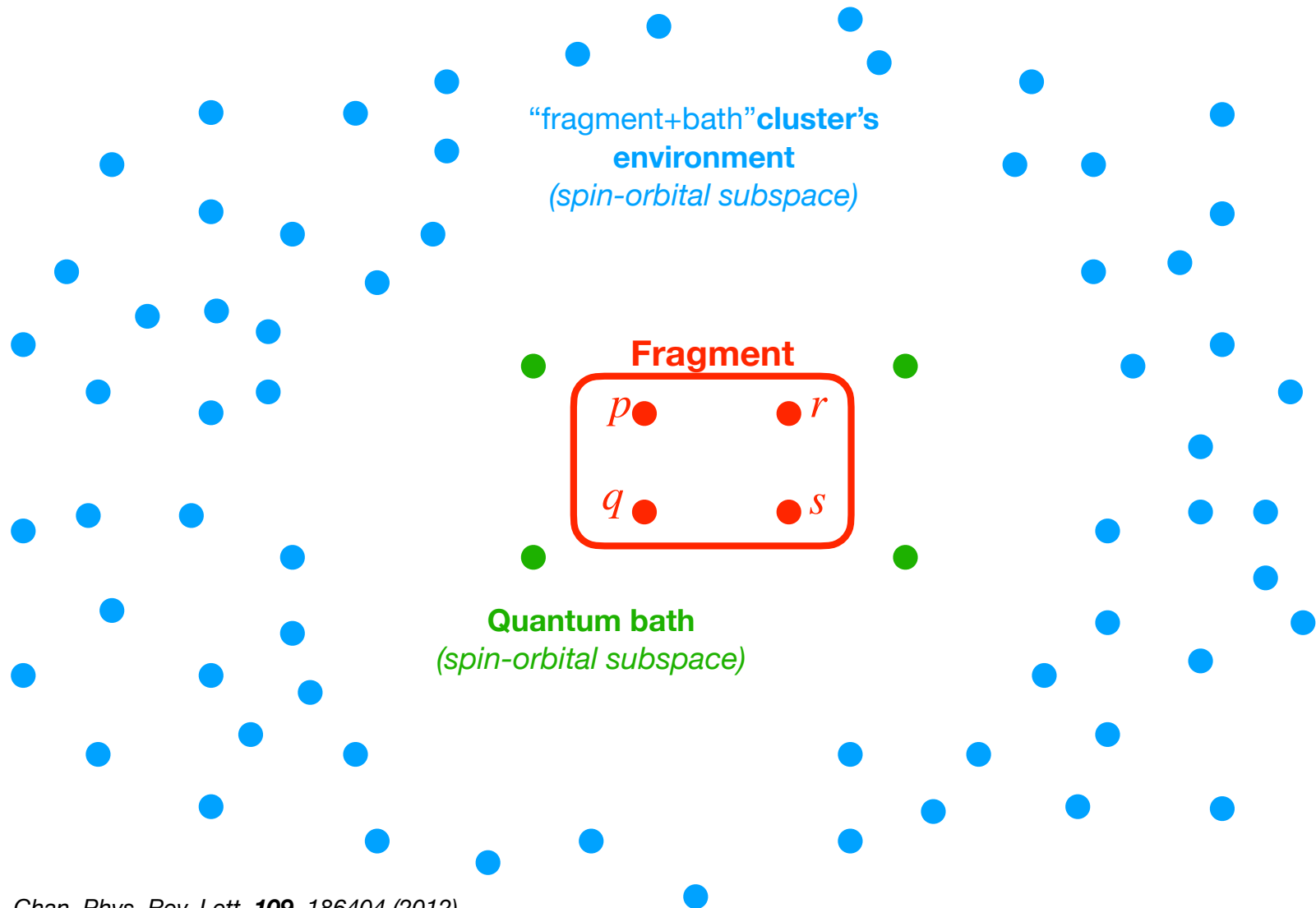
Clusterization through a unitary one-electron transformation

$$\langle \chi_q | \rightarrow \sum_p^{\text{lattice}} \mathcal{U}_{pq} \langle \chi_p | \quad \Leftrightarrow \quad \hat{c}_q \rightarrow \sum_p^{\text{lattice}} \mathcal{U}_{pq} \hat{c}_p \stackrel{\text{notation}}{=} \hat{d}_q$$

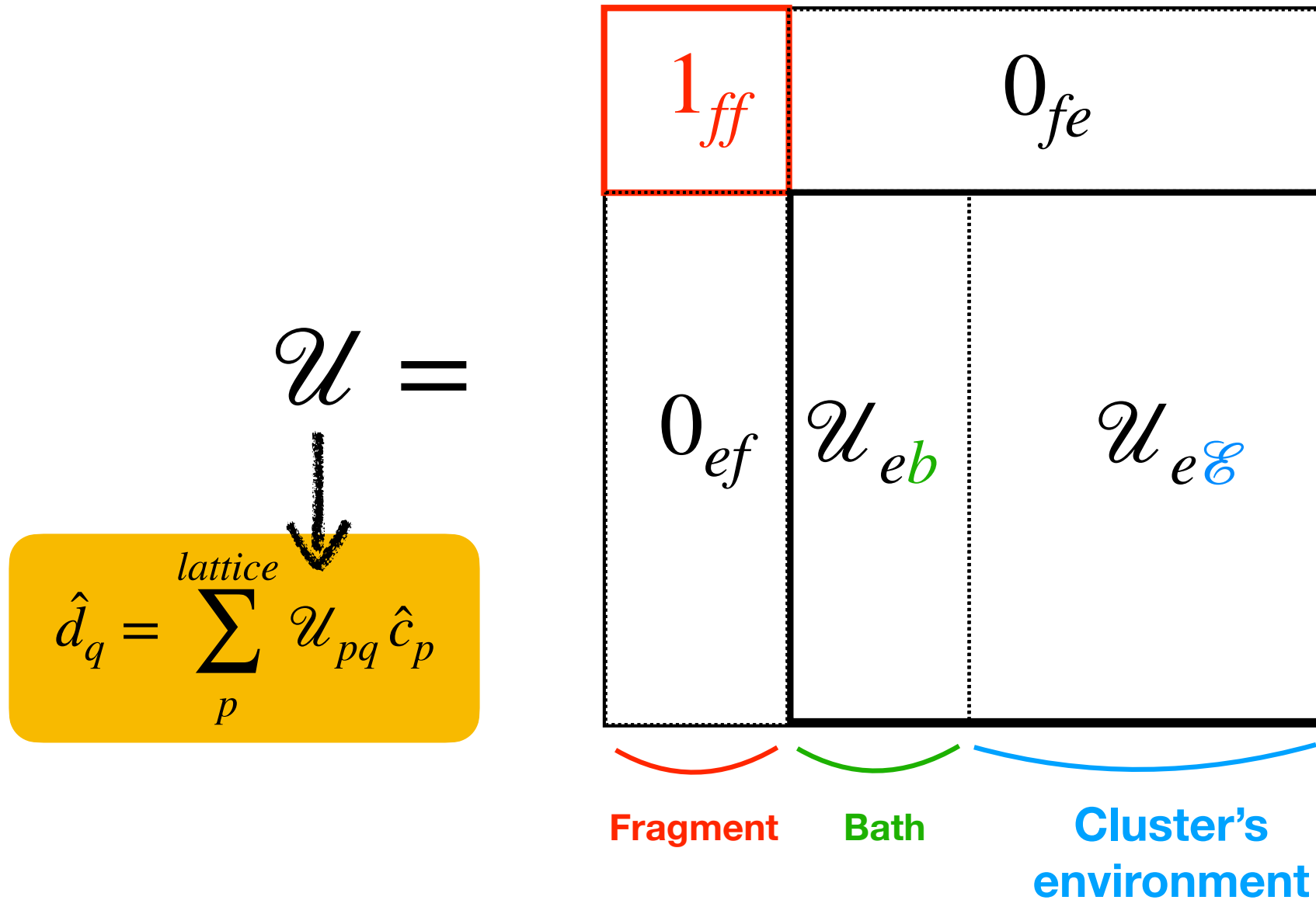


$$\hat{c}_q \rightarrow \sum_p^{\text{lattice}} \mathcal{U}_{pq} \hat{c}_p \stackrel{\text{notation}}{=} \hat{d}_q$$

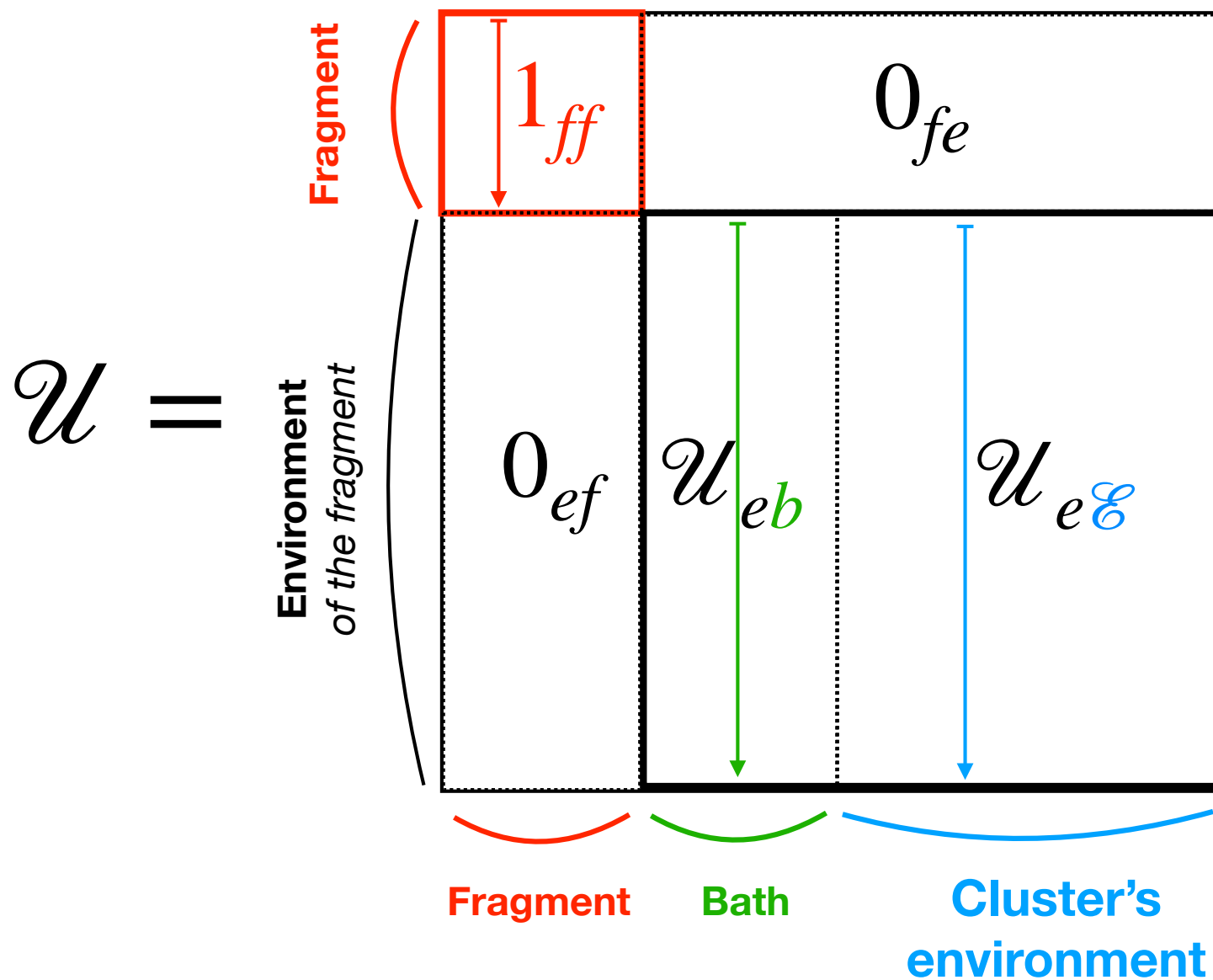
????



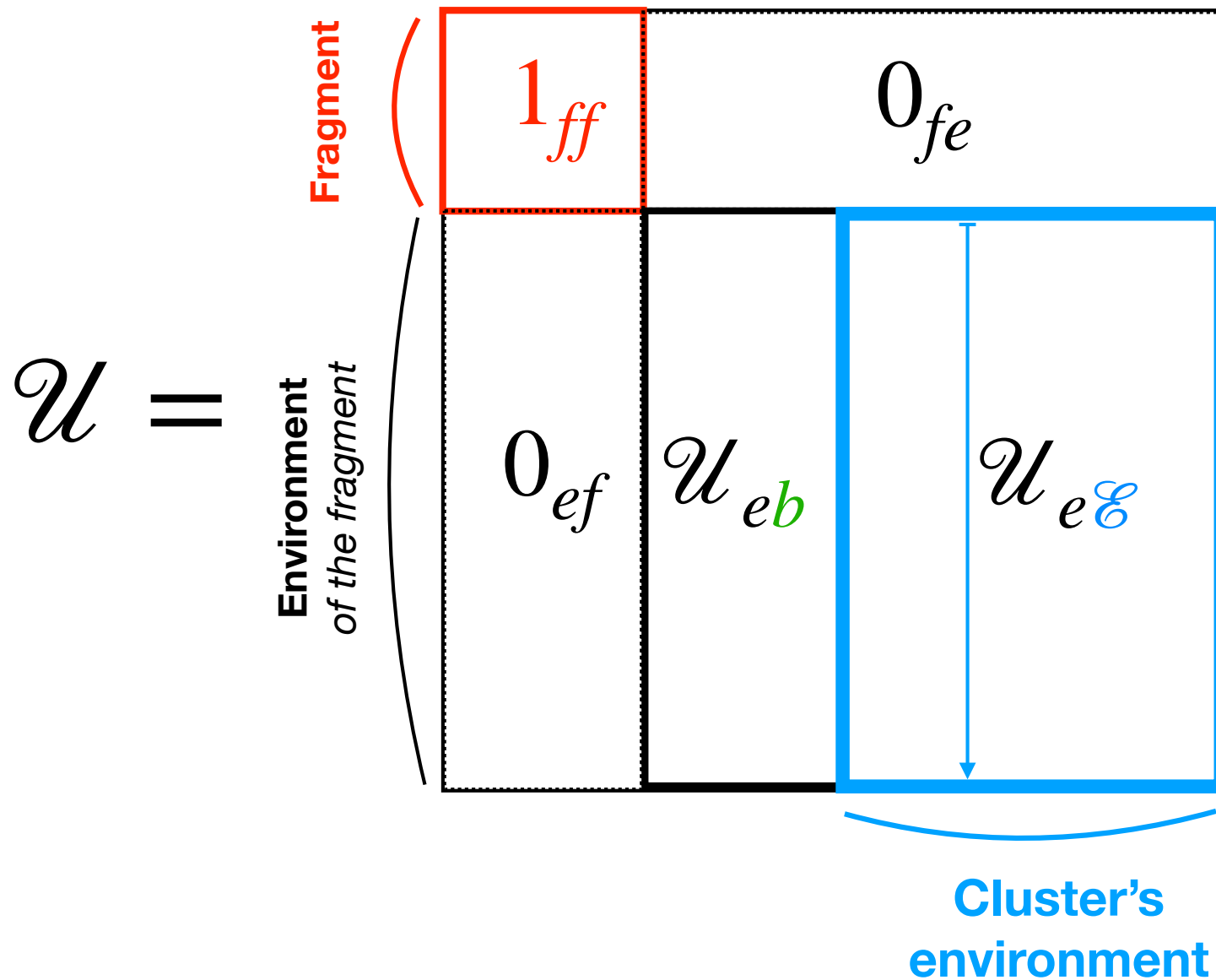
Quantum bath seen as a functional of the density matrix (1RDM)



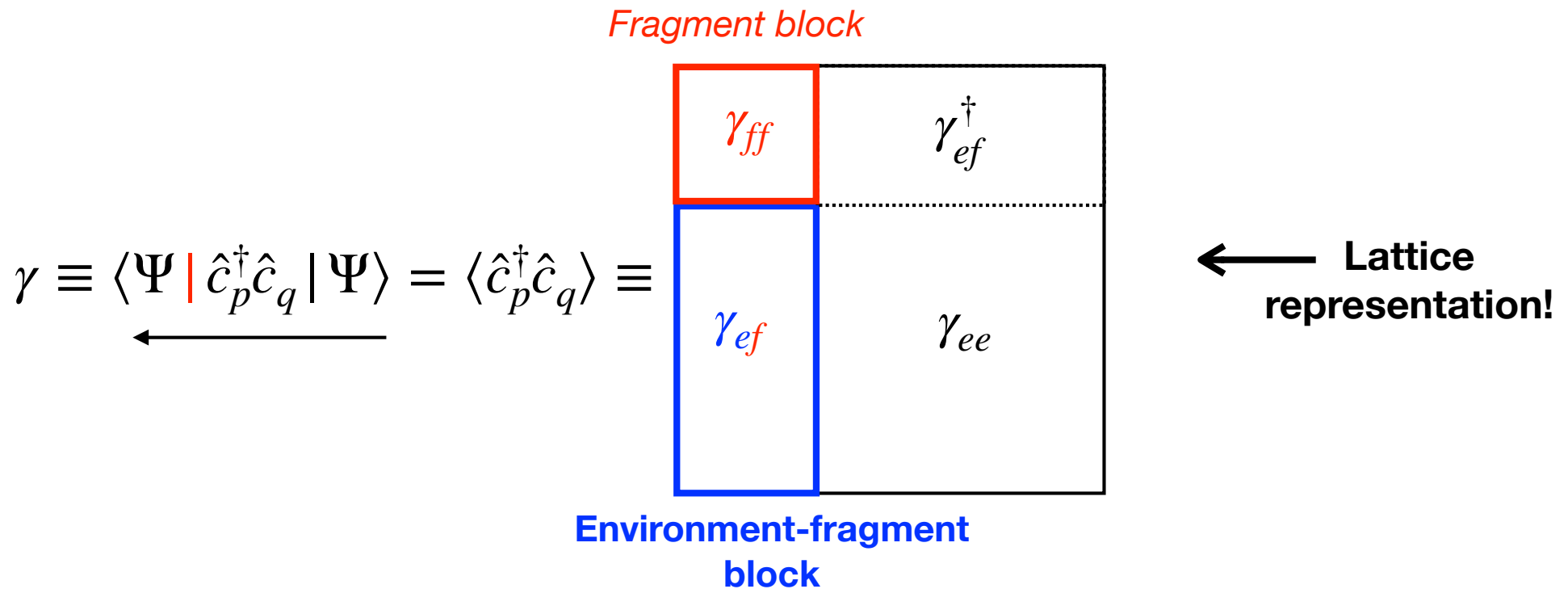
Quantum bath seen as a functional of the density matrix (1RDM)



Quantum bath seen as a functional of the density matrix (1RDM)



Quantum bath seen as a functional of the density matrix (1RDM)



Quantum bath seen as a functional of the density matrix (1RDM)

$$\gamma = \begin{array}{|c|c|} \hline \gamma_{ff} & \gamma_{ef}^\dagger \\ \hline \gamma_{ef} & \gamma_{ee} \\ \hline \end{array}$$

Env.-fragment block

$$\mathcal{U} = \begin{array}{|c|c|c|} \hline 1_{ff} & 0_{fe} & \\ \hline 0_{ef} & \mathcal{U}_{eb} & \mathcal{U}_{e\mathcal{E}} \\ \hline & & \underbrace{\hspace{2cm}} \\ \hline \end{array}$$

Implicit (but much simpler) definition of the cluster's environment:

$$\gamma_{ef}^\dagger \mathcal{U}_{e\mathcal{E}} \equiv 0$$

Cluster's environment

Will be justified later on...

Quantum bath seen as a functional of the density matrix (1RDM)

$$\gamma = \begin{array}{|c|c|} \hline \gamma_{ff} & \gamma_{ef}^\dagger \\ \hline \gamma_{ef} & \gamma_{ee} \\ \hline \end{array}$$

Env.-fragment block

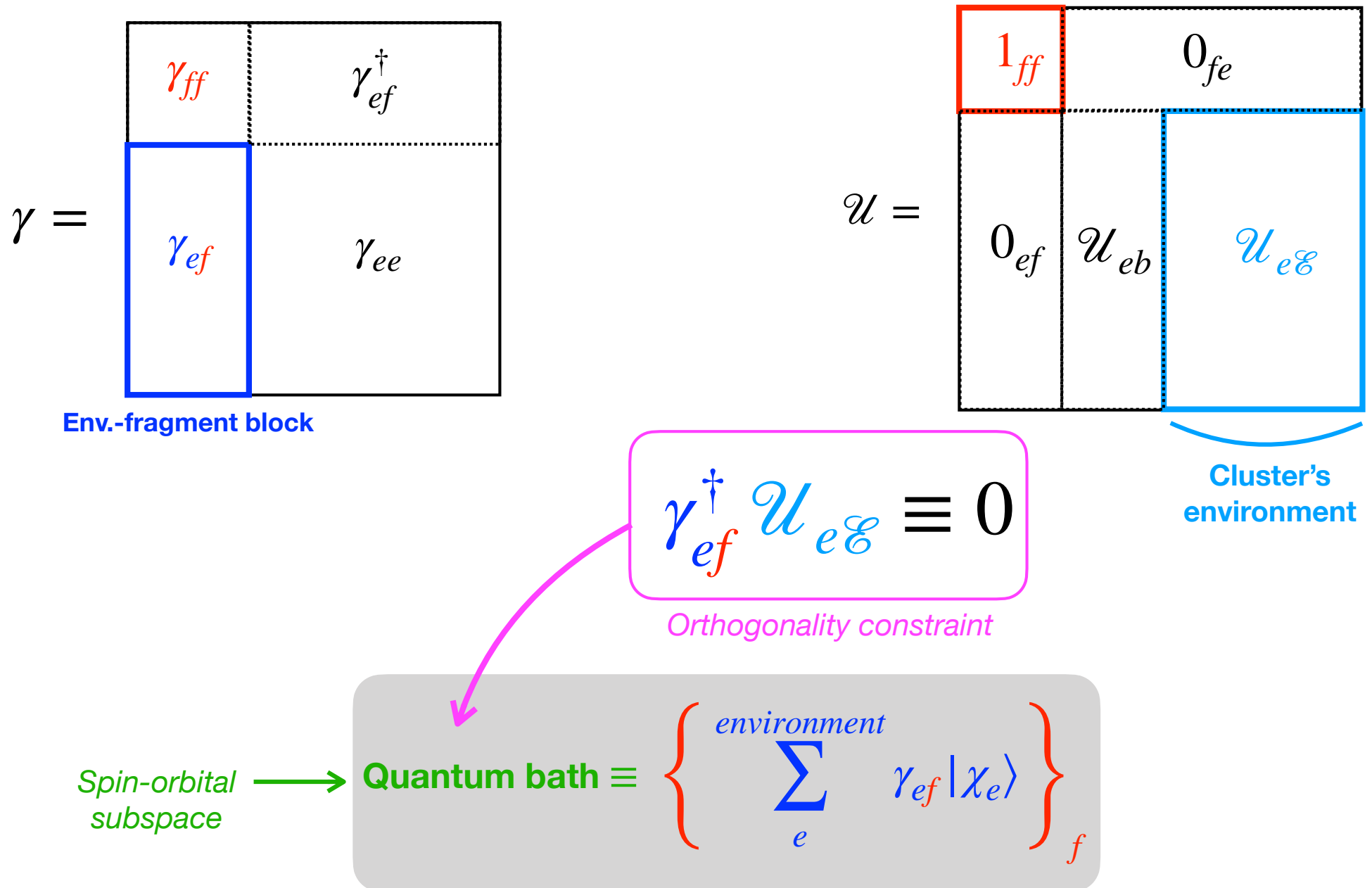
$$\mathcal{U} = \begin{array}{|c|c|c|} \hline 1_{ff} & 0_{fe} & \\ \hline 0_{ef} & \mathcal{U}_{eb} & \mathcal{U}_{e\mathcal{E}} \\ \hline & & \underbrace{\hspace{2cm}} \\ \hline \end{array}$$

Cluster's environment

$$\gamma_{ef}^\dagger \mathcal{U}_{e\mathcal{E}} \equiv 0$$

Orthogonality constraint

Quantum bath seen as a functional of the density matrix (1RDM)



Quantum bath seen as a functional of the density matrix (1RDM)

$$\gamma = \begin{array}{|c|c|} \hline \gamma_{ff} & \gamma_{ef}^\dagger \\ \hline \gamma_{ef} & \gamma_{ee} \\ \hline \end{array}$$

Env.-fragment block

$$\text{Quantum bath} \equiv \left\{ \sum_e^{\text{environment}} \gamma_{ef} |\chi_e\rangle \right\}_f$$

Quantum bath seen as a functional of the density matrix (1RDM)

$$\gamma = \begin{array}{|c|c|} \hline \gamma_{ff} & \gamma_{ef}^\dagger \\ \hline \gamma_{ef} & \gamma_{ee} \\ \hline \end{array}$$

Env.-fragment block

In principle as many bath spin-orbitals
as the dimension of the fragment
(number of "impurities")

$$\text{Quantum bath} \equiv \left\{ \sum_e^{\text{environment}} \gamma_{ef} |\chi_e\rangle \right\}_f$$



Quantum bath seen as a functional of the density matrix (1RDM)

$$\gamma = \begin{array}{|c|c|} \hline \gamma_{ff} & \gamma_{ef}^\dagger \\ \hline \gamma_{ef} & \gamma_{ee} \\ \hline \end{array}$$

Env.-fragment block

To-be orthonormalized
(SVD, Householder transformation, ...)

$$\text{Quantum bath} \equiv \left\{ \sum_e^{\text{environment}} \gamma_{ef} |\chi_e\rangle \right\}_f$$

Quantum bath seen as a functional of the density matrix (1RDM)

$$\gamma = \begin{array}{|c|c|} \hline \gamma_{ff} & \gamma_{ef}^\dagger \\ \hline \gamma_{ef} & \gamma_{ee} \\ \hline \end{array}$$

Env.-fragment block

$$\mathcal{U} = \begin{array}{|c|c|c|} \hline 1_{ff} & 0_{fe} & \\ \hline 0_{ef} & \mathcal{U}_{eb} & \mathcal{U}_{e\mathcal{E}} \\ \hline & & \mathcal{U}_{e\mathcal{E}} \\ \hline \end{array}$$

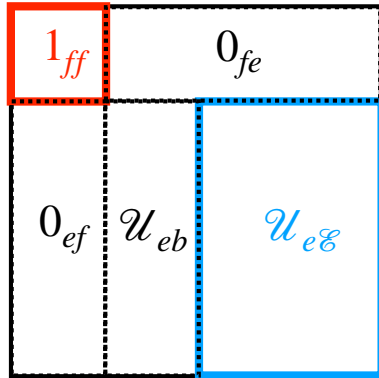
Cluster's environment

$$\gamma_{ef}^\dagger \mathcal{U}_{e\mathcal{E}} \equiv 0$$

Orthogonality constraint



Unitary transformed density matrix



$$\left\{ \langle \hat{d}_p^\dagger \hat{d}_q \rangle \right\} \equiv \tilde{\gamma} = \mathcal{U}^\dagger \gamma \mathcal{U} \equiv$$

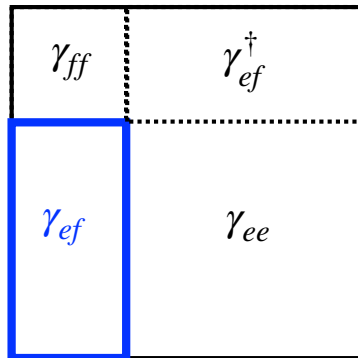
New embedding
representation

Let's visualize the **clusterization** in the 1RDM...



$$\left\{ \langle \hat{c}_p^\dagger \hat{c}_q \rangle \right\} \equiv$$

Original lattice
representation



Unitary transformed density matrix

“**fragment+bath**”
embedding cluster

Orthogonality constraint

$$\tilde{\gamma} = \mathcal{U}^\dagger \gamma \mathcal{U} \equiv$$

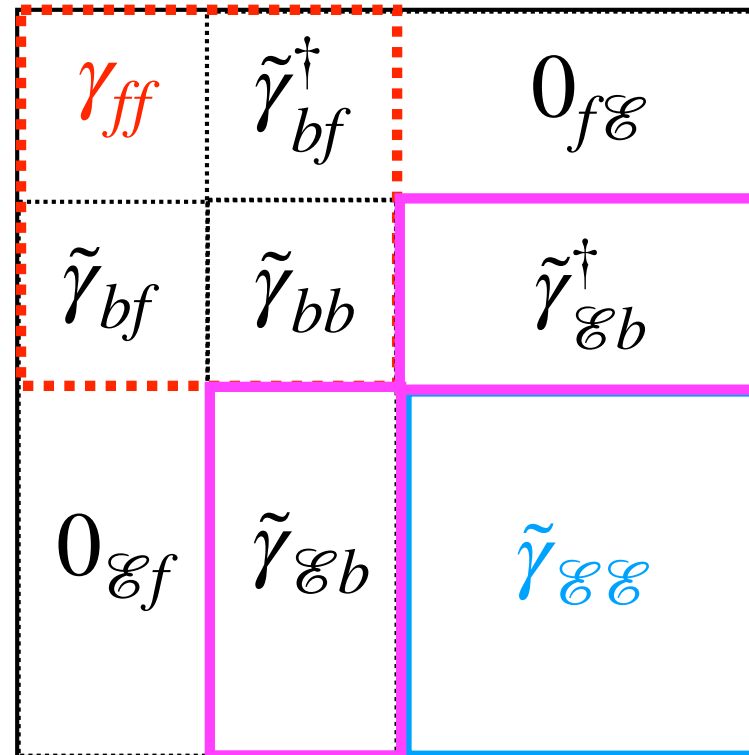
γ_{ff}	$\tilde{\gamma}_{bf}^\dagger$	$\gamma_{ef}^\dagger \mathcal{U}_{e\mathcal{E}} = 0$
$\tilde{\gamma}_{bf}$	$\tilde{\gamma}_{bb}$	$\tilde{\gamma}_{\mathcal{E}b}^\dagger$
$\mathcal{U}_{e\mathcal{E}}^\dagger \gamma_{ef}$ 0	$\tilde{\gamma}_{\mathcal{E}b}$	$\tilde{\gamma}_{\mathcal{E}\mathcal{E}}$

Orthogonality constraint

Unitary transformed density matrix

“**fragment**+bath”
embedding cluster

$$\tilde{\gamma} = \mathcal{U}^\dagger \gamma \mathcal{U} \equiv$$

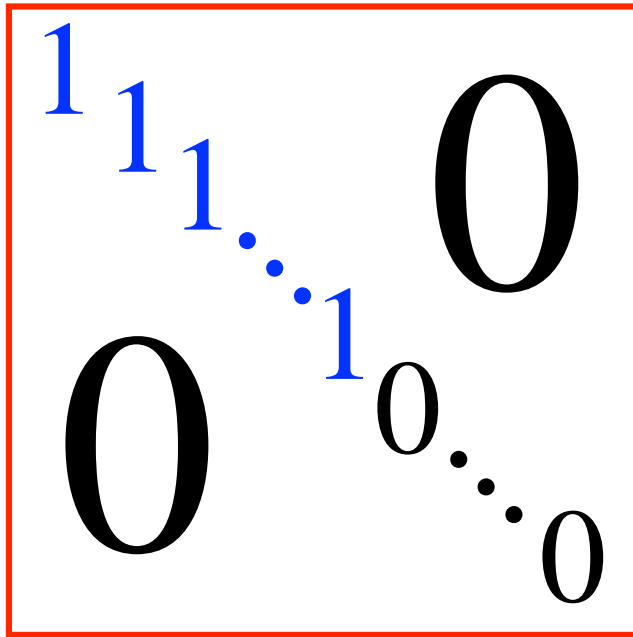


Cluster's
environment

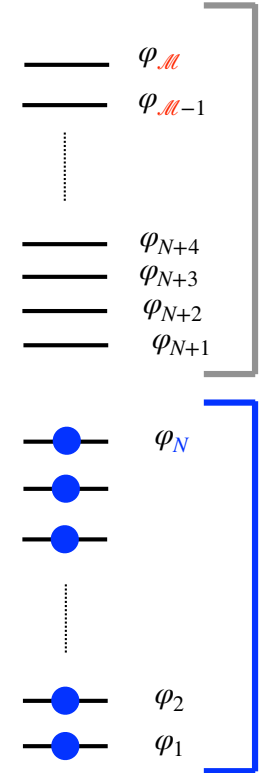
Entanglement

What if the full-system density matrix is idempotent?

$\gamma \equiv$

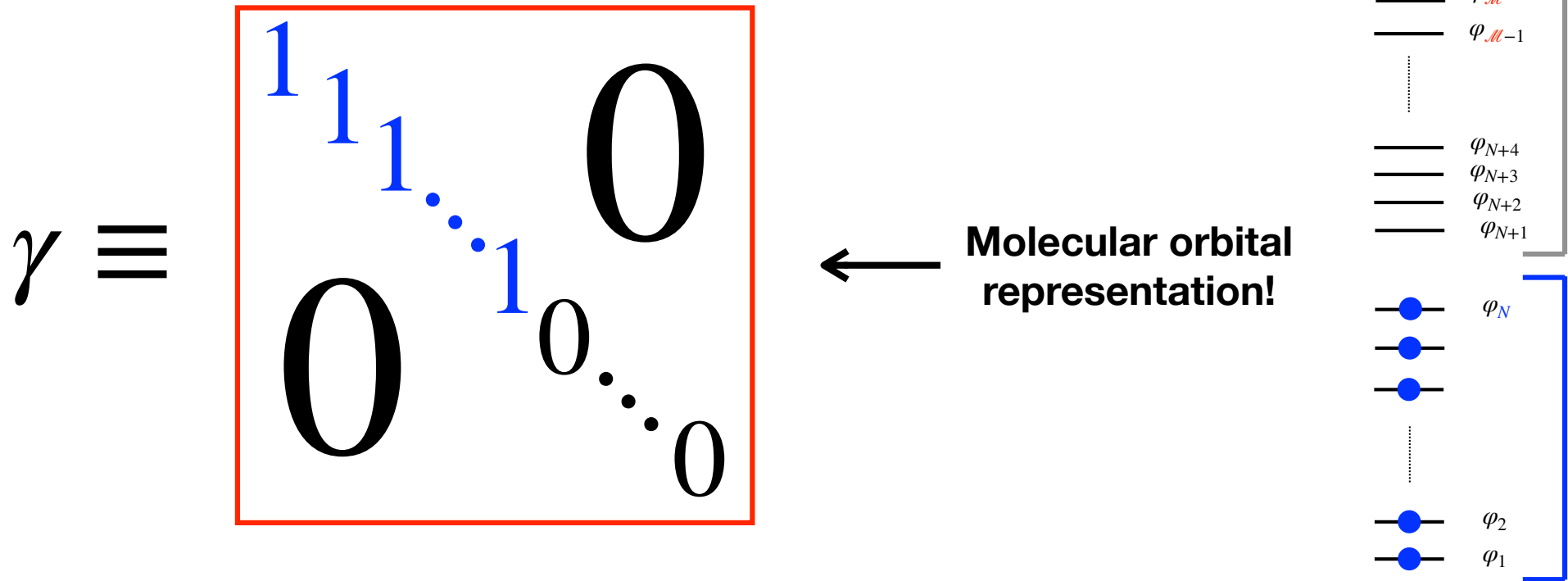


← **Molecular orbital
representation!**



*Mean-field (HF)
or Kohn-Sham DFT*

What if the full-system density matrix is idempotent?



Note that $\text{Tr } \gamma = N$ ← Total number of electrons (in the full system)

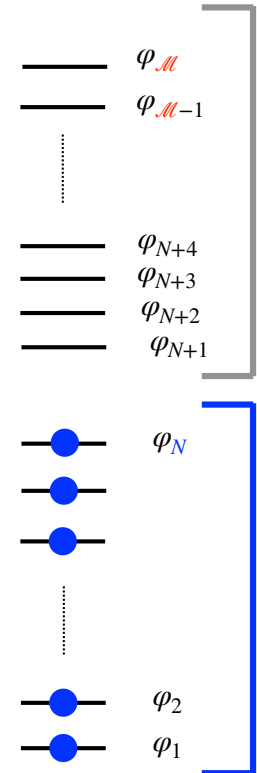
Mean-field (HF)
or Kohn-Sham DFT

What if the full-system density matrix is *idempotent*?

$$\gamma \equiv \begin{array}{cccc} 1 & & & 0 \\ & 1 & & \\ & & 1 & \\ & & & \dots \\ & & & & 1 & & \\ 0 & & & & & & 0 \\ & & & & & \dots & \\ & & & & & & & 0 \end{array}$$

$$\equiv \gamma^2$$

← Molecular orbital representation!

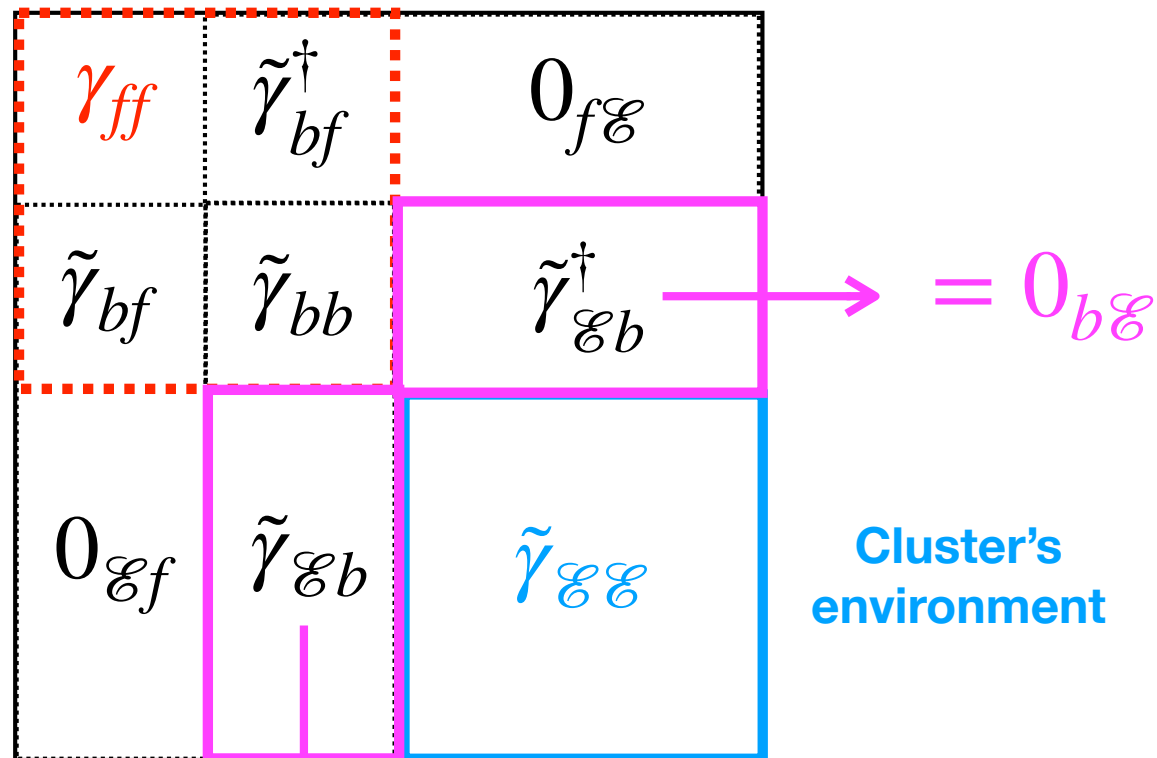


*Mean-field (HF)
or Kohn-Sham DFT*

Unitary transformed density matrix

“**fragment**+bath”
embedding cluster

$$\tilde{\gamma} = \mathcal{U}^\dagger \gamma \mathcal{U} \equiv$$



if $\gamma^2 = \gamma$ then
idempotency

$$0_{Ef} \tilde{\gamma}_{bf}^{-1} = 0_{Eb}$$

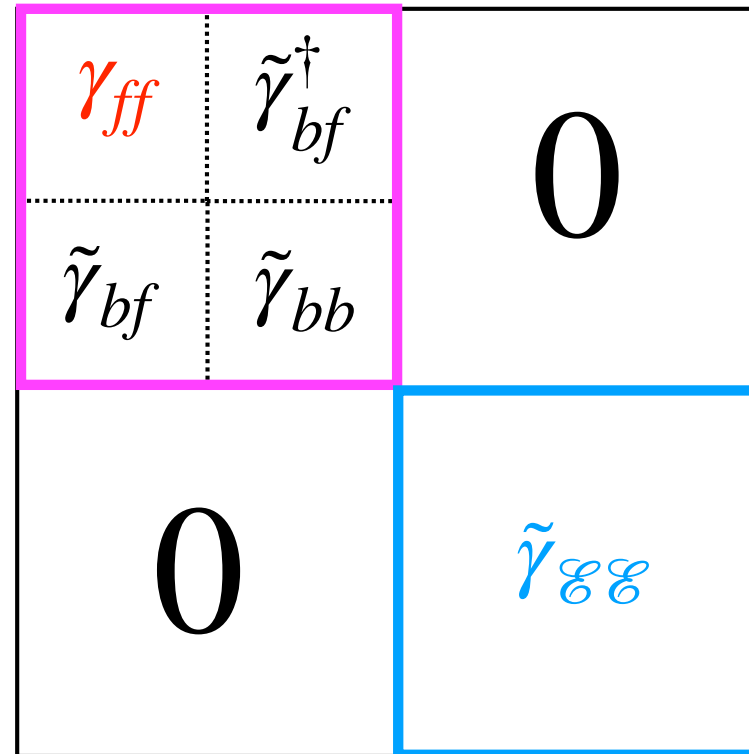
Unitary transformed density matrix

Disconnected
embedding cluster



$$\tilde{\gamma}^2 = \tilde{\gamma} = \mathcal{U}^\dagger \gamma \mathcal{U} \equiv$$

idempotency



Cluster's
environment

Unitary transformed density matrix

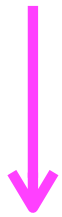
$$\gamma_{ff} + \tilde{\gamma}_{bf}^{-1} \tilde{\gamma}_{bb} \tilde{\gamma}_{bf} = \mathbf{1}_{ff} \leftarrow \begin{array}{l} \gamma^2 = \gamma \\ \text{idempotency} \end{array}$$

	L_f	L_f	
γ_{ff}	$\tilde{\gamma}_{bf}^\dagger$	$\mathbf{0}$	
$\tilde{\gamma}_{bf}$	$\tilde{\gamma}_{bb}$		
$\mathbf{0}$		$\tilde{\gamma}_{\mathcal{E}\mathcal{E}}$	

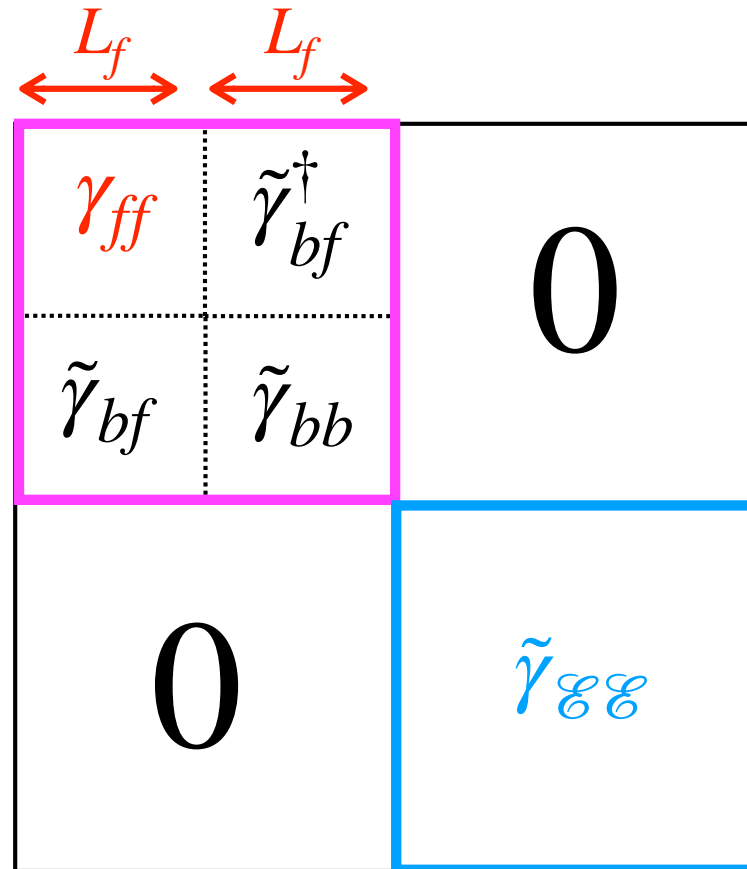
Cluster's environment

Unitary transformed density matrix

$$\gamma_{ff} + \tilde{\gamma}_{bf}^{-1} \tilde{\gamma}_{bb} \tilde{\gamma}_{bf} = \mathbf{1}_{ff} \leftarrow \begin{array}{l} \gamma^2 = \gamma \\ \text{idempotency} \end{array}$$



$$\text{Tr}[\gamma_{ff}] + \text{Tr}[\tilde{\gamma}_{bb}] = L_f$$



Cluster's
environment

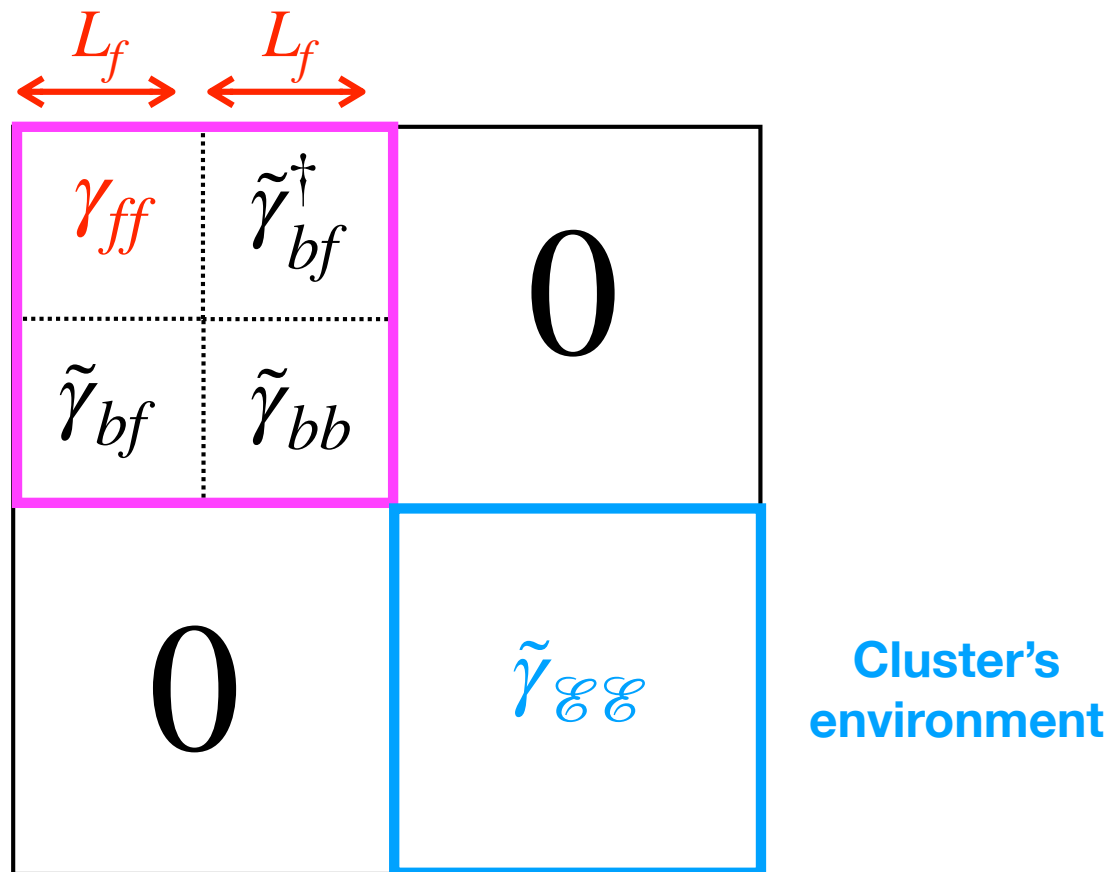
Unitary transformed density matrix

$$\gamma_{ff} + \tilde{\gamma}_{bf}^{-1} \tilde{\gamma}_{bb} \tilde{\gamma}_{bf} = \mathbf{1}_{ff} \leftarrow \begin{array}{l} \gamma^2 = \gamma \\ \text{idempotency} \end{array}$$



$$\text{Tr}[\gamma_{ff}] + \text{Tr}[\tilde{\gamma}_{bb}] = L_f$$

The number of electrons in the cluster equals the number of embedded impurities



Starting a DMET calculation...

Density matrix of the **full system**

$$\gamma = \begin{array}{|c|c|} \hline \gamma_{ff} & \gamma_{ef}^\dagger \\ \hline \gamma_{ef} & \gamma_{ee} \\ \hline \end{array} = ???$$

Starting a DMET calculation...

Density matrix of the **full system**

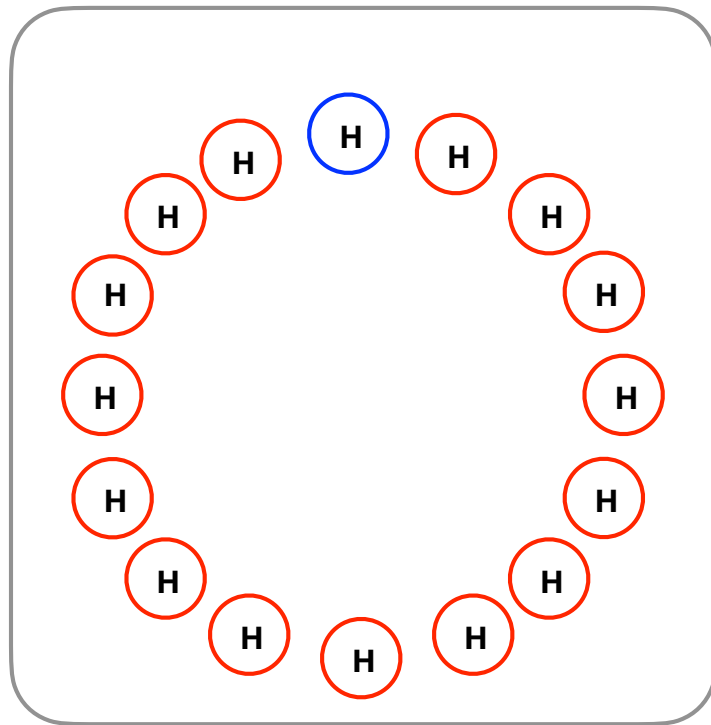
$$\gamma = \begin{array}{|c|c|} \hline \gamma_{ff} & \gamma_{ef}^\dagger \\ \hline \gamma_{ef} & \gamma_{ee} \\ \hline \end{array}$$

Mean-field evaluation in practice

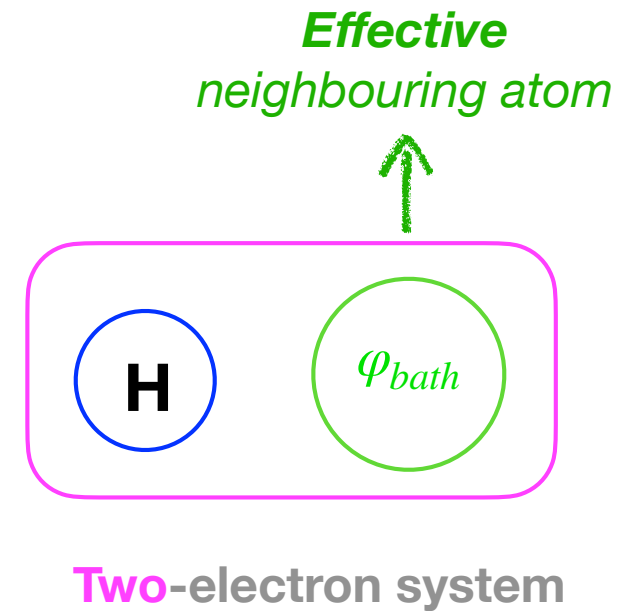
Idempotent ($\gamma^2 = \gamma$)

Illustrative example

Rings of hydrogen atoms (Hubbard model)



N -electron system



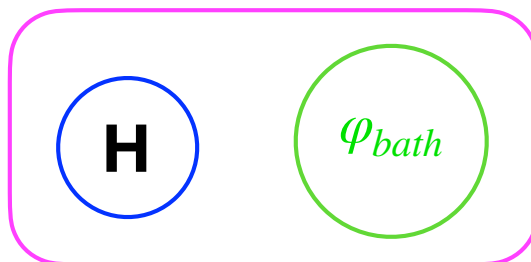
$$\hat{H} = \sum_{\sigma=\uparrow,\downarrow} \sum_{i=0}^{L-1} -t \left(\hat{c}_{i\sigma}^\dagger \hat{c}_{(i+1)\sigma} + \hat{c}_{(i+1)\sigma}^\dagger \hat{c}_{i\sigma} \right) + U \sum_{i=0}^{L-1} \hat{c}_{i\uparrow}^\dagger \hat{c}_{i\downarrow}^\dagger \hat{c}_{i\downarrow} \hat{c}_{i\uparrow}$$

Approximate embedding of interacting Hamiltonians

One-electron Hamiltonian
of the cluster



\hat{h}^c



Exact non-interacting (i.e., for $U = 0$)
embedding

S. Wouters, C. A. Jiménez-Hoyos, Q. Sun, and G. K.-L. Chan, *J. Chem. Theory Comput.* **12**, 2706 (2016).

S. Sekaran, M. Tsuchiizu, M. Saubanère, and E. Fromager, *Phys. Rev. B* **104**, 035121 (2021).

S. Sekaran, M. Saubanère, and E. Fromager, *Computation* **2022**, 10, 45.

Approximate embedding of interacting Hamiltonians

One-electron Hamiltonian
of the cluster



$$\hat{h}^{\mathcal{C}} \text{ single impurity}$$

→

$$\hat{h}^{\mathcal{C}} + \langle pp | pp \rangle \hat{c}_{p\uparrow}^\dagger \hat{c}_{p\downarrow}^\dagger \hat{c}_{p\downarrow} \hat{c}_{p\uparrow}$$

Two-electron repulsion
on the impurity



Exact non-interacting
embedding

Approximate embedding of interacting Hamiltonians

One-electron Hamiltonian
of the cluster

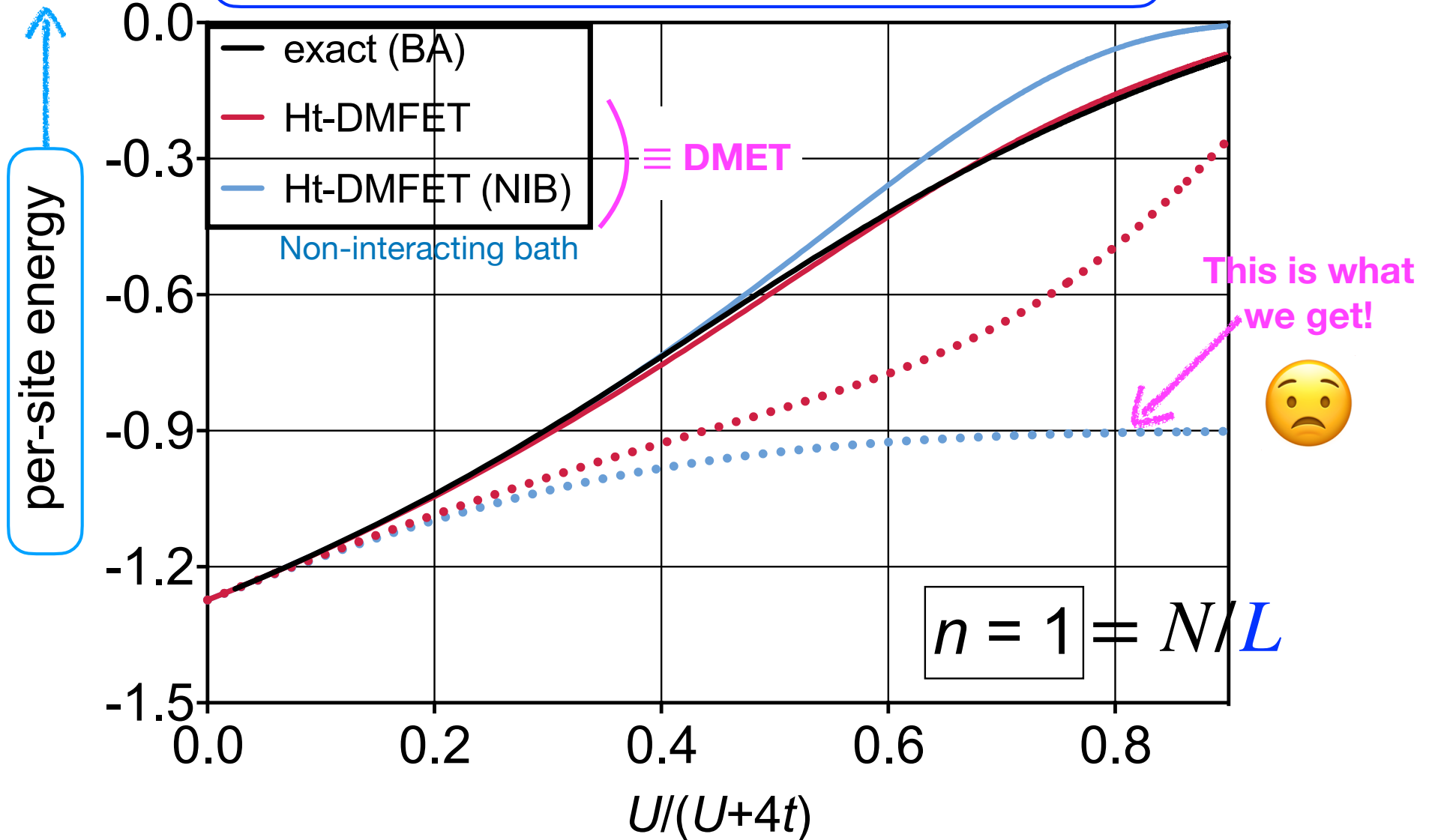
$$\hat{h}^{\mathcal{C}} \xrightarrow{\text{single impurity}} \hat{h}^{\mathcal{C}} + \overbrace{\langle pp | pp \rangle}^U \hat{c}_{p\uparrow}^\dagger \hat{c}_{p\downarrow}^\dagger \hat{c}_{p\downarrow} \hat{c}_{p\uparrow}$$

Exact non-interacting
embedding

Half-filled uniform Hubbard ring with $L = 400$ atomic sites

E/L

$$\hat{H} = \sum_{\sigma=\uparrow,\downarrow} \sum_{i=0}^{L-1} -t \left(\hat{c}_{i\sigma}^\dagger \hat{c}_{(i+1)\sigma} + \hat{c}_{(i+1)\sigma}^\dagger \hat{c}_{i\sigma} \right) + U \sum_{i=0}^{L-1} \hat{c}_{i\uparrow}^\dagger \hat{c}_{i\downarrow}^\dagger \hat{c}_{i\downarrow} \hat{c}_{i\uparrow}$$



Approximate embedding of interacting Hamiltonians

One-electron Hamiltonian
of the cluster

Chemical potential
on the impurity

$$\hat{h}^{\mathcal{C}} \xrightarrow{\text{single impurity}} \hat{h}^{\mathcal{C}} + \overbrace{\langle pp | pp \rangle}^U \hat{c}_{p\uparrow}^\dagger \hat{c}_{p\downarrow}^\dagger \hat{c}_{p\downarrow} \hat{c}_{p\uparrow} - \tilde{\mu}^{imp} \sum_{\sigma=\uparrow,\downarrow} \hat{c}_{p\sigma}^\dagger \hat{c}_{p\sigma}$$

Exact non-interacting
embedding

Approximate embedding of interacting Hamiltonians

One-electron Hamiltonian
of the cluster

$$\hat{h}^c$$

single impurity
→

$$\hat{h}^c$$

$$+ \overbrace{\langle pp | pp \rangle}^U \hat{c}_{p\uparrow}^\dagger \hat{c}_{p\downarrow}^\dagger \hat{c}_{p\downarrow} \hat{c}_{p\uparrow}$$

Chemical potential
on the impurity

$$- \tilde{\mu}^{imp}$$

$$\sum_{\sigma=\uparrow,\downarrow} \hat{c}_{p\sigma}^\dagger \hat{c}_{p\sigma}$$

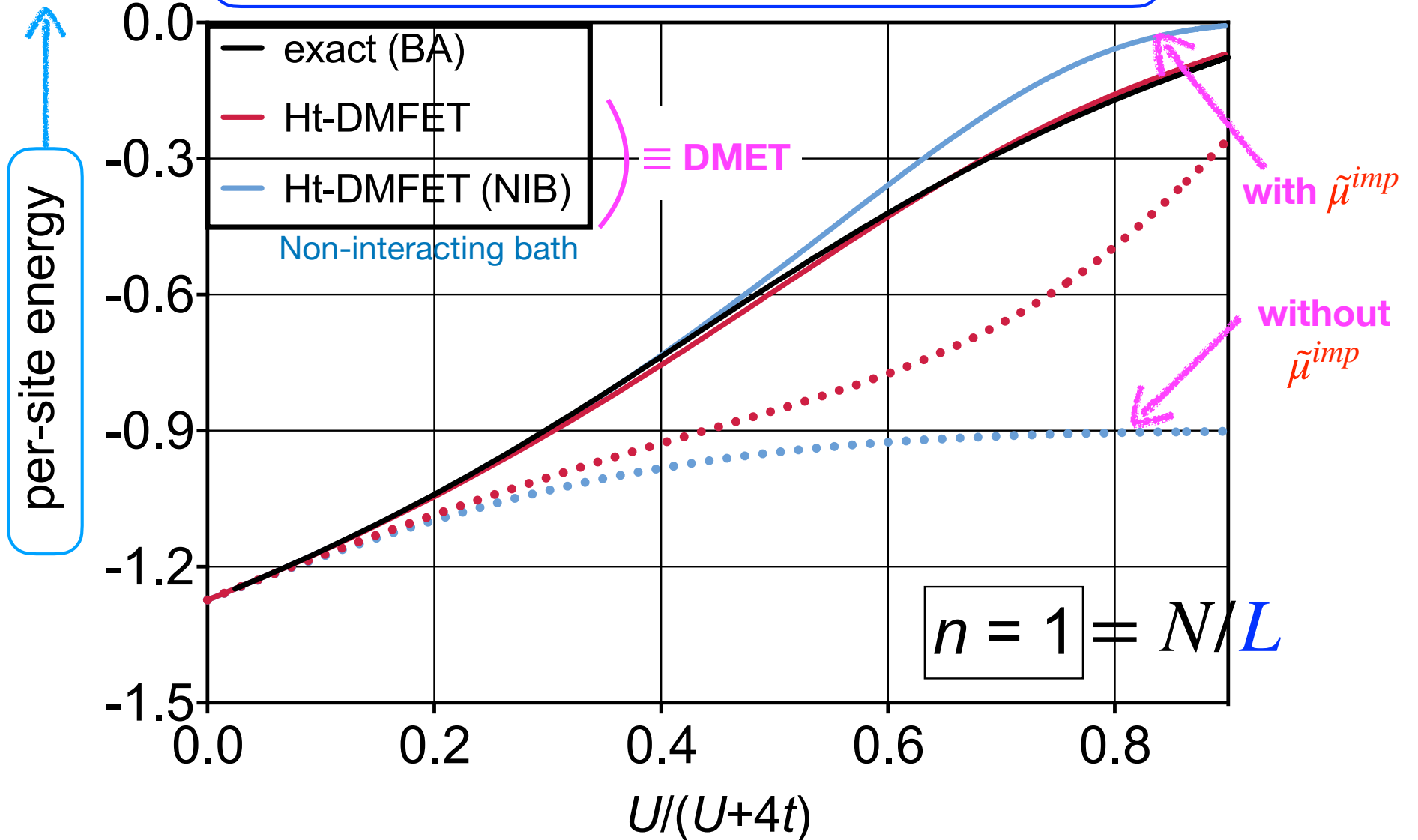
Exact non-interacting
embedding

Approximate interacting
embedding

Half-filled uniform Hubbard ring with $L = 400$ atomic sites

E/L

$$\hat{H} = \sum_{\sigma=\uparrow,\downarrow} \sum_{i=0}^{L-1} -t \left(\hat{c}_{i\sigma}^\dagger \hat{c}_{(i+1)\sigma} + \hat{c}_{(i+1)\sigma}^\dagger \hat{c}_{i\sigma} \right) + U \sum_{i=0}^{L-1} \hat{c}_{i\uparrow}^\dagger \hat{c}_{i\downarrow}^\dagger \hat{c}_{i\downarrow} \hat{c}_{i\uparrow}$$



Self-consistency in DMET
(for a single impurity and a uniform full-size system)


Fixing the number of electrons versus fixing the chemical potential

$$\hat{H} = \sum_{\sigma=\uparrow,\downarrow} \sum_{i=0}^{L-1} -t \left(\hat{c}_{i\sigma}^\dagger \hat{c}_{(i+1)\sigma} + \hat{c}_{(i+1)\sigma}^\dagger \hat{c}_{i\sigma} \right) + U \sum_{i=0}^{L-1} \hat{c}_{i\uparrow}^\dagger \hat{c}_{i\downarrow}^\dagger \hat{c}_{i\downarrow} \hat{c}_{i\uparrow}$$

... and we fix the number of electrons in the system

Fixing the number of electrons versus fixing the chemical potential

$$\hat{H} = \sum_{\sigma=\uparrow,\downarrow} \sum_{i=0}^{L-1} -t \left(\hat{c}_{i\sigma}^\dagger \hat{c}_{(i+1)\sigma} + \hat{c}_{(i+1)\sigma}^\dagger \hat{c}_{i\sigma} \right) + U \sum_{i=0}^{L-1} \hat{c}_{i\uparrow}^\dagger \hat{c}_{i\downarrow}^\dagger \hat{c}_{i\downarrow} \hat{c}_{i\uparrow}$$


$$\hat{H} - \mu \sum_{\sigma=\uparrow,\downarrow} \sum_{i=0}^{L-1} \hat{c}_{i\sigma}^\dagger \hat{c}_{i\sigma}$$

“Grand-canonical” Hamiltonian

Chemical potential

≡ uniform external potential

Fixing the number of electrons versus fixing the chemical potential

$$\hat{H} = \sum_{\sigma=\uparrow,\downarrow} \sum_{i=0}^{L-1} -t \left(\hat{c}_{i\sigma}^\dagger \hat{c}_{(i+1)\sigma} + \hat{c}_{(i+1)\sigma}^\dagger \hat{c}_{i\sigma} \right) + U \sum_{i=0}^{L-1} \hat{c}_{i\uparrow}^\dagger \hat{c}_{i\downarrow}^\dagger \hat{c}_{i\downarrow} \hat{c}_{i\uparrow}$$

$$\hat{H} - \mu \sum_{\sigma=\uparrow,\downarrow} \sum_{i=0}^{L-1} \hat{c}_{i\sigma}^\dagger \hat{c}_{i\sigma}$$

“Grand-canonical” Hamiltonian

Chemical potential

≡ uniform external potential

Uniform density profile (twice the filling):

$$n = \left\langle \sum_{\sigma=\uparrow,\downarrow} \hat{c}_{i\sigma}^\dagger \hat{c}_{i\sigma} \right\rangle = \frac{N}{L}$$

≡ $N(\mu)$: Total number of electrons
 ↓
 N
 ← Total number of sites

Self-consistency through a connection with density functional theory

$$\hat{H}(U) - \mu \sum_{\sigma=\uparrow,\downarrow} \sum_{i=0}^{L-1} \hat{c}_{i\sigma}^\dagger \hat{c}_{i\sigma}$$

True interacting Hamiltonian ($U \neq 0$)

Fixed chemical potential

$$n = n(\mu) = ?$$

Self-consistency through a connection with density functional theory

$$\hat{H}(U) - \mu \sum_{\sigma=\uparrow,\downarrow} \sum_{i=0}^{L-1} \hat{c}_{i\sigma}^\dagger \hat{c}_{i\sigma}$$

True interacting Hamiltonian ($U \neq 0$)



Fixed chemical potential

$$n = n(\mu) = ?$$

Self-consistency through a connection with density functional theory

$$\hat{H}(U) - \mu \sum_{\sigma=\uparrow,\downarrow} \sum_{i=0}^{L-1} \hat{c}_{i\sigma}^\dagger \hat{c}_{i\sigma}$$

True interacting Hamiltonian ($U \neq 0$)



Fixed chemical potential

$$n = n(\mu) = ?$$

$$\hat{H}(U=0) - \mu^{\text{KS}} \sum_{\sigma=\uparrow,\downarrow} \sum_{i=0}^{L-1} \hat{c}_{i\sigma}^\dagger \hat{c}_{i\sigma}$$

“Low-level” non-interacting full-size Hamiltonian that generates the bath through its ground-state idempotent density matrix

Unknown Kohn–Sham chemical potential

↔ Kohn–Sham full-size Hamiltonian

$$-\mu^{\text{KS}} = -\mu + v_{\text{Hxc}} \longleftarrow n^{\text{KS}} = n(\mu) = ?$$

Self-consistency through a connection with density functional theory

$$\hat{H}(U) - \mu \sum_{\sigma=\uparrow,\downarrow} \sum_{i=0}^{L-1} \hat{c}_{i\sigma}^\dagger \hat{c}_{i\sigma}$$

True interacting Hamiltonian ($U \neq 0$)



Fixed chemical potential

$$n = n(\mu) = ?$$

$$\hat{H}(U=0) - \mu^{\text{KS}} \sum_{\sigma=\uparrow,\downarrow} \sum_{i=0}^{L-1} \hat{c}_{i\sigma}^\dagger \hat{c}_{i\sigma}$$

“Low-level” non-interacting full-size Hamiltonian that generates the bath through its ground-state idempotent density matrix

Unknown Kohn–Sham chemical potential

↔ Kohn–Sham full-size Hamiltonian

$$-\mu^{\text{KS}} = -\mu + v_{\text{Hxc}} \longleftarrow n^{\text{KS}} = n(\mu) = ?$$



Self-consistency through a connection with density functional theory

$$\hat{H}(U) - \mu \sum_{\sigma=\uparrow,\downarrow} \sum_{i=0}^{L-1} \hat{c}_{i\sigma}^\dagger \hat{c}_{i\sigma}$$

True interacting Hamiltonian ($U \neq 0$)



Fixed chemical potential

$$n = n(\mu) = ?$$

$$\hat{H}(U=0) - \mu^{\text{KS}} \sum_{\sigma=\uparrow,\downarrow} \sum_{i=0}^{L-1} \hat{c}_{i\sigma}^\dagger \hat{c}_{i\sigma}$$

“Low-level” non-interacting full-size Hamiltonian that generates the bath through its ground-state idempotent density matrix

Unknown Kohn–Sham chemical potential

↔ Kohn–Sham full-size Hamiltonian

$$-\mu^{\text{KS}} = -\mu + v_{\text{Hxc}} \leftarrow n^{\text{KS}} = n(\mu) = ?$$



$$\hat{H}^{\mathcal{C}} = \hat{h}^{\mathcal{C}} + U \hat{c}_{p\uparrow}^\dagger \hat{c}_{p\downarrow}^\dagger \hat{c}_{p\downarrow} \hat{c}_{p\uparrow} - \tilde{\mu}^{\text{imp}} \sum_{\sigma=\uparrow,\downarrow} \hat{c}_{p\sigma}^\dagger \hat{c}_{p\sigma}$$

Impurity-interacting Hamiltonian of the two-electron embedding cluster

Self-consistency through a connection with density functional theory

$$\hat{H}(U) - \mu \sum_{\sigma=\uparrow,\downarrow} \sum_{i=0}^{L-1} \hat{c}_{i\sigma}^\dagger \hat{c}_{i\sigma}$$

True interacting Hamiltonian ($U \neq 0$)



Fixed chemical potential

$$n = n(\mu) = ?$$

$$\hat{H}(U=0) - \mu^{\text{KS}} \sum_{\sigma=\uparrow,\downarrow} \sum_{i=0}^{L-1} \hat{c}_{i\sigma}^\dagger \hat{c}_{i\sigma}$$

“Low-level” non-interacting full-size Hamiltonian that generates the bath through its ground-state idempotent density matrix

Unknown Kohn–Sham chemical potential

↔ Kohn–Sham full-size Hamiltonian

$$-\mu^{\text{KS}} = -\mu + v_{\text{Hxc}} \leftarrow n^{\text{KS}} = n(\mu) = ?$$



$$\hat{H}^{\text{e}} = \hat{h}^{\text{e}} + U \hat{c}_{p\uparrow}^\dagger \hat{c}_{p\downarrow}^\dagger \hat{c}_{p\downarrow} \hat{c}_{p\uparrow} - \tilde{\mu}^{\text{imp}} \sum_{\sigma=\uparrow,\downarrow} \hat{c}_{p\sigma}^\dagger \hat{c}_{p\sigma}$$

Impurity-interacting Hamiltonian of the two-electron embedding cluster



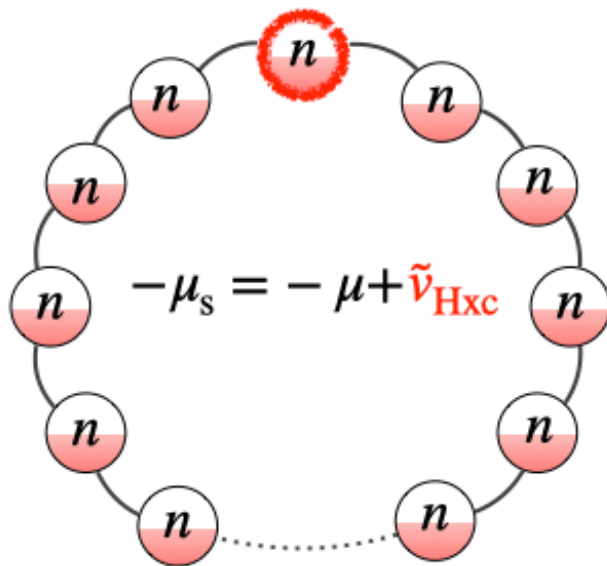
Local potential-functional embedding theory (LPFET)

$$v_{\text{Hxc}}(n) \approx \tilde{\mu}^{\text{imp}}(n)$$



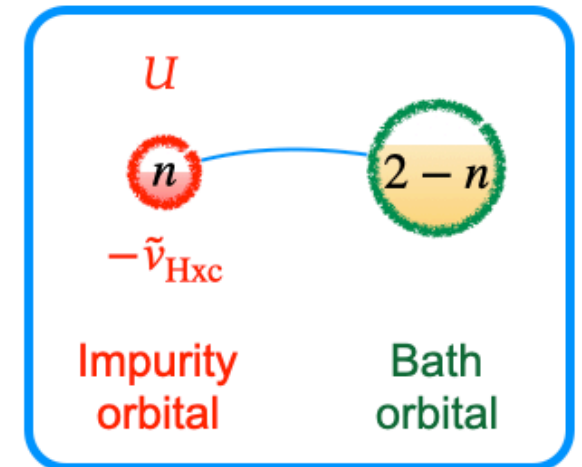
Self-consistency loop

Kohn-Sham lattice

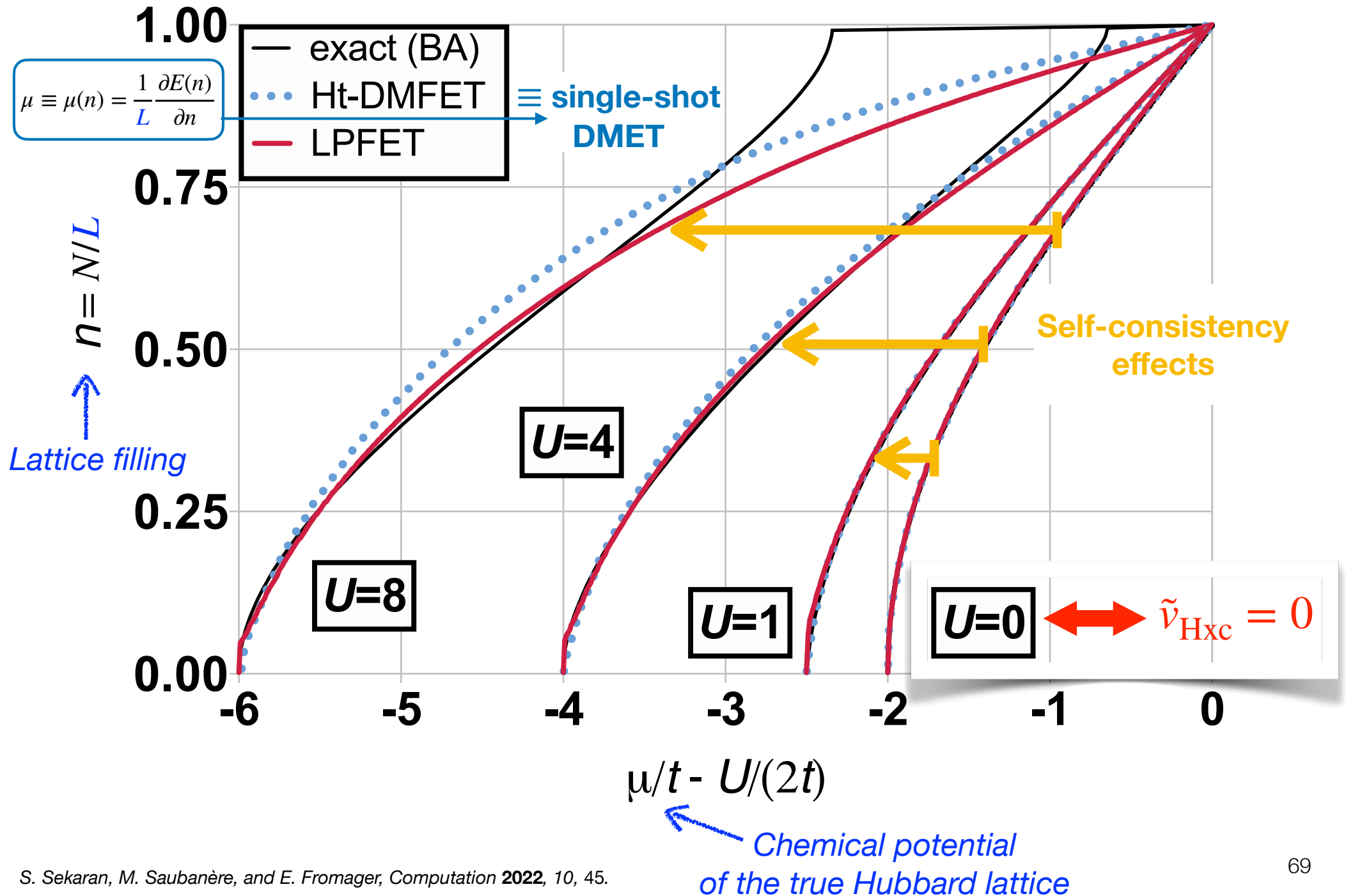


$$n_{\Phi_0^{\text{KS}}}[\tilde{v}_{\text{Hxc}}] = n_{\Psi_0^{\text{cluster}}}[\tilde{v}_{\text{Hxc}}]$$

Embedding cluster

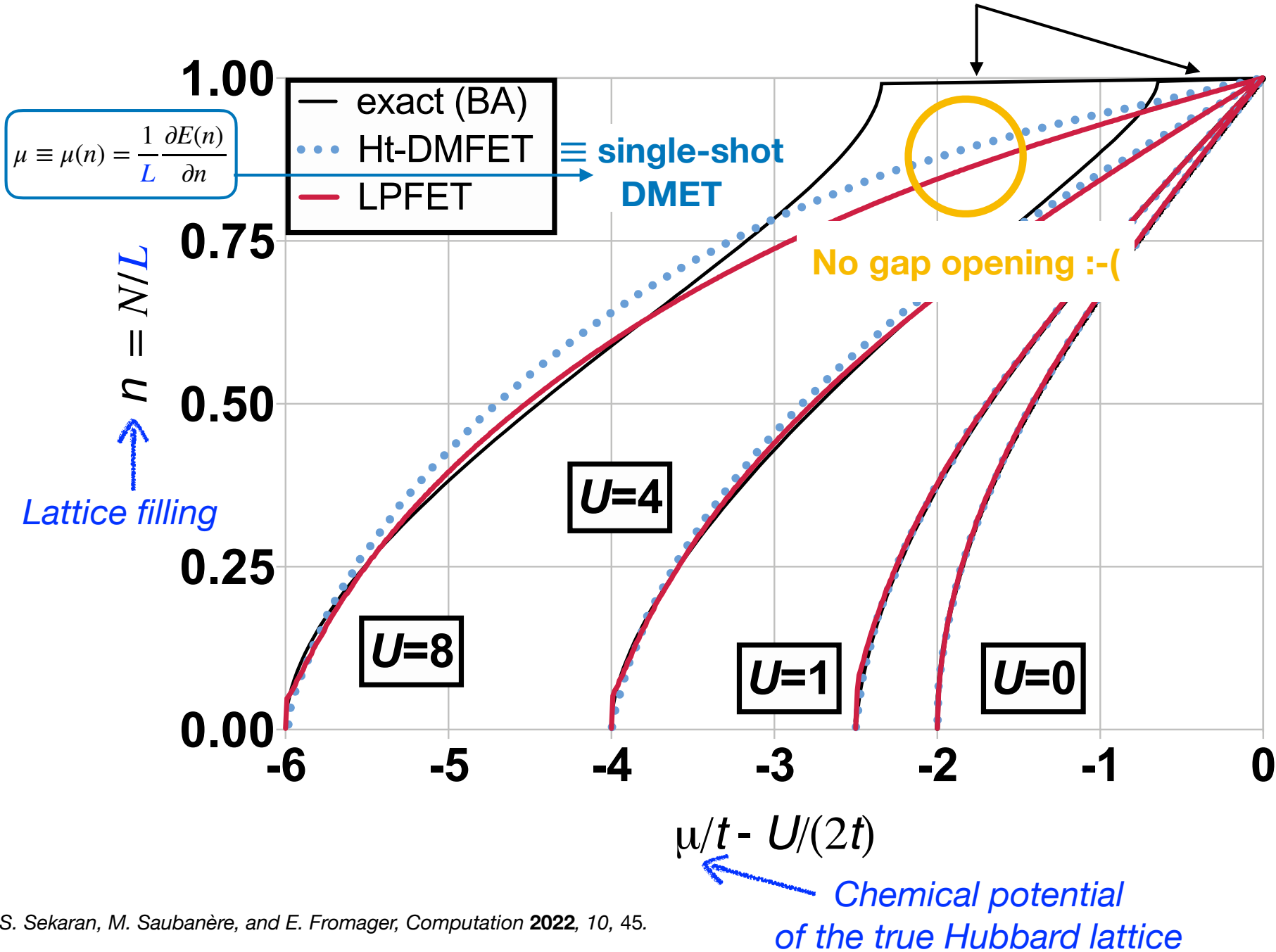


Local potential-functional embedding theory (LPFET)

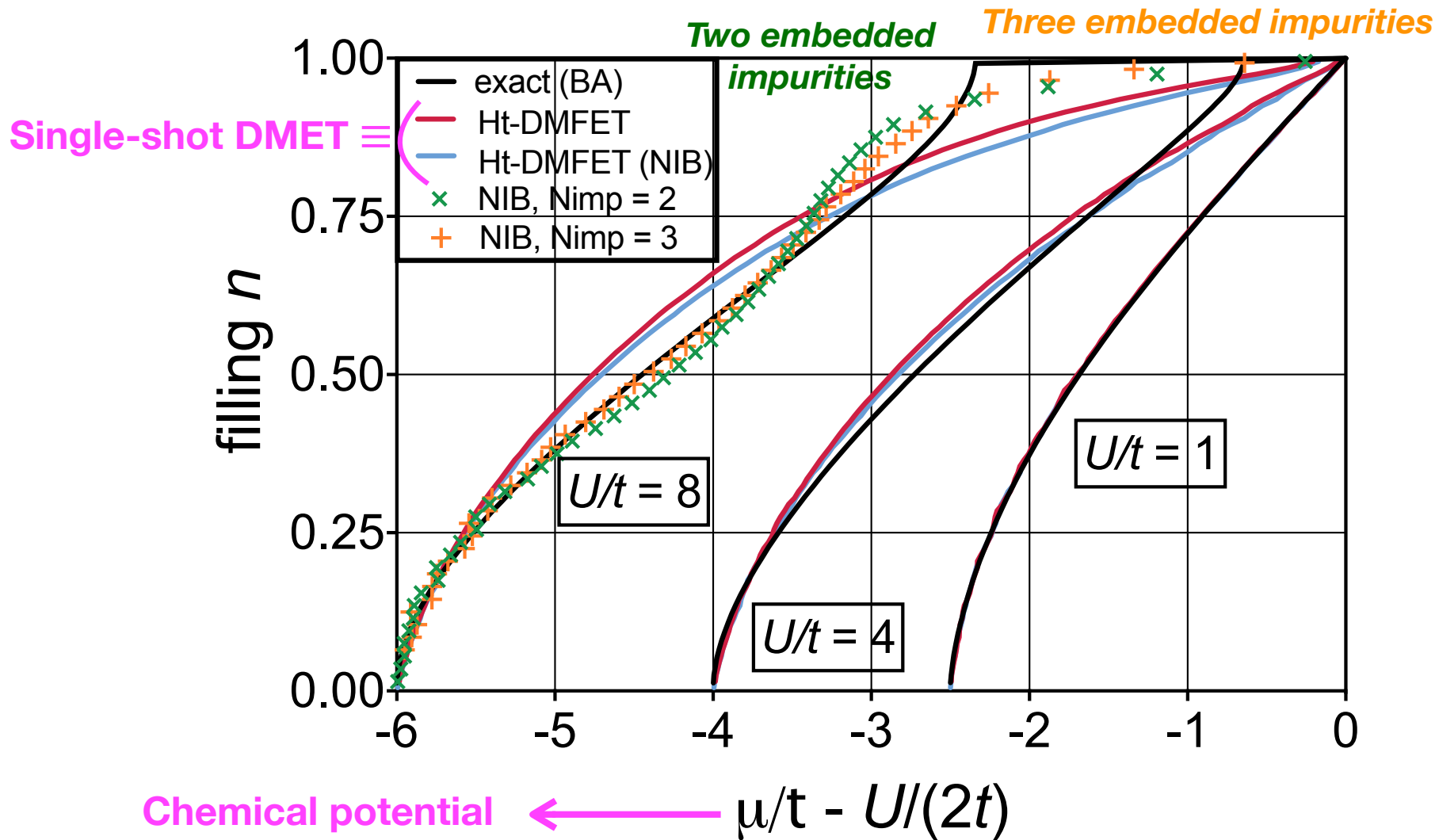


Local potential-functional embedding theory (LPFET)

Density-driven **Mott-Hubbard transition**






Mott-Hubbard density-driven transition and multiple impurities



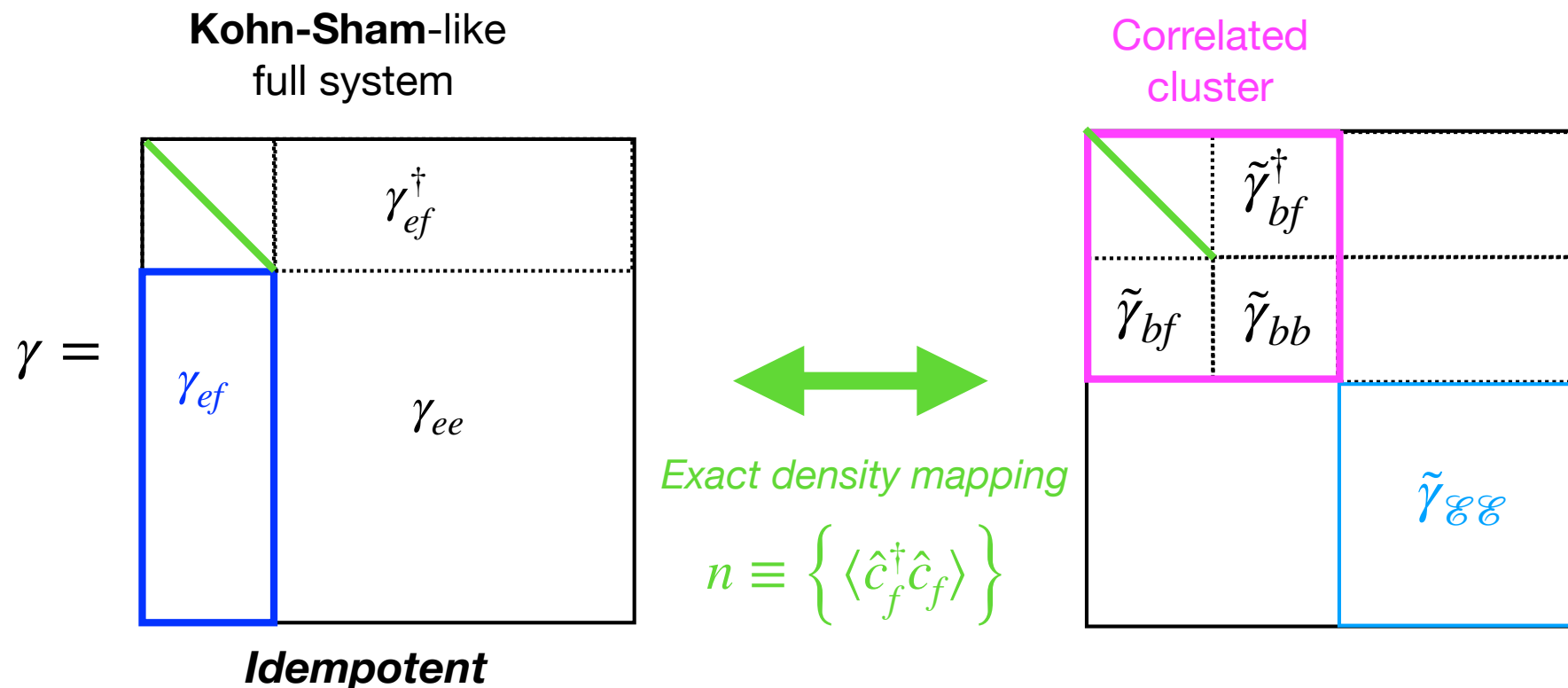
$$\mu \equiv \mu(n) = \frac{1}{L} \frac{\partial E(n)}{\partial n}$$

Article

Local Potential Functional Embedding Theory: A Self-Consistent Flavor of Density Functional Theory for Lattices without Density Functionals

Sajanthan Sekaran ^{1,*} , Matthieu Saubanère ²  and Emmanuel Fromager ¹ 

Citation: Sekaran, S.; Saubanère, M.; Fromager, E. Local Potential Functional Embedding Theory: A Self-Consistent Flavor of Density Functional Theory for Lattices without Density Functionals. *Computation* **2022**, *10*, 45. <https://doi.org/10.3390/computation10030045>

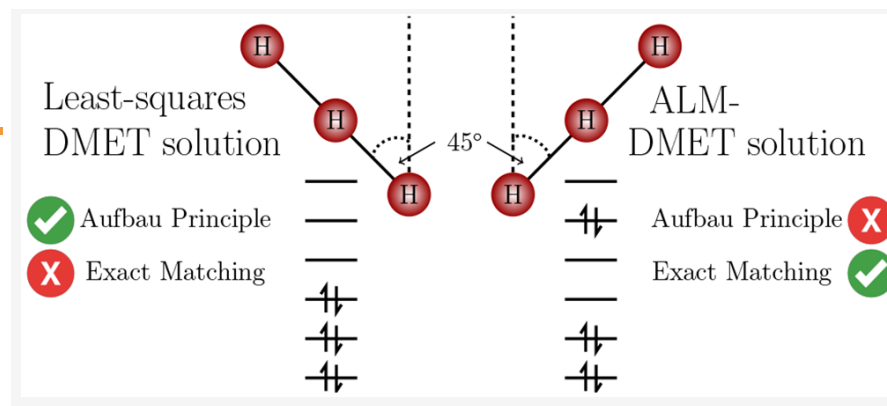


Pure State v -Representability of Density Matrix Embedding Theory

Fabian M. Faulstich,[⊥] Raehyun Kim,[⊥] Zhi-Hao Cui, Zaiwen Wen, Garnet Kin-Lic Chan, and Lin Lin*



Cite This: *J. Chem. Theory Comput.* 2022, 18, 851–864



Kohn-Sham-like
full system?

$$\gamma = \begin{array}{|c|c|} \hline \gamma_{ff} & \gamma_{ef}^\dagger \\ \hline \gamma_{ef} & \gamma_{ee} \\ \hline \end{array}$$

Idempotent

?

↔

Density matrix mapping

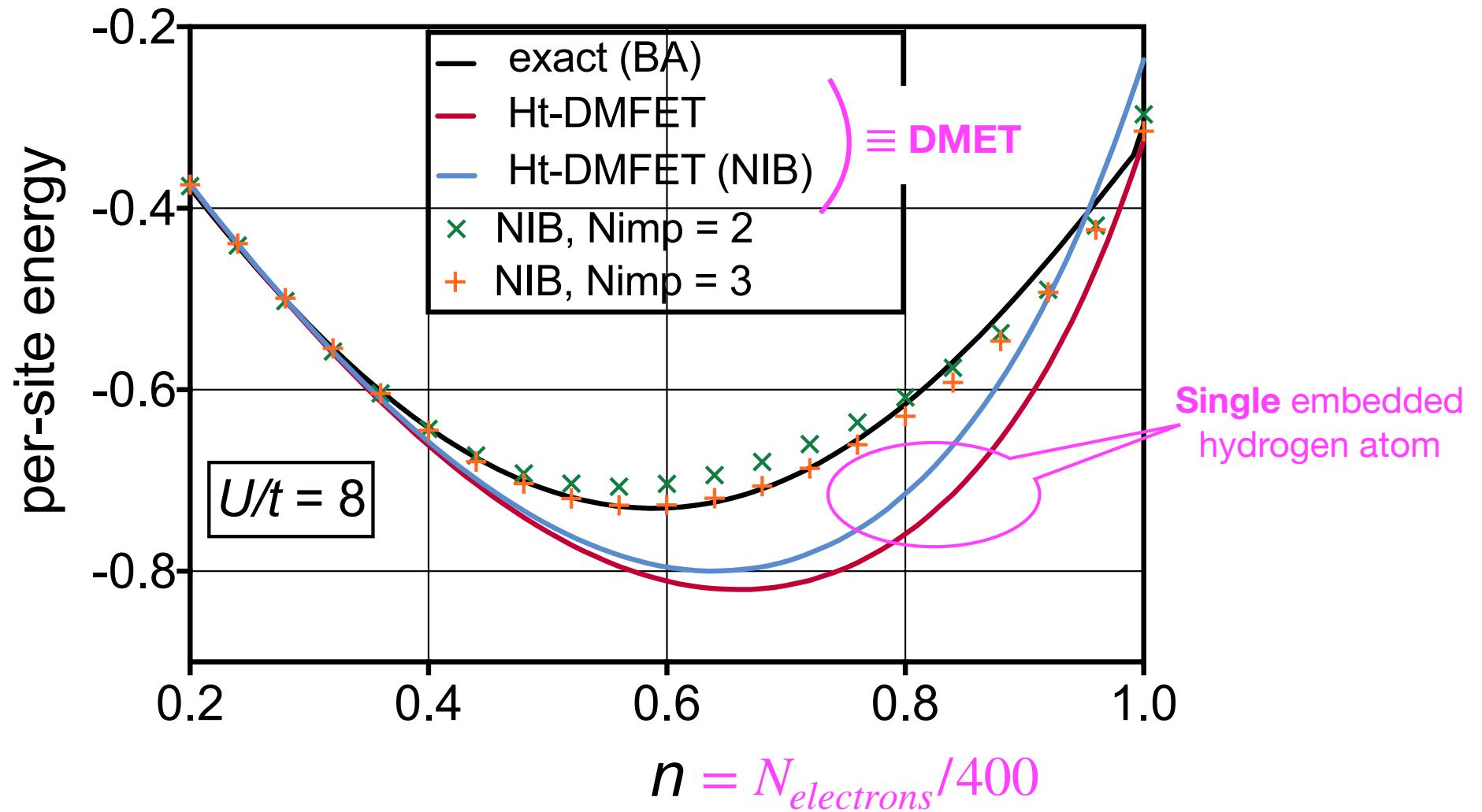
Correlated
cluster

$$\tilde{\gamma} = \begin{array}{|c|c|c|} \hline \gamma_{ff} & \tilde{\gamma}_{bf}^\dagger & \\ \hline \tilde{\gamma}_{bf} & \tilde{\gamma}_{bb} & \\ \hline & & \tilde{\gamma}_{ee} \\ \hline \end{array}$$

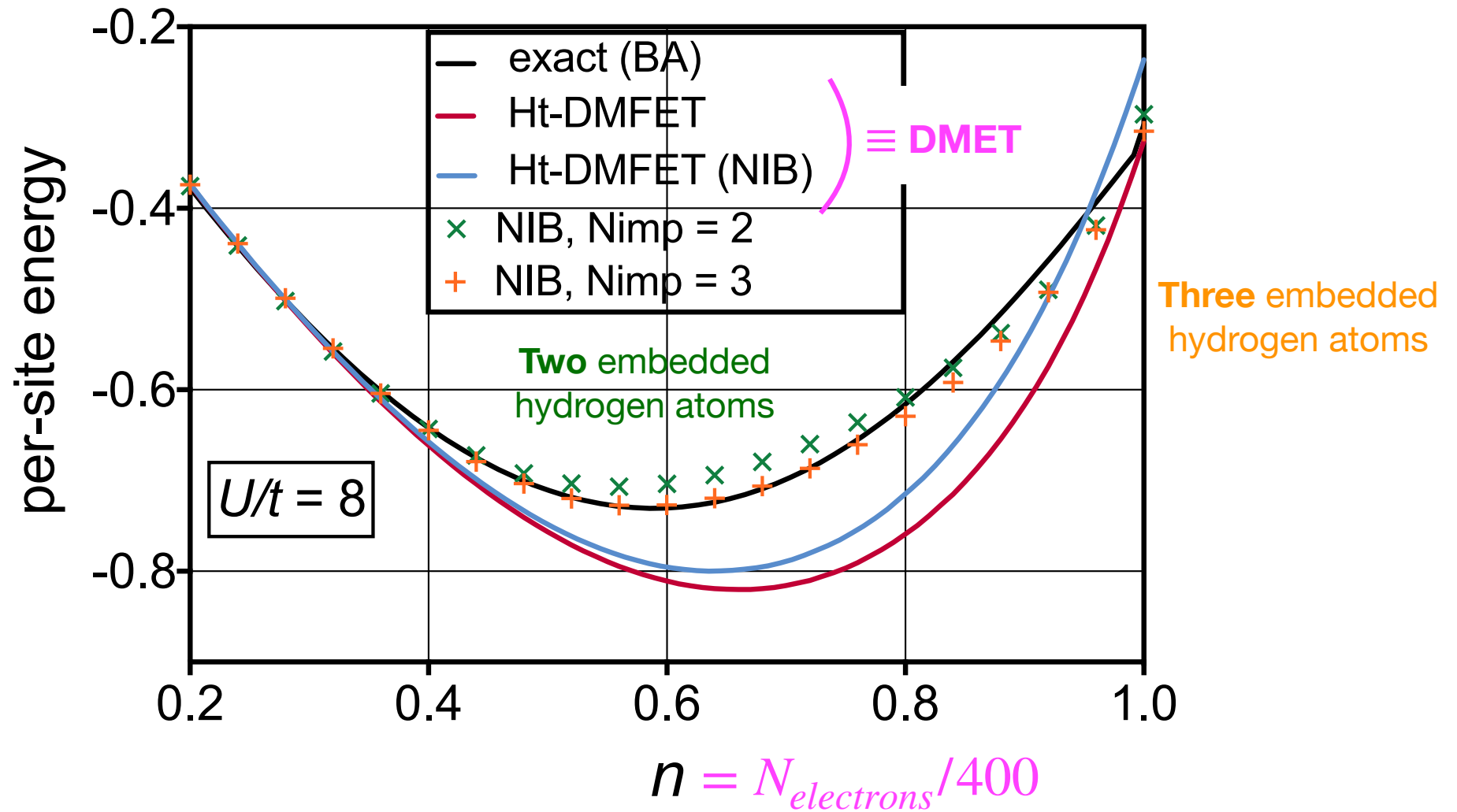
Non-idempotent

Complements

(Hubbard) model of a stretched 400-atom hydrogen ring



(Hubbard) model of a stretched 400-atom hydrogen ring

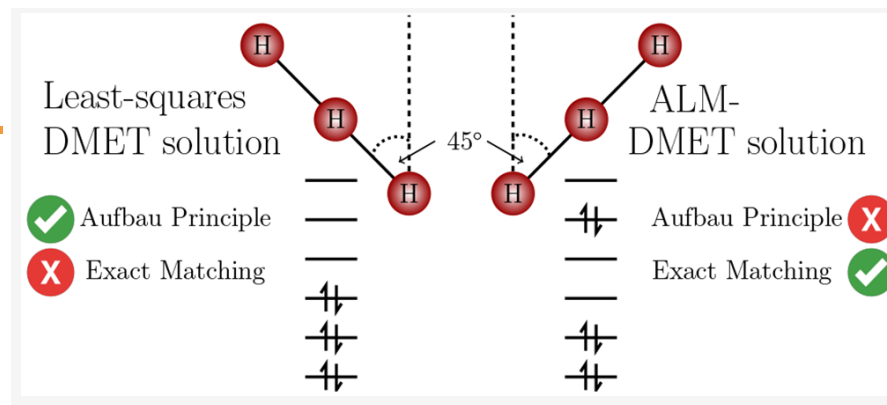


Pure State v -Representability of Density Matrix Embedding Theory

Fabian M. Faulstich,[⊥] Raehyun Kim,[⊥] Zhi-Hao Cui, Zaiwen Wen, Garnet Kin-Lic Chan, and Lin Lin*



Cite This: *J. Chem. Theory Comput.* 2022, 18, 851–864



Kohn-Sham-like
full system?

$$\gamma = \begin{array}{|c|c|} \hline \gamma_{ff} & \gamma_{ef}^\dagger \\ \hline \gamma_{ef} & \gamma_{ee} \\ \hline \end{array}$$

Idempotent

?

↔

Density matrix mapping

Correlated
cluster

$$\tilde{\gamma} = \begin{array}{|c|c|} \hline \gamma_{ff} & \tilde{\gamma}_{bf}^\dagger \\ \hline \tilde{\gamma}_{bf} & \tilde{\gamma}_{bb} \\ \hline \end{array}$$

Non-idempotent

Non-Hermitian quantum mechanics?

Non-Hermitian but idempotent density matrix, static self-energy, ...

<https://www.youtube.com/watch?v=8zgMa-MhoZg>

<https://www.youtube.com/watch?v=mDkzmSJwwkQ&t=726s>



Another approach to
*one-electron reduced
density matrix functional theory...*

Non-Hermitian quantum mechanics?

Non-Hermitian but idempotent density matrix, static self-energy, ...

<https://www.youtube.com/watch?v=8zgMa-MhoZg>

<https://www.youtube.com/watch?v=mDkzmSJwwkQ&t=726s>



Non-Hermitian quantum mechanics?

Non-Hermitian but idempotent density matrix, static self-energy, ...

<https://www.youtube.com/watch?v=8zgMa-MhoZg>

<https://www.youtube.com/watch?v=mDkzmSJwwkQ&t=726s>



Using an enlarged bath (ghost orbitals)?

N. Lanatà, Phys. Rev. B 108, 235112 (2023).

C. Mejuto-Zaera, Faraday Discuss., 2024, DOI: 10.1039/D4FD00053F

Quantum embedding for molecules using auxiliary particles – The ghost Gutzwiller Ansatz[†]

Carlos Mejuto-Zaera^{*a}

Faraday Discussions

Accepted Manuscript

DOI: 10.1039/D4FD00053F

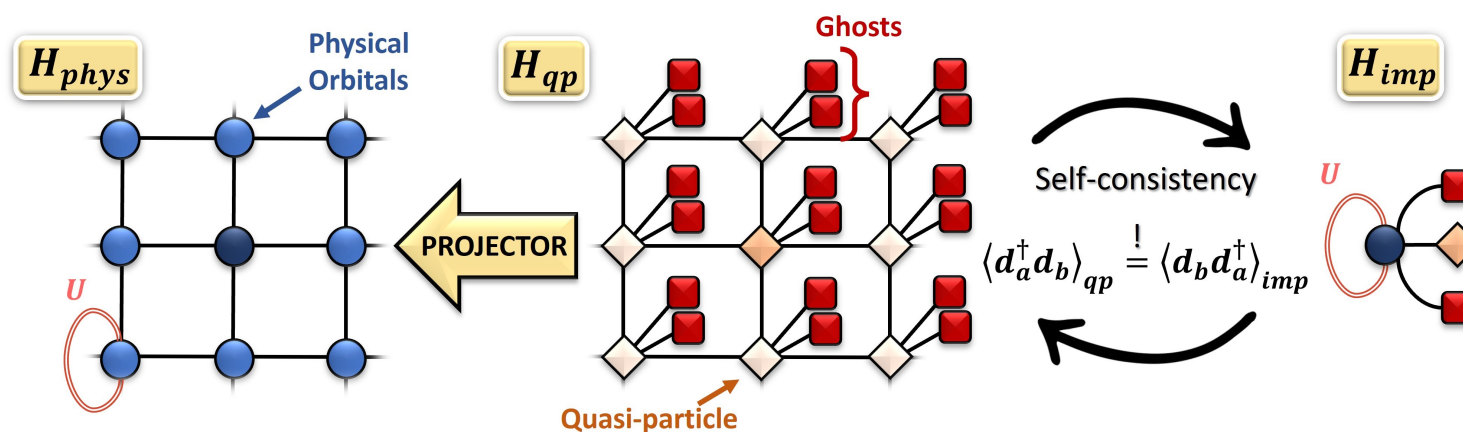


Fig. 1 Schematic representation of the Gutzwiller and ghost Gutzwiller Ansatz. Gutzwiller corresponds to the limit with no ghost orbitals. Local interactions are marked with the symbol U . See text for details.

Effective Reconstruction of Expectation Values from Ab Initio Quantum Embedding

Max Nusspickel, Basil Ibrahim, and George H. Booth*



Cite This: *J. Chem. Theory Comput.* 2023, 19, 2769–2791



Read Online

ACCESS |



Metrics & More

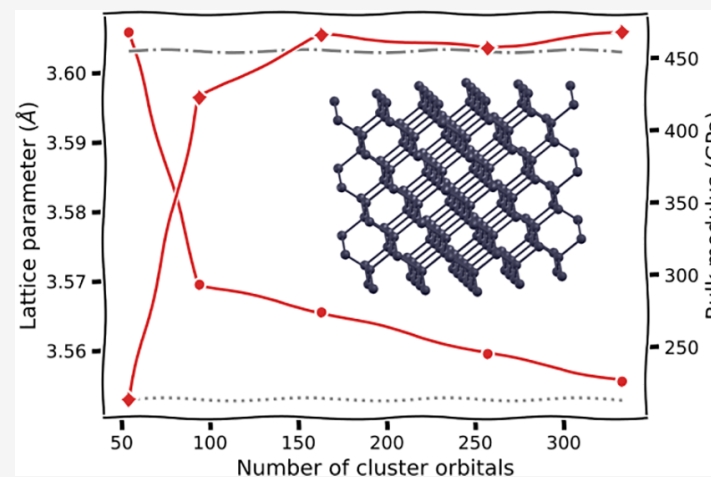


Article Recommendations



Supporting Information

ABSTRACT: Quantum embedding is an appealing route to fragment a large interacting quantum system into several smaller auxiliary “cluster” problems to exploit the locality of the correlated physics. In this work, we critically review approaches to recombine these fragmented solutions in order to compute nonlocal expectation values, including the total energy. Starting from the democratic partitioning of expectation values used in density matrix embedding theory, we motivate and develop a number of alternative approaches, numerically demonstrating their efficiency and improved accuracy as a function of increasing cluster size for both energetics and nonlocal two-body observables in molecular and solid state systems. These approaches consider the N -representability of the resulting expectation values via an implicit global wave function across the clusters, as well as the importance of including contributions to expectation values spanning multiple fragments simultaneously, thereby alleviating the fundamental locality approximation of the embedding. We clearly demonstrate the value of these introduced functionals for reliable extraction of observables and robust and systematic convergence as the cluster size increases, allowing for significantly smaller clusters to be used for a desired accuracy compared to traditional approaches in *ab initio* wave function quantum embedding.



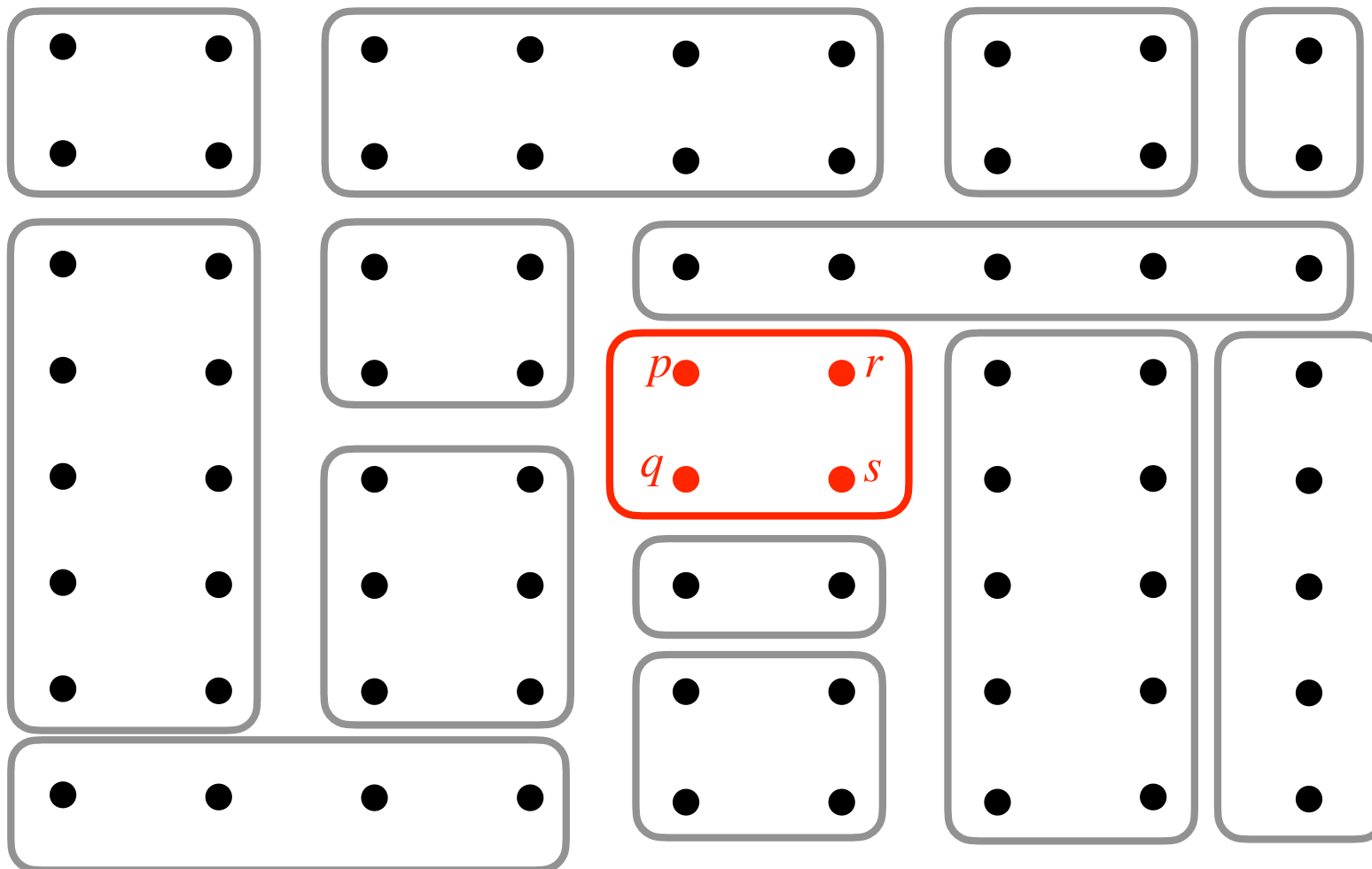
Local evaluation of the energy (in a **localised** spin-orbital basis)

So-called “lattice representation”

$$\langle \hat{H} \rangle = \sum_{pq} h_{pq} \langle \hat{c}_p^\dagger \hat{c}_q \rangle + \frac{1}{2} \sum_{pqrs} \langle pq | rs \rangle \langle \hat{c}_p^\dagger \hat{c}_q^\dagger \hat{c}_s \hat{c}_r \rangle$$

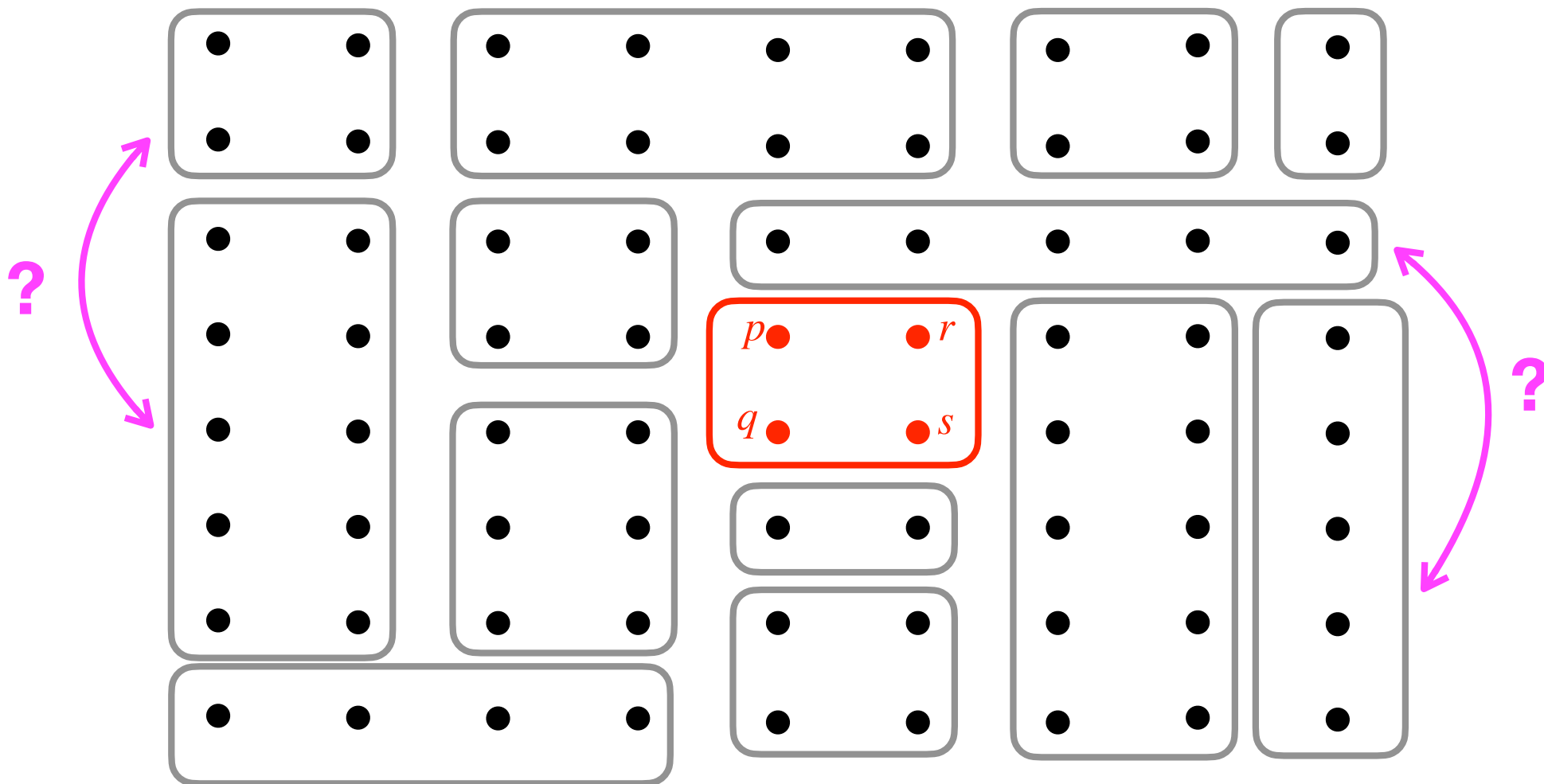
One-electron density matrix (1RDM)
Two-electron density matrix (2RDM)

Fragmentation



$$\langle \hat{H} \rangle = \sum_{pq} h_{pq} \langle \hat{c}_p^\dagger \hat{c}_q \rangle + \frac{1}{2} \sum_{pqrs} \langle pq | rs \rangle \langle \hat{c}_p^\dagger \hat{c}_q^\dagger \hat{c}_s \hat{c}_r \rangle$$

One-electron density matrix (1RDM)
Two-electron density matrix (2RDM)



N-representability problem

$$\langle \hat{c}_p^\dagger \hat{c}_q \rangle_{clusters} \stackrel{?}{=} \langle \Psi | \hat{c}_p^\dagger \hat{c}_q | \Psi \rangle$$

$$\langle \hat{c}_p^\dagger \hat{c}_q^\dagger \hat{c}_r \hat{c}_s \rangle_{clusters} \stackrel{?}{=} \langle \Psi | \hat{c}_p^\dagger \hat{c}_q^\dagger \hat{c}_r \hat{c}_s | \Psi \rangle$$

